

FOREWORD

This report was prepared by the University of Brussels, Belgium, under USAF Contract No. AF61(052)-225. The contract was initiated under Project No. 7350, "Refractory Inorganic Non-Metallic Materials," Task No. 735001, "Non-Graphitic." The work was administered under the direction of the Directorate of Materials and Processes, Deputy for Technology, Aeronautical Systems Division, Wright-Patterson Air Force Base, Ohio. Mr. F. W. Vahldiek was the project engineer.

ABSTRACT

The available spectroscopic and thermochemical data for the dissociation energy of the group IV-group VI MeX molecules are reviewed and the best present values proposed. Correlations of several excited molecular states with atomic products are proposed.

This technical documentary report has been reviewed and is approved.

W. G. RAMKE

Chief, Ceramics and Graphite Branch Metals and Ceramics Laboratory

Materials Central



TABLE OF CONTENTS

THE	DISSOCIATION ENERGY OF THE GROUP IV-GROUP VI MOLECULES]
	Thermochemical Data	2
	Spectroscopic Data	16
	Comparison of Thermochemical and Spectroscopic Data	2]
	Acknowledgments	27
	References	29

THE DISSOCIATION ENERGY OF THE GROUP IV-GROUP VI MOLECULES

Accurate convergence limits for one or more electronic states of most of the group IV (Me: C, Si, Ge, Sn, Pb)group VI (X: 0, S, Se, Te) molecules are known. This is in particular the case for CO (1), a molecule not further discussed here and SiS, SiSe, GeS, SnO, SnS, SnSe and PbS, molecules for which the best known convergence limit is that of the excited E state. Accurate spectroscopic values $D_{\alpha}^{"}$ for the dissociation energy of the ground state of these molecules can be derived therefrom, provided the correlation with atomic products is known. It has been postulated (2,3) that the E states of these molecules are analogous and correlate with the same atomic products. Recent rotational analysis for Pb0⁽⁴⁾ and Sn0⁽⁵⁾ showed that for the E state of these molecules, the most likely correlation is $Me(^{3}P_{1})+O(^{3}P_{1})$. Likewise the comparison of thermochemical and spectroscopic data for SnS and PbS (6) showed that their E state most probably correlates with $Me(^{3}P_{1}) + S(^{3}P_{1}, 0)$.

The purpose of this paper is to present a comparison between the spectroscopic data and recent thermochemical determinations of the dissociation energies of the group IV-group Manuscript released by the authors February 1963 for publication as an ASD Documentary Report.

VI molecules, several of which were performed mass spectrometrically and to derive in each case the best present value for the dissociation energy. The comparison gives further support to the correlation of the E state with $Me(^3P_1) + X(^3P_1)$, which is then applied to all molecules of this group for which reliable data for this state are known. Other available spectroscopic data are also examined.

THERMOCHEMICAL DATA

The thermochemical data were in general obtained mass spectrometrically or from total pressure measurements supplemented with mass spectrometric analysis of the composition of the vapor. The mass spectrometric technique (7) and the instruments (8-10) used have been described previously.

The system studied, the temperature intervals covered and the composition of the vapor are summarized in Table I. This table shows that above MeX compounds, including SiO, GeO and SnO known to disproportionate to 1/2 MeO₂(s)+1/2Me(s,1), the predominent molecule is MeX, polymers (and elements for GeTe, SnSe, SnTe, PbS, PbSe and PbTe) representing in general 10^{-1} to 10^{-4} in pressure, except for SnO and PbO, where (SnO)₁₋₄ and (PbO)₁₋₄ have comparable pressures. The MeX₂ compounds vaporize mainly to MeX(g)+1/2X₂(g). MeX₂ molecules were observed above SiO₂ (11), SiS₂, GeTe and SnTe.

Contrails

TABLE I: Gaseous Species in Equilibrium with Group IV-Group VI Compouds Standard Heats of Formation and of Sublimation

System	Composition of the vapour	Temperature range investi- gated (°K)	Ref.	-4H ⁹ 98,f kcal/mole	Ref.	AH998, subl(b) kcal/mole	1(b)
1/2 SiO ₂ +1/2 Si	Si0, Si ₂ 0 ₂	1200 - 1460	11	108.9*0.1	13	85.5±1.5	20
\sin_2	SiO, O, SiO ₂ , O ₂	1750 - 1950	11	217.7±0.2	13	198.5*1.5	21
Sis	Sis	850 - 1140		38,7 ±3,0		64.3±3.0	
Sis ₂	Sis, S ₂ , SiS ₂	850 - 1140		61.5±3.0		102.5±3.0	
GeO	(GeO,(GeO) ₂ ,(GeO) ₃)	!		58,9*2,5		50.6±2.5	
1/2GeO ₂ +1/2Ge	GeO, (GeO) ₂ , (GeO) ₃	750 - 950		64.2±0.3	1 4	55.9±1.6	
GeO ₂	GeO, 0 ₂	1020 - 1520		128,3±0,5	≒	(120.0*2.0)	
GeS	GeS	1		16,2±1,3	15	38,7±0,6	22
GeTe	GeTe, Te 2, GeTe 2	456 - 078		4.2*2.0	16	46.6±2.0	23,24
1/2SnO ₂ +1/2Sn	SnO,(SnO) ₂₋₄	1100 - 1400		E9.8±0.1	16	71.1*1.3	
Sn0 ₂	Sn0, 0 ₂ ,(Sn0) ₂	1200 - 1500		138.7±0.2	16	139,7*2,0	
SnS	SnS,(SnS) ₂	835 - 1005	9	25.1*1.2	17	52.6±1.6	9
SnSe	SnSe,(SnSe) ₂ , Se ₂	790 - 1010		19.9*2.0		51,1*2.0	a,25
SnTe	SnTe, Te, Te2, SnTe2,	830 - 990		14.6±0.3	18	52.5±1.6	a,25
Pb0	Pbo, (Pbo),	940 - 1180		52.4±0.2	19	66.5*2.5	
PbS	PbS, (PbS), Pb, S2	865 - 1140	9	22.5±0.5	19	55.7*1.6	9
PbSe	PbSe, Pb, Se ₂	1180 - 1450	12	18.0*2.0	18	53.9±1.7	26
PbTe	$PbTe_{b}b_{J}e_{1}e_{2}$	923 - 1171	12	16.6±0.5	18	53.8±2.0	27



Legend of Table I:

- a. this work
- b. ΔH° refers to the equilibrium 298, subl

 $MeX(s) \rightarrow MeX(g)$ or

 $MeX_2(s) \rightarrow MeX(g)+1/2X_2(g)$

The dissociation energies $D_{O}^{O}(\text{MeX})$ were calculated from thermochemical cycles of the type

Additional determinations based on the equilibrium $\text{MeX(g)} \rightarrow \text{Me(s, 1 or g)} + 1/2 \text{ X}_2(\text{g}) \text{ were made for GeTe, SnSe}$ SnTe, PbS⁽⁶⁾, PbSe⁽¹²⁾ and PbTe⁽¹²⁾.

