# NN ITERATIVE METHOD IN DYNAMIC ETRUCTURAL ANALYSES WITH \&ONPROPORTIOMAL DAKPING 

Wan T. Taai', Joseph T. Leang ${ }^{2}$


#### Abstract

A new method in dynamic analyses of structures with nonproportional damping is proposed. By decomposing the nonproportional damping matrix into two portions, the diagonal and offdiagonal, the iterative technique can be employed through use of the classic method of solving a large dynamic structural system with real modal coordinates. Explicitly, the diagonalized damping matrix is retained to form a system of discrete differential equations with proportional damping. The off-diagonal portion of damping forces is treated as a correcting forcing function. The iteration is to use the off-diagonal damping induced forces for the load correction in the subsequent computational step. This enables the structural responses to be simply determined while the effect of off-diagonal damping forces is included.


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# NN ITRRATIVE KBTHOD IN DYNAKIC 8TRUCTURAL ANALYEES WITH NONPROPRTIONAL DAMPING 

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## INTRODUCTIOA

This article presents a new method in dynamic analyses of structures with non-proportional damping. Computed iteratively, this method ensures highly accurate responses yet low cost analyses when the structures are subjected to dynamic forcing functions.

In dynamic analyses of space vahicles, the mass and stiffness matrices of each substructure are usually generated by different contractors. Each substructure may contain its own component modal damping acquired from component testing or empirical data. When all substructures are coupled together into a system for dynamic analyses, a difficulty arises. Generally, the system damping matrix cannot be transformed into a diagonal matrix by using the same transformation matrix as for generalizations of the system mass and stiffness matrices. Since the transformed damping matrix is not diagonalized, the structural responses cannot be determined by taking the advantage of using the real mode superposition technique. In order to avoid this difficulty, the off-diagonal elements of the transformed damping matrix are usually neglected in the wake of their smallness compared to the corresponding diagonal elements. Known as triple-matrix-product (TMP) method in the aerospace industry, this approach has been widely employed as a standard method in analyses of space vehicles. Using this approach, the interface loads are usually accurate for design purposes. However, the responses in some payload components may be grossly incorrect when the full scale payload-orbitar coupled systems are exercised.

The proposed iterative method is to improve the accuracy of payload responses yot to retain the advantage of using the modal superposition technique for cost saving. This method decomposes the transformed system modal damping matrix into two portions, a diagonal damping matrix and an off-diagonal damping matrix. The damping force induced by the off-diagonal damping elements are treated as a correcting force vector to modify the applied forcing function. This correction can be repeatedly applied until the results are within an acceptible range of error. Using this iterative approach, the desired goal can be reached and the impact to the currently applied TMP method can be minimized.

The iterative method in treating the nonproportional damping matrix has been applied by the first author of this paper to analyze small scale of structures since 1988. Independently, the same method given by Udwadia and Esfandiari (1990), may be the first article related to this method published in the open literature. The convergent characteristics of the method is the primary emphasis
of the article. prior to using the iterative method, several approaches have been proposed in treating structural systems with non-proportional damping. Primarily due to the development of nuclear power plants in the 1970s, replacements of the nonproportional damping matrix with diagonal matrix have been extensively investigated. Those which have been more commonly employed are: (1) using the diagonal elements of the nonproportional damping matrix, i.e., the triple-matrix-product (TMP) method, (2) replacing the non-proportional damping matrix by a diagonal matrix containing each element with a factor of the critical damping to each mode, i.e., the system damping method, and (3) obtaining each diagonal damping element by using the algebraic sum of the corresponding row of the non-proportional damping matrix. Errors acquired from these approximations had also been examined. Those who had been involved in these methodology developments included Clough and Moftahedi (1976), Cronin (1976), Duncan and Taylor (1979), Hasselman (1976), Thompson, Calkin, and Caravani (1974), and Warburton and Soni (1977.) A different approach in synthesizing the diagonal system damping elements from the component modal testing was given by Tsai (1989.) It is known that a nonproportionally damped systam can be completely generalized through the complex mode transformations. Many investigators have been involved in developments of this method, for instance, Beliveau and Soucy (1985), and Veletsos and Ventura (1986.) Although, the complex mode transformation is the exact method in treating nonproportionally damped structural systems, the real mode transformation seems to be still more favorable to most application engineers for two reasons: (1) It is less costly in the full scale transient analysis when an approximate method is employed. (2) It is easier to capture the image of physical behavior when the real mode transformation is applied. Therefore, the approximations developed in the 1970s are etill favorably used. The TMP and the system damping approaches have been particularly favored in the aerospace industry.