The heats of sublimation $\Delta H_{\rm sub}({\rm MeX})$ are summarized in Table II. The heats of sublimation of the group IV elements Me (28) and of the group VI molecules X_2 as well as the heats of formation of the compounds MeX or MeX₂ were in general taken from the litterature. The values used for the dissociation energies $D_{298}^{\rm O}$ of D_2 , D_2 , D_2 , D_3 , and D_4 were D_2 , D_3 , D_4 , D_5 , D_5 , D_6 , D_7 , D_8 , D_8 , D_8 , D_8 , D_9 ,

(It may be noted that the agreement between spectroscopic and thermochemical values of the dissociation energies of the group of molecules discussed in the following chapters is an indirect confirmation of the $D_O^O(S_2)$, $D_O^O(S_2)$ and $D_O^O(Te_2)$ values used).



For all the gaseous species Me, X_2 and MeX, values for free energy function given in the literature (28,32) were used. For those stable compounds MeX and MeX, for which values are not given in the literature (32,33) estimates were made as follows. For solid SnTe, PbSe and PbTe (32), the entropies at 298°K are known. For these compounds only the heat capacity, while for SiS, GeTe and SnSe both the entropy at 298°K and the heat capacity were estimated from the corresponding quantities for the elements using approximations discussed by Kubaschewski and Evans (18). The corresponding quantities for GeS were derived from thermodynamic data for the reactions GeS(s) \rightarrow GeS(g) (22) and GeS(s) \rightarrow Ge(s) + 1/2 S₂ (15). The values obtained for the free energy functions are at 298°K, 12.0(SiS), 18.6(GeS), 20.3(GeSe), 22.2(GeTe), 22.5 (SnSe) and at 900°K 16.2(SiS), 23.0(GeS), 24.8(GeSe), 26.8 (GeTe), 28.0(SnSe), 29.8(SnTe), 30.0(PbSe) and 31.8(PbTe) cal/degree/mole.

The thermochemical data for each molecule are briefly discussed below

CS Thermochemical determinations (34) of the heat of formation of CS(g) were based on the reactions

C(graphite) + CS₂(g) + 2 CS(g)
MnS(s) + C(graphite) + CS(g) + Mn(g)

Using the latter reaction for which the enthalpy change $\Delta H_{298}^{o} \text{ was measured to be } 178.5 \pm 0.5 \text{ kcal/mole}^{(34)}, \ \Delta H_{298}^{o}, f^{(MnS(s))} = 49.5 \pm 0.5 \text{ kcal/mole}^{(35)} \text{ and } \Delta H_{298}^{o}(\text{sub Mn}) = 66.6 \pm 1.0 \text{ kcal/mole}^{(35)}$

mole $^{(28,30)}$, Freeman $^{(37)}$ calculated a value of 62.5 ± 3.0 kcal/mole for the standard heat of formation of CS(g). This leads to a value D_0^0 (CS)=173.6±3.5 kcal/mole $^{(7.53\pm0.15 \text{ eV})}$.

- CSe Thermochemical data were obtained (38) for the reaction $C(graphite)+Se(g) \rightarrow CSe(g)\Delta H_0^0=-31.4\pm2.0 \text{ kcal/mole}$ whence $D_0^0(CSe)=138.4\pm2.5 \text{ kcal/mole}$ (6.00±0.10 eV)
- CTe This molecule has hiterto not been observed. Its dissociation energy is estimated here as D₀=111±9 kcal/mole (4.8±0.4 eV) by comparison with the other IV-VI molecules.
- Sio The mass spectrometric data (11) confirm earlier conclusions that the main vaporization processes are

1/2 $SiO_2(s) + 1/2 Si(s,1) + SiO(g)$ I $SiO_2(s) + SiO(g) + 1/2 O_2(g)$ II

The heats of these reactions were taken as $\Delta H_{298}=85.5$ $\pm 0.5^{(20)}$ and $\Delta H_{298}^{\circ}=198.3\pm 1.5^{(21)}$ kcal/mole respectively. To derive the heat of formation of SiO(g), the reaction $^{(20)}$ SiO₂(s)+H₂(g) + SiO(g)+H₂O(g)(III), $\Delta H_{298}^{\circ}=134.8\pm 2.0$ kcal/mole was further taken into account. (For a more complete review, see Kubaschewski and Evans and Schick $^{(35)}$). The determination for reaction I and III are in good agreement with one another, while those for reaction II lead to slightly higher values. The recently redetermined standard heat of formation of SiO_2 , $\Delta H_{298f}^{\circ}(SiO_2(s))=-217.7\pm 0.2$ kcal/mole $^{(13)}$ was used

and the heat of sublimation of silicon taken as $\Delta H_{298,sub}^{0}$ (Si) = 107.9±1.0, this being the average of several determinations (40). The resulting mean value for the dissociation energy of SiO(g) is $D_{0}^{0}(SiO) = 191.9\pm3.0 \text{ kcal/mole } (8.32\pm0.13 \text{ eV}).$

No independant thermochemical data for this molecule were yet obtained. The spectroscopic value for the dissociation energy D" = 6.38±0.06 eV was used to calculate the heats of formation of SiS(s) and SiS₂(s) from mass spectrometric study of mixtures of SiS and SiS₂ and from total pressure for SiS₂(41). The heats of formation obtained in this manner from the reactions

SiS(s) \rightarrow SiS(g) $\Delta H_{298}^{\circ} = 63.4 \pm 3 \text{ kcal/mole}$ SiS₂(s) \rightarrow SiS(g)+1/2S₂(g) $\Delta H_{298}^{\circ} = 101.5 \pm 3 \text{ kcal/mole}$ are ΔH_{298f}° (SiS(s)) = -38.8 \pm 3.0, and ΔH_{298f}° (SiS₂(s))= -61.9 \pm 3.0 kcal/mole

The heat of formation of condensed SiS shows this compound not to be metastable up to about 1250°K, in agreement with observations by Retzlaff and Kohlmeyer (42). The heat of formation of SiS₂ is quite different from the discordant values given in the literature, which have anyway to be corrected for the recent values for the heat of formation of SiO . The average value

selected by Kubaschewski and Evans $^{(18)}$, 61.6±6.0 kcal/mole (corrected for $\Delta H_f(\mathrm{SiO}_2)$) is in agreement with the present value. The latter also makes it possible to explain the apparent boiling points of SiS_2 at about 1400°K as being that of the Si + SiS_2 system.

GeO Earlier thermochemical values were based on vapor pressure determination for GeO (43) and Ge + GeO (14,43)

The evaporation behavior of Ge + GeO and GeO were examined mass spectrometrically. The evaporation processes are

 $1/2\text{Ge}(s) + 1/2\text{GeO}(s) + 1/n(\text{GeO})_n(g) \quad n=1,2,3$ $\text{GeO}_2(s) + \text{GeO}(g) + 1/2O_2(g)$

The mass spectrometric results for GeO_2 obtained here are in disagreement with an earlier investigation ⁽⁴⁴⁾, also by mass spectrometry, which seems to be characteristic of $Ge + GeO_2$ rather than of GeO_2 . The enthalpy for the reaction

 $1/2 \text{GeO}_2(s) + 1/2 \text{Ge}(s) + \text{GeO}(g) \Delta \text{H}_{298}^{\circ} = 55.9 \pm 1.6 \text{ kcal/mole}$ was calculated from mass spectrometric pressure measurements made independently, as well as from the total pressures given in the literature (14,43) correcting these for the presence of $(\text{GeO})_2$ and $(\text{GeO})_3$. A similar procedure was adopted for GeO(s), for which a heat of formation $\Delta \text{H}_{298}^{\circ} = -58.9 \pm 3.0 \text{ kcal/mole}$ was

derived from the earlier data (43) in satisfactory agreement with direct determinations (45). The best average of all data gives $D_0^0(\text{GeO}) = 156.6 \pm 3.0 \text{ kcal/mole}$ (6.79 \pm 0.13 eV). The enthalpy of the reaction $\text{GeO}_2(s) \rightarrow \text{GeO}(g) + 1/2 O_2(g)$ was not measured. As a result of interaction with SiO_2 crucible the pressures in the literature (46) are low. The enthalpy given in Table I was calculated from the heat of formation of $\text{GeO}_2(s)$ and the above dissociation energy of GeO(g).

GeS The mass spectrometric analysis of the vapor showed only the monomer GeS in the gas phase.

The enthalpy of the reaction

GeS(s) \rightarrow GeS(g) ΔH_{298}° = 38.7 \pm 0.6 kcal/mole is hence very well known from a number of pressures measurements which are all in very good agreement except that of Kenwonthy and coworkers which seems to be a factor of approximately 3 too high. The heat of formation of GeS was evaluated from the reaction (15)

 $2 \text{ GeS(s)} + 2 \text{ Ge(s)} + S_2(g)$

and yields

 ΔH_{298f}° (GeS) = -16.2±1.3 kcal/mole Inserting these values in the cycle given above, one obtains D_0° (GeS) = 133.3±3.5 kcal/mole (5.78±0.15 eV).

- Liu Ch'Un-Hua, Pashinkin and Novoselova (24)

 Liu Ch'Un-Hua, Pashinkin and Novoselova . Using the spectroscopic value of the dissociation energy D"(GeSe) = 113.1±0.7 kcal/mole discussed below, it can be shown that the decomposition pressure of Se₂ is low compared to that of the main species GeSe(g). The total pressure measurements yield therefore a value of the heat of sublimation ΔH° (GeSe) = 298sub 44.6±1.5 kcal/mole obtained by End (44.2) and 3rd law (44.8) procedures. Combining this value with D"(GeSe), a heat of formation ΔH° (GeSe(s)) = -14.4±3.0 kcal/mole is obtained.
- GeTe The mass spectrometric investigation of the decomposition of the vapor showed the presence of GeTe,

 Te₂, (GeTe)₂ and GeTe₂. The pressure of the Te₂ molecule, relative to that of the main species GeTe(g)

 is 10⁻¹ at 880°K. Using total vapor pressure measurements (23,24) to evaluate the partial pressures,

 a heat of sublimation ΔH° (GeTe) = 46.6±2.0

 kcal/mole and a heat of formation ΔH° (GeTe(s)) =

 -4.2±2.0 kcal/mole were calculated from 3rd law procedures (2rd law: ΔH° (GeTe) = 47.8 (23), 46.7 (24)).

 These values lead to the dissociation energy D° (GeTe)=

 92.4±3.0 kcal/mole (4.00±0.13 V.)



SnO The heats of both reactions

 $SnO_2(s) \rightarrow SnO(g) + 1/2O_2(g)$ and $1/2Sn(1) + 1/2SnO_2(s) \rightarrow 1/n(SnO)_n(g)$ n=1,2,3,4were studied mass spectrometrically. The values obtained are $\Delta H^0_{298} = 139.7 \pm 2.0$ and 71.1 ± 1.3 kcal/mole (for n=1) respectively. Whence $D^0_0(SnO) = 127.2 \pm 2.0$ kcal/mole (5.53 ± 0.10 eV).

Total pressure determinations were examined for comparison. The vapor pressure measurements of $\mathrm{Sn0}_2(s)$ by the Knudsen method (47) were not included, since the sample interacted with the crucible material (16) $\mathrm{Si0}_2$. The pressure determination by the transport $\mathrm{method}^{(16)}$ gives too high values because the polymers $(\mathrm{Sn0})_{2-4}$ contribute to the weight loss. Estimating their contribution from the mass spectrometer data yields $\mathrm{D}_0^0 = 130.5 \pm 3.0$ kcal/mole (5.64 ± 0.13 eV) in agreement with the direct determinations.

SnS The thermochemical value for the dissociation energy based on mass spectrometry and vapor pressure data is $D_0^0(SnS) = 110.1 \pm 3.0 \text{ kcal/mole (4.78 \pm 0.13 eV)}$. For a detailed discussion see reference (6).

Since The composition of the vapor as well as partial pressures were determined mass spectrometrically. The pressure of the main component, Since, was found in good agreement with pressure determinations by the effusion method (25). These data lead to a value for the heat of sublimation: AHO (Since)=51.1±2.0 kcal/298sub mole (3rd law). From the partial Se pressure, a heat of formation AHO (Since) = -19.9±1.5 kcal/mole was derived compared to the estimated value -16.5±2 kcal/mole (18). These values lead to DO(Since)=94.8±4.0 kcal/mole (4.11±0.17eV).

SnTe The composition of the vapor and partial pressures were determined as above. The SnTe pressure was again in good agreement with total pressure measurements by the effusion method (25) The Te₂ partial pressure gave a value of -14.7±2.0 for the heat of formation of SnTe(s), in good agreement with AH = -14.6±0.3 kcal/mole (18), thereby confirming the values of the free energy function of SnTe(s) used. The 2nd and 3rd law values, 47.5 and 52.3 kcal/mole respectively, for the heat of sublimation are in less good agreement. Preference is given to the 3rd law value, which leads to D°(SnTe) = 79.8±3.0 kcal/mole (3.46±0.17 eV).

Pb0 The mass spectrometric analysis of the vapor above
Pb0(s) showed the presence of (Pb0)
1-4. The enthalpy
change for the reaction

 $PbO(s) \rightarrow PbO(g)$

was calculated as being ΔH_{298} = 65.8±1.5 kcal/mole from total pressure measurement given in the literature (48) taking the composition of the vapor into account. This value combined with the heat of formation of PbO(s), ΔH_{298f} = -52.4±0.2 kcal/mole (19) leads to $D_0^0(PbO)$ = 92.2±3.0 kcal/mole (4.00±0.13 eV).