In addition to the derivation of the iterative method, this paper emphasizes on case applications of the method to a full scale payload/orbiter dynamic analysis, i.e., the IUS/TDRS payload for the 26th space transportation system (STS-26) manifest. The direct integration method is employed as the basis to substantiate the validity of the new method. The reason that the TMP method is inadequate for payload response computations is extensively discussed. Recommendations in transient analyses for nonproportionally damped structures are provided.

## ITERATIVE METHOD

Let $M, C$, and $K$ be the physically coupled mass, damping, and stifiness matrices, $P$ the applied forcing vector, $X$ the response vector, and dots the derivatives with respect to time; the governing
differential equation of the dynamic system is given by

$$
\begin{equation*}
M \ddot{X}+C \dot{X}+K X=P \tag{1}
\end{equation*}
$$

Eq. (1) can be solved in a simple manner by transforming the physical response coordinates $X$ into a ystem of generalized response coordinates $x$ by introducing a transformation matrix $\varphi$. Namely

$$
\begin{equation*}
x=\varphi x \tag{2}
\end{equation*}
$$

Through use of Eq. (2), along with the correlations

$$
m=\varphi^{T} M \varphi, \quad c=\varphi^{T} C \varphi, \quad k=\varphi^{T} K \varphi, \quad p=\varphi^{T} P \quad(3 a, b, c, d)
$$

Eq. (1) then reduces to

$$
m \ddot{x}+c \dot{x}+k x=p
$$

Where, $m$ is a unit matrix, $c$ is a tully populated matrix, $k$ is a diagonal matrix containing eigenvalues in the diagonal elements, and $p$ is a generalized forcing vector.

Eq.(1') could be easily solved by using the real mode superpositions if $c$ were a diagonal matrix. In order to take the advantage of expressing all responises with the superpositions of real modes, let $c$ be expressed by the sum of a diagonal matrix $c_{d}$ and an off-diagonal matrix $c_{0}$. Namely,

$$
\begin{equation*}
c=c_{d}+c_{0} \tag{4}
\end{equation*}
$$

Eq. (1') can then be rewritten by

$$
\begin{equation*}
m \dot{x}+c_{d} \dot{x}+k x=p-c_{0} \dot{x} \tag{5}
\end{equation*}
$$

Thus, all the coefficient matrices on the left hand side of Eq.(5) are diagonal. The contribution of each generalized mode can be directly determined without coupling with the other modes. The technique of modal superpositions can then be applied to simplify the analysis inasmuch as the forcing function vector on the right hand side is explicitly given. The industrial practices generally assume that the effect of $c_{0}$ is negligible since its elements are generally smaller than the corresponding elements of $c_{d}$. The approach using this assumption is the method commonly referred in the aerospace industry as the triple-matrix-product (TMP). In fact, the words triple-matrix-product (TMP) does not comprise any meaning of stripping the off-diagonal damping elements. Nevertheless, this article uses the commonly accepted definition that the TMP method implies the applications of the diagonalized damping matrix.

Now, the iterative method comes to play by treating the response vector $x$ on both sides of Eq. (5) as if they were independent at different stages of computations, $x_{n}$ and $x_{n-1}$, where $n$ is the number of iterations. Eq. (5) is then rewritten into the form

$$
\begin{equation*}
\dot{x}_{n}+c_{d} \dot{x}_{n}+k x_{n}=p-c_{0} \dot{x}_{n-1}, \quad n=0,1,2, \ldots \tag{6}
\end{equation*}
$$

The left hand side represents the classic modal system of equations whereas the right hand side is the forcing function corrected by the off-diagonal damping induced forces. Using this system of equations, the modal DOFs can be easily determined since every equation is associated with a single DOF.