- Mass spectrometric (6) and vapor pressure data combined with $\Delta H_{298f} = 22.5 \pm 0.5$ kcal/mole (18) lead to $D_0^0(PbS) = 78.9 \pm 0.25$ kcal/mole (3.42 ± 0.12 eV). For a detailed discussion see reference (6).
- The knowledge of the dissociation energy of this molecule is not satisfactory. The mass spectrometric investigation of the vapor by Porter (12) shows PbSe to be the main gaseous species. Hence one calculates from the published pressure measurements (24) a heat of sublimation of 53.9 kcal/mole (2nd law: 53.2), which combined with the estimated heat of formation ΔH^O_{298f}(PbSe)=-18.0±2.0 kcal/mole (18) gives D^O(PbSe)=65.0 kcal/mole (2.82 eV). Direct measurements made in double oven type Knudsen cells by Porter (12) give

for the reaction PbSe(g) \Rightarrow Pb(g) + Se(g) values ranging from 60.0 to 64.5 kcal/mole, with an average of 62.9 kcal/mole. Simultaneous measurements for the reaction Pb(g) + Se₂(g) \Rightarrow PbSe(g) + Se(g) give however $\Delta H_{298}^{0} = 2.7 \pm 1.0 \text{ kcal/mole}$, which combined with $D_{298}^{0}(\text{Se}_{2}) = 75.4 \pm 2.0 \text{ kcal/mole}$ (30) would give $D_{298}^{0}(\text{PbSe}) = 72.7 \pm 3.0 \text{ kcal/mole}$ ($D_{0}^{0} = 72.0 \pm 3.0$). Until a direct determination of the heat of formation of PbSe(s) is available, an average value $D_{0}^{0}(\text{PbSe}) = 63 \pm 5 \text{ kcal/mole}$ (2.95 \pm 0.22 eV) is selected.

PbTe The dissociation energy of PbTe was measured mass spectrometrically by Porter⁽¹²⁾, who obtained D_o⁰(PbTe) = 51.4±2.0 kcal/mole. This figure is in agreement with the value D_o⁰(PbTe) = 55.0±2.0 kcal/mole (2nd law: 52.9) that can be derived from the vapor pressure measurements⁽²⁵⁾ (main component PbTe) and the heat of formation of PbTe(s), $\Delta H_{298f}^{0} = -16.6 \pm 0.5$ kcal/mole ⁽¹⁹⁾. The average value D_o⁰(PbTe) = 53.1±3.0 kcal/mole (2.30±0.17 eV) is selected.

The thermochemical values D_0^0 for the dissociation energies of the group IV-group VI molecules have been assembled in Table V.

SPECTROSCOPIC DATA

The spectroscopic data considered in the present paper have been collected in table II and were all taken from the literature. Only the following comments seem necessary for the purposes of this paper. For the molecules SiS, GeS, GeSe, GeTe, SnO and SnSe graphical extrapolations for the E state could be made. For GeS, GeSe, Sno, SnS, SnTe and PbS, the extrapolation were confirmed or supplemented by direct observation of convergence limits (continua) for the same or other states. For GeTe, Pb0 and possibly PbSe the value given for the convergence limit of the E state corresponds to the wavelength down to which vibrational structure lines could be observed. For the molecule GeTe the value of the convergence limit is somewhat uncertain, since it was estimated by analogy with GeS and GeSe. It should however not be lower than 4.38 eV, since the vibrational levels are known up to that energy. For the molecule SiTe, for which the linear Birge Sponer limit for the E state was not calculated because of the presence of large anharmonic and cubic terms, the extrapolation was nevertheless made here to serve as a guide for obtaining a spectroscopic value for the dissociation energy, since no other data are available.

	j			TABLE II.		Spectroscopic Data				
Molecule		State voo	3	3 %	γs	ω ² /4×ω	Conve	Convergence L	Limit	
MeX		cm_1	cm ⁻¹	cm ⁻¹	cm-1	9	linear ev	graph. ev	observ. ev	Ref
SS	׫	38804.8	1285.1	6.5 10.3		7.88 3.46	7.88			6+
⊕ \$⊕	×∢	0 35134.6	1035.9	. 88 88		6.82	6,82			20
Sio	хоын	0 42640.4 52579.9 64497.2	1241.44 852.71 675 1116.5	5.92 6.44 4.15 7.22		8.08 3.50 3.40 3.40	8.07 8.78 9.91 13.34			ო
Sis	ХQЫ	0 34910.1 41750.8	749.5 513.12 403.54	2.56 2.93 1.40	-0.0329*	6.92 2.79 3.61	6.92 7.12 8.79	44.9)		51
9 Sis 1 7	хош	32360.2 38370.3	580°,0 399°,8 308°,8	1.78 1.93 1.95	-0.032	5.86 2.56 1.52	5.86 6.57 6.28			, k
SiTe	хаы	0 28590.4 33871	481.2 338.6 242	1.30 1.70 (3.68)	1.30 1.70 (3.68)(+0.13)	5.52 2.09 (0.50)	5.52 5.64 (4.70)			6
0 e 0	хоыч	37599.5 49396.5 67385	985.5 650 504.3 809.3	4.23 4.23 4.8 5.66		7.02 3.12 1.65	7.01 7.77 7.76 11.95			ო
ഗ ഉ	×AM	0 32589.2 38752.0	575.8 375.0 310.3	1.80		5.70 2.89 2.08	5.70 6.96 6.89	5.79	5.79	5.4
	×	zw = 1.15								



TABLE II. (Continued)

1 cm-1 cm-1 cm-1 cm-1 ev	Molecule	State	00°	3	æx	yε	m׆/m	Conve	Convergence Lin	Limit observ.	Ref.
GeSe X 0 408.7 1.36 3.80 3.80 3.80 5.25 E 5.53 6.35 5.25 E 5.56 6.35 5.25 E 5.56 6.35 5.25 E 5.56 6.35 5.25 E 5.56 5.60 5.25 1.02 1.02 1.71 5.15 5.15 6.04 9.4.38 E 5.367.1 217.7 1.02 1.71 6.157 6.25 5.62 5.62 5.62 5.62 5.62 5.62 5.62			ı	1	cm-1	cm ⁻¹	e۸	ev	94 to 1	ev	
GeTe X 0 323.9 0.25 5.98 5.25 5.25 GeTe X 0.89 0.25 5.98 5.98 5.98 5.98 5.98 5.98 5.98 5.9	\ \frac{1}{2}	×	0	08	1.36			•			
GeTe X 0 323.9 0.25 5.98 5.98 5.25 5.25 5.00 2.25 5.98 5.98 5.98 5.98 5.98 5.98 5.98 5.9)	; c	0776.	69	ထ			•		4	
GeTe X 0 323.9 0.25 5.98 5.98 5.98 5.98 5.98 5.98 5.98 5.9		ы	5367.	17.	•		•	•	• 5		
Sno X 0 822.4 3.73 5.62 5.62 5.62 5.62 5.62 5.62 5.62 5.62	Ę	>	c	23	0.25		0				54
Sno X 0 8224 3.73 5.62 5.62 5.62 5.62 5.63 5.63 5.63 5.63 5.64 5.74 (4.67) 4.49 \$4.38 5.60 5.65 5.62 5.62 5.62 5.62 5.62 5.62 5.62	ט פעד	۲,	7 0 0	יר היי	0 C		~				
Sno X 0 822.4 3.73 5.62 5.62 5.62 5.62 5.62 5.62 5.62 5.62			31393)	70.	(1.2)		. 7	•		ຕຸ	
SnS X 0 331.2 0.736 1.765 5.06 SnS X 0 331.2 0.736 5.696 5.906 5.906 5.906 5.906 5.906 5.906 5.906 5.906 5.006	O C	×	0	22.	7		φ.	9•			ო
SnS X 0 331.2 0.736 4.54 5.06 6.40.6.40 SnS X 0 185.9 1.358 5.39 5.39 5.39 6.596 5.00 E 32939.6 295.05 1.09 -0.012 2.49 6.58 6.06 E 2697.0 331.2 0.736 4.63 4.54 E 30571.6 125.6 5.36 4.54	:		25325	75)							
E 36138 508.0 2.9 2.76 7.23 5.68 6.40-6.44 D 57491 F 58752 724 21 SnS			29505.	82.	•			•	.0±0		
SnS X 0 485.9 1.358 5.39 5.39 5.39 SnS X 0 485.9 1.358 5.39 5.39 5.39 E 32939.6 295.05 1.09 -0.012 2.49 6.58 5.06 P 52174.5 (395) T 56470-890 SnSe X 0 331.2 0.736 4.63 4.63 A 19299.3 223.2 0.88 B 21795 (220) C 22695 (220) C 22695 (220) C 22695 (220) E 30671.6 196.6 0.77 -0.0016 1.56 5.36 4.54		ш	6138	08.	•		-	• 2	ဖ္	()	=======================================
SnS X 0 485.9 1.358 5.39 5.39 5.39 SnS X 0 485.9 1.358 5.39 5.39 5.39 E 32939.6 295.05 1.09 -0.012 2.49 6.58 5.06 E 52174.5 (395) G 55928 T 56470-890 SnSe X 0 331.2 0.736 4.63 E 30671.6 196.6 0.77 8.05 E 2269 5.35 C 2269 5.35 E 30671.6 196.6 0.77 -0.0016 1.56 5.36 4.54			5177							0 - 0	† †
X 0 485.9 1.358 5.39 5.39 5.39 D 28180.8 330.08 1.265 2.51 6.01 5.01 E 52939.6 295.05 1.09 -0.012 2.49 6.58 5.06 Q 55928 1.09 -0.012 2.49 6.58 5.06 F 56470-890 56470-890 4.63 4.63 4.63 4.63 X 0 331.2 0.736 4.15 4.15 4.15 B 21795 (220) 2.28 5.69 4.54 C 22695 225.1 0.69 -0.0016 1.56 5.36 4.54 F (47830) (290) -0.0016 1.56 5.36 4.54		ды	875	2	21		7	0			
D 28180.8 330.08 1.265	S.	×	O	85.	35		•	က			5.5
E 32939.6 295.05 1.09 -0.012 2.49 6.58 5.06 %5.2174.5 (395) p 52174.5 (395) q 55928 r 56470-890 X 0 331.2 0.736 A 19299.3 223.2 0.88 B 21795 (220) C 22695 (220) D 27496.6 225.1 0.69 E 30671.6 196.6 0.77 -0.0016 1.56 5.36 4.54 F (47830) (290)	:	; A	8180.	30.0	. 26		•	0			•
P 52174.5 (395) Q 55928 T 56470-890 X 0 331.2 0.736 A 19299.3 223.2 0.88 B 21795 (220) C 22695 (220) D 27496.6 225.1 0.69 E 30671.6 196.6 0.77 -0.0016 1.56 5.36 4.54 F (47830) (290)		щ	2939.	95.0	٥.	-0.012	•	S	o.		20
q 55928 r 56470-890 X 0 331.2 0.736 A 19299.3 223.2 0.88 B 21795 (220) C 22695 (220) D 27496.6 225.1 0.69 E 30671.6 196.6 0.77 -0.0016 1.56 5.36 4.54 F (47830) (290)		ር	2174.	ი მ						σ	
X 0 331.2 0.736 4.63 4.63 A 19299.3 223.2 0.88 B 21795 (220) C 22695 (220) D 27496.6 225.1 0.69 E 30671.6 196.6 0.77 -0.0016 1.56 5.36 4.54 F (47830) (290)		_다 년	5928 6470-89	~						0	
A 19299.3 223.2 0.88 1.76 4.15 B 21795 (220) C 22695 (220) D 27496.6 225.1 0.69 E 30671.6 196.6 0.77 -0.0016 1.56 5.36 F (47830) (290)	, ,	>	c	5	7.3		ပ	9			57,58
21795 (220) 22695 (220) 27496.6 225.1 0.69 2.28 5.69 30671.6 196.6 0.77 -0.0016 1.56 5.36 (47830) (290)	מווים	< ∢	9299.	223.	00		. 7	٦.			
27496.6 225.1 0.69 2.28 5.69 30671.6 196.6 0.77 -0.0016 1.56 5.36 (47830) (290)		щc	179	22							
30671.6 196.6 0.77 -0.0016 1.56 5.36 (47830) (290)		ם כ	7496.	225.	0.69		• 2	9			
		ыr	0671.	96	0.77	-0.0016	٠.	ო.	च • च		

TABLE II. (Continued)

Molecule State V	Sta	ate Voq	3	w X		ω2/4×ω	Collinear	Convergence	Limit observ.	Ref.	1
		E D		cm_T	Cm L	6	e >				
SnTe	×	0	59.			4.18	4.18			57.58	1
	¥	6803.	œ	•		2	ന	•			
	m	0380.	30.	Š	-0.013	0	ယ္	(3.34) ^D			
	ပ	1397.	18.	σ,		S.	Ч.				
	A	5404.	79.			ᠴ.	ဖ				
	Щ	7580.	35.	•		~	9		(4.07)		É
	ധ	8465.	8			സ	ထ		•		9
	I	045	8	0.3		4.17	7.82				j G
	н	0789.	01.			0	0				ŧ,
	F4	4018.	229.7	_	-0.003	m	_	6.56			Ž
	ה	2 tr 5	30								'n
Pbo	×	0	21.	. 7		က	6			4	Œ,
	∢	20088	451.7	3.33		1,89	4.38			•	É
19	ď										12
•	മ	22531	£65	2.26		3,33	•				E.
	ပ	410	က	6		• 2	(5.26)				
	ပ်	572	475)								
	Ω	9 t 0	630,5	(2.92)		(5.99)	(6,77)				
	Щ	1 68	+ + +						>4.86	59	
PbS	×		28.1	.20			4.74			90	
	A	18768.0	261,09	0,365		5.80	٦.))	
	щ	1774.	82.1	. 85			5				
	ပ	3150.	03.9	÷ ±3			ω.				
	ပ်	4952.	83.9	.17			5.24				
	A I	9586.	99,3	.57		•					
	리 [나	47739		(7.8)					4.55	59	