In Eq. (6) , $\dot{x}_{-1}=0$ when $n=0$. The response $x_{0}$ is the TMP result. The off-diagonal damping force obtained from the TMP result is applied to modify the applied force and the first iterative response $x_{1}$ for $n=1$ is then determined from Eq. (6). Analogously, the second iterative response $x_{2}$ for $n=2$ is computed by using the correcting force obtained from the $x_{1}$ response. This procedure can be repeatedly applied until the $n$th iterative response $x_{n}$ reaches an acceptible accuracy.

The iteration can be stopped when the response converges to a desired accuracy. Using various examples of three degree-offreedom, Udwadia and Esfandiari showed that the results of six iterations were almost identical to the exact solutions. For practical applications, such a highly accurate result may not be necessary. The analysis may be terminated by setting an accuracy criterion that the converging rate of the modal accelerations is within a specified admissible error, $\varepsilon$. Namely,

$$
\begin{equation*}
\left|\frac{x_{n-1}}{x_{n}}-1\right|<\varepsilon, \quad n=1,2, \ldots \tag{7}
\end{equation*}
$$

For most design purposes, the modal accelerations may be accurate enough to assume an admissible error of 5\% (0.05). This value is suggested on the basis of common practices, not a sophisticatedly computed number. Experiences indicate that responses may generally be converged to $\varepsilon<0.05$ when two iterations are performed.

It must be noted that the provided error limit for the modal accelerations does not assure of the physical accelerations to be always within the same limit. However, their limits are generally agreeable to each other in most applications. It must also be noted that the error determined by Eq. (7) is to judge the acceleration computed in the $(n-1)$ th iterative analysis with the nth iterative acceleration as the basis. In fact, when the result of the nth iteration is used for final response evaluation, the accuracy is higher than the specified admissible error since the updated result is supposed to be judged with the $(n+1)$ th iterative values.

## APPLICATIONS OF ITERATIVE METHOD

The dynamic liftoff analysis for the STS-26 flight manifest was used to demonstrate applications of the iterative method. The manifest consists of the substructures: tracking data relay
satellite (TDRS), inner upper stage (IUS) booster, orbiter, solid rocket booster (SRB), and external tank (ET). The spacecraft TDRS was integrated to the IUS which was in turn secured to the orbter cargo bay through payload trunnions. The primary purpose of this flight was to deploy the payload IUS/TDRS.

Originally, all components were individually modelled as substructures, each was represented by hundreds or thousands of degree-of-freedoms (DOFs). After several stages of substructural coupling and condensations, the final model used for the liftoff dynamic analyses was 520 DOFs. The system was subjected to a dynamic forcing function, LR1200, one of the conditions being used for liftoff transient analyses. The analysis was performed over a range of 11 seconds to cover the complete liftoff event. To assure of obtaining an accurate result, the small time interval of 0.001 seconds was used over the entire time span of transient analyses. In the dynamic loads analysis, 1\% of the critical damping was assumed for all the IUS/TDRS payload modes. For the orbiter, $1 \%$ for frequencies below 10 Hz and $2 \%$ for frequencies above 10 Hz were assumed. No damping was assumed at the payload/orbiter interface DOFs. The analyses were performed by using various approaches for comparisons. The results are summarized in Tables 1-5 in which column 1 represents the items of interest. Columns 2-6 represent the minimum and maximum values obtained from various methods of transient analyses.