TABLE II. (Continued)

							Conv	Convergence Limit	Limit	
Molecule	St	State	3	3	γĸ	ω ² /4×ω	linea	graph	observ.	Ref.
		cm ⁻¹	G E	cm cm	T Eo	e v	e >	ev	ev	
									r Il:	69.63
ር ነ	×	0	77.	0.51						•
2	. ∢	8661.	66.	0.14						
	<u>, na</u>	0959	• ₩±	0.43						
	ပ		183.0	0.25		5 t • 1	10.7	q(36 ii)	q(
	Ω	74.	90.	0.53	+00.00-			*	5.7 ≥4.32	
	ltj [ri	45194.5	224.8	0.50		3.13	8.73			
1	;	4	0	7		3,17				59,62
PDTe	< <	ב פ פ	1 C	201		4.77	7.05			
20	C A		α - 1	91110		J. 46				
0	ם ב		10.0			0,40				
) t	• 7 + 4 - 7	י טנ) -		0.92				
	ų U	45515.5 46515.5	159.6) 		0 0 0				
	}	 -								

a. In most cases the letter label is that of the ref. quoted; it may be different from that used in ref.58. b. Numerically calculated using the formula $E=\omega v+\kappa\omega v^2+y\omega v^3$.



COMPARISON OF THERMOCHEMICAL AND SPECTROSCOPIC DATA.

This comparison will be started with the E state, since this is the only state for which in general accurate values for the convergence limit are available.

It had already been concluded, on the basis of the agreement between the convergence limits for this state and the Birge Sponer value for the dissociation energy of the ground state, that this state correlates with $Me(^3P) + X(^3P)^{(3)}$.

Inspection of the multiplet splittings given in Table III, shows that, in order to decide on the basis of a comparison of spectroscopic and thermochemical data, which Me(³P₀,1,2) and X(³P₂,1,0) sublevel are involved the more favorable cases would be CSe, CTe, SiSe, SiTe and perhaps GeSe and GeTe for X(³P) and SnO, SnS, PbO and PbS for Me(³P). Unfortunately, the available data preclude such comparison for CSe, CTe, SiSe, SiTe and GeSe. This comparison can however be made for GeTe, SnO, SnS, PbO, PbS and also SnSe, SnTe and to a certain extent PbSe.

It has already been shown for SnS and PbS that such a comparison leads to the conclusion that the E state correlates with $Me(^3P_1) + X(^3P_{1,0})^{(6)}$.



TABLE III. Excitation energy of ${}^3P_{0,1,2}$ and 1D_2 states of the elements of Group IV and VI(in ev).

	³ P ₀	³ P ₀	³ P ₂	¹ B ₂
С	0	0.002	0.005	1.27
Si	0	0.01	0.03	0.78
Ge	0	0.07	0.18	0.89
Sn	0	0.21	0.43	1.07
Рb	0	0.97	1.32	2.68
0	0.03	0.02	0	1.97
S	0.07	0.05	0	1.15
Se	0.31	0.25	0	1.19
Te	0.58	0.58	0	1.31

This conclusion is confirmed by a similar comparison for the molecules SnO, SnSe and PbO and to a certain extent GeTe, SnTe and PbSe, as shown in Table IV.

The only molecule for which there seems to be a discrepancy between the dissociation energy so derived from the spectroscopic data of the E state and the thermochemical value is SnTe. It may however be noted that the spectroscopic data for the E state of all tellurides are not satisfactory. For SnTe only a region of apparently continuous absorption is known and interpreted as a possible convergence limit of the E state leading to the value 4.07 e.V. which according to the authors (57) "should be treated with considerable reserve".

The accuracy of the thermochemical data makes it difficult to distinguish between the $X(^3P_1)$ and $X(^3P_1)$ sublevels, the separation of which is maximum 0.06 e.V. for X = Se. From the retational analysis of several states of the molecule PbO, Barrow, Deutsch and Travis (4) have concluded on the basis of the Ω correlation rule and the non-crossing rule, that the most likely correlations for several 0^+ states of this molecule are as indicated in Table IV. Professor Barrow has kindly informed us that a similar analysis for SnO also gives as most likely correlation for the E state, $Sn(^3P_1) + O(^3P_1)$.

The interpretation of the rotational analysis and the comparison of the thermochemical and spectroscopic data are thus in agreement and in fact confirm and supplement one another.

The conclusion that the E state of the group IV-group VI molecules correlate with Me(³P₁) + X(³P₁) has therefore been extended in Table V to the molecules SiS, SiSe, SiTe GeS and GeSe for which accurate convergence limits of the E state and no thermochemical data are available. The dissociation energies of SiS and GeSe so obtained were already used in the section dealing with thermochemical data to derive the heats of formation of the corresponding compounds.

If the correlation of the E state thus seems established it is less easy to derive the correlations of the other states. The fact that most combinations of $Me(^3P_1)$ + $X(^3P_1)$ sublevels can lead to one of several Ω = 1 molecular states make it difficult to propose off hand a correlation for the B(1) and D(1) states. The convergence limit of the B state for SnTe, 3.34 e.V., calculated from the formula $E = \omega v + x\omega v^2 + y\omega v^3$ with $\omega = 230.3$, $x\omega = 1.53$, $y\omega = -0.013 cm^{-1}$ would indicate that this state correlates with $Sn(^3P_0) + Te(^3P_1)$. The B state in PbTe, known up to 2.65 e.V. can however not correlate with $Pb(^3P_0) + Te(^3P_1)$ at 2.23±0.10 e.V. and should at least correlate with $Pb(^3P_0) + Te(^3P_1)$ or any combination



 $Pb(^3P_1)$ + $Te(^3P_{2,1,0})$. If the convergence limit of SnTe is considered to be low, but not off by more than 0.5 e.V; and if the B states of both molecules is analogous as would be indicated by the same ratio ω_B/ω_X , then the probable correlation of the B state would be $Me(^3P_0)$ + $X(^3P_1)$. (Table IV).

TABLE IV. Correlation of the X, A, B, C, D and E States of the Group IV-Group VI MeX molecules.

Atomic Me	States X	Possible Ω values	likely mole cular states
3 _P 0	3 _{P2}	0 ⁺ , 1, 2	x(o ⁺)
-	3 _{P1}	0-, 1	B(1)?
	³ P ₀	o +	A(0 ⁺)
3 _P 1	3 _{P2}	0-(2), 0+, 1(3), 2(2), 3	c(o ⁺)
	3 _P 1	0-, 0+(2), 1(2), 2	E(0 ⁺)
	3 PO	0, 1	D(1)?
³ P ₂	3 _{P2}	$0^{-}(2), 0^{+}(3), 1(4), 2(3), 3$	(2),4
	3 _P 1	0-(2), 0+, 1(3), 2(2), 3	
	³ P ₀	0 ⁺ , 1, 2	

It is even less easy to speculate about the correlation of the D state. In PbTe, the difference between the convergence limit of this D state, 3.67*0.20 e.V. obtained from a linear Birge Sponer extrapolation and the thermoche-



mical value for the dissociation energy $D_0^{0}=2.30\pm0.17$ is 1.37 ± 0.37 e.V. It may be noted that the uncertainty quoted on the Birge Sponer extrapolation corresponds to an uncertainty as high as *70%, which still would not affect the conclusion since the linear dissociation energy of the D state itself is only 0.4 e.V. This difference compared to the excitation energies of the different $^{3}P_{0,1,2}$ + $^{3}P_{2,1,0}$ combinations for Pb and Te, 0, 0.58, 0.58, 0.97, 1.55, 1.55 1.32, 1.90 and 1.90 e.V., indicates that possible correlations are $Pb(^{3}P_{1}) + Te(^{3}P_{1,0})$ or $Pb(^{3}P_{2}) + Te(^{3}P_{2})$. A selection between those two correlations cannot be made on the basis of the highest vibrational level of the D state observed in any of the MeX molecules. A possible argument for the ${}^{3}P$ + ${}^{3}P$ correlation could be that the linearly extrapolated convergence limit of the D state lies either above or below the similarly evaluated limit of the E state when the excitation energy $Me(^{3}P_{1}) + X(^{3}P_{1,0})$ is either larger or smaller than $Me(^{3}P_{2}) + X(^{3}P_{2})$ for 8(SiSe, SiTe, GeSe, GeTe, SnO, SnS, SnSe, SnTe) out of 12 molecules of which two (GeO, GeS) have both limits essentially at the same level and two (SiO, SiS) present the inverse situation.

The F states, have already been shown (3) to correlate with $Me(^1D_2) + X(^1\mathfrak{p}_2)$. They have not been observed for all molecules of this group and will not be discussed here. For SnSe the combination $Sn(^1D_2) + S(^1D_2)$ leads to

one shallow and one repulsive state, from which accurate values of the dissociation energy can be derived (see footnote c, Table V).

For SnTe, higher molecular states G, H and I are known. Perhaps could these states be correlated with $\mathrm{Sn(^3P_2)} + \mathrm{Te(^3P_{2,1,0})}$. The term values ν_{oo} increase regularly for the A to I states and an abrupt increment for the F state. The same situation occurs for the other molecules of this group as can be seen seen from Table II.

ACKNOWLEDGMENTS

The authors wish to thank Professor P. Goldfinger for his interest and encouragement, Professor R. Barrow for communicating his results on SnO and for a stimulating correspondance, Professor F. Klanberg for his results on the vapor pressure of SiS and Mr. G. Verhaegen for assistance with the mass spectrometric experiments.



TABLE V. Comparison of spectroscopic D_O^n and thermochemical Dovalue of the dissociation energies of group IVgroup VI molecules (in e.V.)

Molecule	Convergence limit E state	Excitation energy Me(³ P ₁)+X(³ P ₁)	D"	Do
CO CS CSe CTe			11.111#0.06	7.53±0.15 6.00±0.15 (4.8 ±0.4)(a
SiO SiS SiSe SiTe	6.44±0.06 5,66±0.15 (4.70±0.25)	0.06 0.26 0.59	6.38±0.06 5.40±0.15 (4.11±0.25)	8.32 ±0.15
Ge0 GeS GeSe	5.79±0.03 5.25±0.09	0.12 0.32	5.67±0.03 4.93±0.03	6.79±0.13 5.78±0.15
GeTe	4.49 +0.30 -0.10	0.65	3.84+0.30	4.00±0.15
Sn0	5.68±0.05	0.23	5.45±0.05 5.37(b)	5.53±0.13
SnS	5.06*0.10	0.26	4.80±0.10 4.74(c) 4.83(c)	4.78±0.12
SnSe SnTe	4.54±0.06 (4.07)?	0.46 0.79	4.08±0.06 (3.28)?	4.11 [±] 0.17 3.45 [±] 0.17
Pb0 PbS PbSe PbTe	>4.86 4.55±0.20 >4.32 ?	0.99 1.02 1.22	> 3.87 3.53±0.20 > 3.10?	4.00*0.13 3.42*0.12 2.95*0.22 2.30*0.17

a estimated

c from shallow q and repulsive r state respectively, assuming both to correlate with $Sn(^{\circ}D_2)$ + $S(^{\circ}D_2)$

b from absorption continuum at 1926 to 1937 Å, correlated with Sn(1D2)+0(3P). It was assumed here that the O(3P2) sublevel is involved. If the absorption occured from v"=1, which could be the case as a result of the large interatomic equilibrium distance in the shallow was a state of the large interatomic equilibrium distance in the shallow was at the same state. librium distance in the shallow upper state, the value obtained D"=5.47 eV would be in perfect agreement with D"=5.45 eV.

REFERENCES.

- (1) A.G. Gaydon, "Dissociation Energies", Chapman and Hall Ltd (1953).
- (2) E.E. Vago and R.F. Barrow, Vol.Comm. Victor Henry, Liège Desoer (1948).
- (3) R.F. Barrow and H.C. Rowlinson, Proc.Roy.Soc., London, A 224, 374 (1954).
- (4) R.F. Barrow, J.L. Deutsch and D.N. Travis, Nature, 191, 374 (1961).
- (5) R.F. Barrow, private communication.
- (6) R. Colin and J. Drowart, J.Chem. Phys., 37, 1120 (1962).
- (7) M.G. Inghram and J. Drowart in "High Temperature Technology" (McGraw Hill Book Co, Inc., New York, 1960).
- (8) J. Drowart and R.E. Honig, J. Phys. Chem., 61, 980 (1957).
- (9) J. Drowart and P. Goldfinger, J.Chem. Phys., 55, 721 (1958).
- (10) M. Ackerman, F.E. Stafford and J. Drowart, J.Chem. Phyd. 33, 1784 (1960).
- (11) R.F. Porter, W.A. Chupka and M.G. Inghram, J.Chem.Phys. 23, 216 (1955).
- (12) R.F. Porter, J.Chem.Phys., 34, 583 (1961).
- (13) C.N. Cochran and L.M. Foster, J.Phys.Chem., 86, 380 (1962); W.D. Good, J.Phys.Chem., 66, 380 (1962); S.S. Wise, J.L. Margrave, J.L. Feder and W.N. Hubbard, J.Chem.Phys. 66, 381 (1962).
- (14) W.L. Jolly and W.M. Latimer, J.Amer.Chem.Soc., 74,5757 (1953).
- (15) K. Ono and K. Sudo, Sci.Repts.Research Inst., Tohoku Univ. (A) 7, 385 (1955).
- (16) J.C. Platteeuw and G. Meyer, Trans.Faraday Soc., <u>52</u>, 1066 (1956).
- (17) A.W. Richards, Trans. Faraday Soc., <u>51</u>, 1193 (1955).