Initially, the dynamic system was analyzed by using the TMP method. The results are sumarized in column 2 of Tables 1-5 for various component responses. The accelerations along the $X$ direction at the tip of SA antenna ribs (node ll5) was excessively high, a minimum of -25.7 g and a maximum of 27.2 g as shown in column 2 of Table 1. Since this result was not acceptible and such a strong response was very unusual, a similar analysis was performed by using 1\% of the critical system damping. As shown in column 6 of Table 1 , the response for the same item reduces significantly to a minimum of -9.6 g and a maximum of 9.2 g . Clearly, the comparison between these two sets of results indicates a controversial conclusion that the system with a higher damping value has a stronger response than the one with a lower damping value. This conclusion violates the nature of mechanics that a structure has less responses with higher damping.

When the iterative method is applied, the peak accelerations of the same item become -6.8 g and 7.4 g for one iteration, and -7.2 g and 7.7 g for two iterations. These results are respectively shown in columns 3 and 4 of Table 1. They are significantly different from that of the TMP method. The responses of several other items are also significantly changed between the TMP and iterative methods. For instance, the $Y$-acceleration of the $C$-band reflector CG in Table 1 and the member force at the LTM row 149 in Table 2, their responses using the iterative method appear to be less than one-half the TMP results. The difference between these selected items reveal a fundamentally severe deviation between the TMP and iterative
methods, although the responses of many other items are in good agreement. Physically, the results obtained from the iterative method make better sense than from the TMP method when they are compared to those using the system damping method, as shown in column 6 of Tables 1-5.

To make sure that the iteratively computed results are reasonably accurate, the first iterative result is compared to the second one, i.e., column 3 compared to column 4 of Tables 1-5. Among all interested items, the maximum difference between these two iterative analyses is only $4.6 \%$ occured in the X -acceleration at the tip of SA antenna ribs. Furthermore, in order to assure of correct results obtained from the iterative analyses, a direct integration method using the fully populated damping matrix is also performed. The results are shown in column 5 of Tables 1-5. All iterative responses are in very good agreement with the direct integration results. In general, the results of the second iteration are much closer to that of the direct integration than those of the first iteration. In certain particular items such as the tip of SA antenna ribs, the second iterative value has slightly more deviation than the first iterative result when both are compared to the direct integration result. It just happens on the way of converging process to the final result, but does not indicate any inaccuracy in the iterative method.

A question has been raised regarding the magnitude of the interface loads between using the system damping and iterative methods. Specifically, the load of -8909 lbs in the $1 \%$ system damping analysis appear to be weaker than -9024 lbs of the iterative analysis for the X -interface load at $\mathrm{X}=1155.53$ inches (node 43.) This may not be surprised since the damping matrix established for the iterative method is more complex than that for the system damping matix. The one for the iterative analysis consists of zero damping value at the inteface nodes as well as $1 \%$ and $2 \%$ for the Craig-Bamptom form of substructure modelling. But the $1 \%$ system damping implicitly include damping values at the interface nodes as well as the other DOFs, as shown in the reversed expression of Eq. (3b). This explains the reason that some of the interface loads are stronger in the iterative analyses than in the analysis of using 1\% system damping. The discrepancy for most of other quantities appears to be in the right order that the responses using the system damping method (1\% damping) is slightly greater than those using the iterative method ( $1 \%$ and $2 \%$ damping for orbiter.)

## REASONS FOR THE RESPONSE DISCREPANCY OF TMP METHOD

The reason for such a significant deviation in the TMP method has been interpreted as the consequence of modal response superpositions from two modes that have two closely spaced modes. This can be explained by considering the modal contribution for an item at a particular time slice. For instance, the modal contributions
of the acceleration at the tip of $S A$ antenna ribs at $t=8.3$ second are shown in Figures 1 and 2 respectively for the TMP and oneiterative analyses. Although both modal contribution plots appear to be different, the major contributions occur at the same frequency of 25.4 Hz in either analyses. Near the interested frequency, the contributions from Figure 1 of the TMP approach are both negative whereas the contributions from Figure 2 of the iterative method are mixed with a negative and a positive value. Owing to this type of misrepresentation in the TMP modal contributions, the TMP computed responses become significantly different from those of the iterative analysis.