- (18) O.Kubaschewski and E.LL. Evans, "Metallurgical Thermochemistry", Pergamon Press (1958).
- (19) F.D. Rossini, D.D. Wagman, W.H. Evans, S. Levine and I. Jaffee, Nat.Bur.Standards, Circ.500 (1952).
- (20) H.F. Ramstad, F.D. Richardson and P.J. Bowles, Trans. AIME, 221, 1011 (1961).
- (21) A.N. Nesmayanov and L.P. Firsova, Zhur.Fiz.Khim. 34, 1032 (1960).
- (22) R.F. Barrow, P.G. Dodsworth, G. Drummond and E.A.N.S. Jeffries, Trans. Faraday Soc., 57, 1480 (1955); U.I. Davydov and N.P. Diev, Zhur.Neorg.Khim., 2, 2003, (1957); H. Spandau and F. Klanberg, Z.anorg.u.Allgem. Chemie, 295, 291 (1958); K. Sudo, Sci.Repts.Research Inst., Tohoku Univ., (A) 12, 54 (1960); H. Kenworthy, A.G. Starliper and A. Ollar, U.S. Bur.Mines Rept.Invest. 5190 (1956).
- (23) C. Hiramaya, J. Phys. Chem., <u>66</u>, 1563 (1962).
- (24) Liu Ch'ün-Hua, A.S. Pashinkin and A.V. Novoselova, Zhur.Neorg.Khim, 7, 965 (1962).
- (25) Ya.M. Nesterova, A.S. Pashinkin and A.V. Novoselova, Zhur.Neorg.Khim., 6, 2014 (1961).
- (26) M.V. Lonomosov, Zhur. Neorg. Khim., 4, 2657-60 (Through C.A. 16971 g(1960)).
- (27) U.P. Zlomanov, B.A. Popovkin and A.N. Novoselova, Zhur.Neorg.Khim., 4, 2661-4 (1959).
- (28) D.R. Stull and G.C. Sinke, Advances in Chem. Ser. 18 (1956).
- (29) D.G. Maraden, J.Chem.Phys., 31, 1144 (1959); L. Brewer J.Chem.Phys., 31, 1143 (1959); R. Colin, P. Goldfinger and M. Jeunehomme, Nature, 187, 408 (1960).
- (30) This laboratory, unpublished work.
- (31) R. Colin, Ind.Chim.Belg., 26, 51 (1961).
- (32) K.K. Kelley, U.S. Bureau of Mines, Bull. 584, (1960); K.K. Kelley and E.G. King, U.S. Bureau of Mines, Bull 592 (1961); J.P. Coughlin, Contribution to the data on theoretical metallurgy, XII, U.S. Gov. Printing Office, Washington (1954).

- (33) JANAF Interim. Thermochemical Tables. The Dow Chemical Co., Michigan (1961).
- (34) H. Schäfer and H. Wiedemeier, Z.anorg.Allgem.Chem., 290, 841 (1958); "Chemical and Thermodynamic Properties at High Temperature" XVIIIth Intern.Congress of Pure and Applied Chemistry, Montreal, Canada (1961).
- (35) A.D. Mah, U.S. Bureau Mines Rept. Invest. N°5600 (1960).
- (36) P.L. Woolf, G.R. Zellars, E. Foerster and J.P. Morris, U.S. Bur.Mines Rept. Invest. no5634 (1960).
- (37) R.D. Freeman, "Thermodynamic Properties of Binary Sulfides", Research Foundation Report n°60, Oklahoma State Univ. (1962).
- (38) M. Jeunehomme, Thesis, Université Libre de Bruxelles Brussels (1962).
- (39) R.L. Schick, Chem.Rev., 60, 331 (1960).
- (40) R.E. Honig, J.Chem.Phys., 22, 1610 (1954); J. Drowart and G. DeMaria in "Silicon Carbide"; J.R. O'Connor and J. Smilten, Eds.Pergamon Press, New York 1960;p.16; R.L. Batdorf and F.M. Smits, J.Appl.Phys., 34, 659(1961); S.G. Davis, D.F. Anthrop and A.W. Searcy, J.Chem.Phys. 34, 659 (1961); P. Grieveson and C.B. Alcock in "Special Ceramic" Heywood and Co, Ltd, London 1961.
- (41) H.S. Spandau and F. Klanberg, personnal communication.
- (42) E.J. Kohlmeyer and H.W. Retzlaff, Z.anorg.Chem. <u>158</u>, 743 (1959).
- (43) W. Bues and H. von Wartenberg, Z.anorg.allgem.Chem. 266, 281 (1951).
- (44) S.A. Shuchkarev and G.A. Semenov, Doklady Akad.Nauk. SSSR, 120, 1059 (1958).
- (45) Jolly and Latimer, quoted by L. Brewer, Chem. Rev. 52, 1-62 (1953).
- (46) V.I. Davydov, Zhur. Neorg. Khim. 2, 1460 (1957).
- (47) B.K. Veselovskii, J.Appl.Chem. URSS <u>16</u>, 397 (1943).
- (48) E. Feiser, Metal Erz., 26, 269 (1929); A.W. Richards,

- Trans. Inst. Mem. Met. 65, 151 (1956)(through ref.(18)); A.N. Nesmeyanov, L.P. Firsova and E.P. Işakova, Zhur. Fiz.Khim., 34, 1200 (1960); Zhur.Fiz.Khim. 34, 1699 (1960).
- (49) A. Lagerqvist, H. Westerlund, C.W. Wright and R.F. Barrow, Arkiv.Fysik, 14, 387 (1959).
- (50) R.K. Laird and R.F. Barrow, Proc.Phys.Soc., <u>66A</u>, 836 (1953).
- (51) A. Lagerqvist, N. Nilheden and R.F. Barrow, Proc.Phys. Soc., 65A, 419 (1952).
- (52) S.J.Q. Robinson and R.F. Barrow, Proc. Phys. Soc. <u>67A</u>, 95 (1954).
- (53) E.E. Vago and R.F. Barrow, Proc. Phys. Soc., 58, 533(1946).
- (54) G. Drummond and R.F. Barrow, Proc. Phys. Soc., <u>65A</u>, 277 (1952).
- (55) A.E. Douglas, L.L. Howe and J.R. Morton, J.Mol.Spectroscopy, 7, 161 (1961).
- (56) R.F. Barrow, G. Drummond and H.C. Rowlinson, Proc. Phys. Soc. <u>68A</u>, 685 (1953).
- (57) E.E. Vago and R.F. Barrow, Proc. Phys., 58, 707 (1946).
- (58) G. Herzberg, Spectra of Diatomic Molecules, Van Nostrand (1961).
- (58) E.E. Vago and R.F. Barrow, Proc. Phys. Soc., <u>59</u>, 449 (1947).
- (60) G.D. Rochester and H.G. Howell, Proc.Roy.Soc., London 148, 157 (1935).
- (61) J.W. Walker, J.W. Staley and A.W. Smith, Phys.Rev., 53, 140 (1938).
- (62) R. Grove and N. Ginsburg, Spectrochim. Acta, <u>16</u>, 730 (1960).