Although the above interpretation is mathematically correct, there remains a clout regarding the true driving source that causes such a strong impact to the component responses in the present illustration. A careful study indicates that the light weight flexible components are driven by wrong forces when the TMP approach is applied. That is the true reason to induce such a significant impact at the component responses. This can be directly interpreted by using Eq. (5). Since the generalized mass $m$ is a unit matrix, Eq. (5) can be rewritten into an alternate form to express the acceleration in terms of $p, k, c_{0}, c_{d}, x$, and $\dot{x}$. The acceleration for the ith modal DOF is given by

$$
\begin{equation*}
\bar{x}_{i}=p_{i}-k_{i} x_{i}-\left(c_{d}\right)_{i} \dot{x}_{i}-\sum_{j}\left(c_{0}\right)_{i j} \dot{x}_{j} \tag{8}
\end{equation*}
$$

In the TMP analysis, the last term associated with $c_{0}$ has been entirely neglected. This may be justified when the applied force $p_{i}$, stiffness force $k_{i} x_{i}$, and diagonal damping force $\left(c_{d}\right)_{i} \dot{x}_{i}$ are much greater than the total off-diagonal damping force. For the internal components of a payload, the DOFs are generally not subjected to any directly applied forces. Instead, the component DOFs are driven by the combined action of stiffness and damping forces. When the stiffness is relatively small like the $S A$ antenna ribs ${ }^{1}$, the associated stiffness force is small and the damping force becomes an important part of the driving force. Furthermore, the off-diagonal damping force becomes the dominated portion in the total driving force. Specifically, the sum of several hundreds of off-diagonal damping force components may override the stiffness and diagonal damping forces to influence the final responses of the transient analysis, although each off-diagonal damping force component may be small compared to the counterpart of the stiffness and diagonal damping forces.

The component stiffness is small compared to the other portion of the structural system. However, the component frequency may not be small since the corresponding component mass is usually small too.

The above interpretation can be substantiated by the TDRS member loads shown in Table 2. The item at the LTM row 149 has a peak member force reduced from 21.7 lbs in the TMP approach down to 11.9 lbs in the one-iterative analysis. Similarly, the small IUS motor (node 3457) has the peak Y-acceleration changed from 0.42 g to 0.23 g as shown in Table 4. The changes are up to 100\%. On the contrary, the changes in the interface loads at the bridge points between the IUS/TDRS and orbiter as shown in Table 3 and the orbiter bridge acceleration as shown in Table 5 are much less, a maximum of 13\% only. Therefore, the TMP may still be applicable if all substructural components are stiff. However, when the component is flexible, it may be subjected to a wrong driving force when the offdiagonal damping force is neglected. As a result, the component response is incorrect. It is particularly sensitive to the component of small mass since it is more responsive to any variation of the driving force.

CONCLUSIONS

1. The commonly referred TMP method has assumed that the offdiagonal damping elements are small and negligible; and that uses of the diagonal damping elements are sufficient to capture accurate component responses of structures. This has been proved to be incorrect. In fact, the TMP method may produce a structural response totally different from the true result. Therefore, the currently applied TMP method should not be used.
2. A new method using iterative procedure is proposed for transient analyses of dynamic structures with non-proportional damping. This method can provide an accurate result within small number of iterations. The validity of this method has been substantiated by using the direct integration method through the illustrative analyses for the STS-26 flight.
3. The proposed iterative method is cost effective. On the basis of analyses for various structure sizes, the cost of using each additional iteration is about $15 \%$ more than the cost of using the TMP approach. Generally, two iterations may result in an accurate response for design purposes. If two iterations are used, the expected computing cost may increase about 30\%. This cost is not a significant impact when it can assure of obtaining a reliable response in all payload components. Therefore, the iterative method is a viable approach for transient analyses when the transformed system damping matrix is non-proportional.

## RECOMMENDATIONS

To ensure accurate component responses of structures, the offdiagonal damping elements must not be neglected. In order to retain the off-diagonal damping without significantly increasing computational cost, the iterative method may be used.

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TABLE 1 STS-26 TDRS ACCELERATION ( $g$ )

|  |  |  | Diagon | TMP | 1 iter |  | 2 Iter |  | full dam | atrix | 1\% system | mping |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | ITEMS |  | Min | Max | Min | Max | Min | Max | Min | Max | Min | Max |
| $\stackrel{\stackrel{4}{8}}{\stackrel{1}{\bullet}}$ | SaL Antenna | * 13 x | -1.869 | 2.678 | -1.842 | 2.812 | -1.830 | 2.830 | -1.836 | 2.831 | -1.742 | 2.823 |
|  | SCL Antema | * 13 r | -2.732 | 2.759 | -2.755 | 2.818 | -2.759 | 2.824 | -2.761 | 2.825 | -2.729 | 2.782 |
|  | SGL Antemma | \# 132 | 0.300 | 3.355 | 0.220 | 3.288 | 0.216 | 3.290 | 0.216 | 3.289 | 0.189 | 3.352 |
|  | SGL Feed | * $15 \times$ | -2.868 | 3.590 | -2.850 | 3.774 | -2.842 | 3.782 | -2.843 | 3.789 | -2.798 | 3.652 |
|  | SGL Feed | * 15 Y | -2.192 | 2.094 | -2.226 | 2.128 | -2.233 | 2.137 | -2.236 | 2.139 | -2.181 | 2.057 |
|  | SGL feed | \# 152 | -2.380 | 4.473 | -2.298 | 4.353 | -2.306 | 4.350 | -2.309 | 4.351 | -2.317 | 4.394 |
|  | C-Band Antenna | * $17 \times$ | -4.607 | 4.534 | -4.067 | 3.831 | -3.973 | 3.785 | -4.012 | 3.800 | -3.952 | 3.905 |
|  | C-Band Anterna | \# 17 r | -1.628 | 1.618 | -1.208 | 1.240 | -1.207 | 1.240 | -1.207 | 1.239 | -1.217 | 1.249 |
|  | C-Band Antenna | \#172 | -2.883 | 5.601 | -0.532 | 3.593 | -0.524 | 3.602 | -0.532 | 3.605 | -0.564. | 4.039 |
|  | C-Band Reflector Cg | \# 18 Y | -5.371 | 5.482 | -2.297 | 2.268 | -2.300 | 2.275 | -2.301 | 2.282 | -2.652 | 2.941 |
|  | Top C-Band Antenna | * $20 \times$ | -17.004 | 15.478 | -14.429 | 13.682 | -14.471 | 13.738 | -14.391 | 13.689 | -14.779 | 14.034 |
|  | Top C-Band Antenna | * 20 Y | -7.164 | 7.223 | -4.977 | 5.120 | -5.002 | 5.144 | -5.002 | 5.138 | -5.034 | 5.261 |
|  | Top C-Band Antenna | * 20 z | -5.117 | 8.104 | -1.120 | 4.128 | -1.134 | 3.935 | -1.135 | 3.786 | -1.769 | 5.133 |
|  | Propellant Tank cG | \# 75 X | -0.814 | 1.028 | -0.730 | 0.912 | -0.728 | 0.911 | -0.728 | 0.911 | -0.721 | 0.917 |
|  | Propellant Tank cG | \# 75 r | -0.510 | 0.453 | -0.405 | 0.378 | -0.416 | 0.379 | -0.422 | 0.380 | -0.505 | 0.477 |
|  | Propel lant Tank CG | \# 752 | 0.160 | 3.347 | 0.160 | 3.354 | 0.157 | 3.360 | 0.158 | 3.358 | 0.130 | 3.442 |
|  | +Y Solar Panel Hinge | \# $83 \times$ | -4.582 | 3.507 | -3.064 | 2.399 | -3.001 | 2.377 | -2.898 | 2.370 | -3.557 | 3.091 |
|  | +Y Solar Panel Hinge | \# 83 r | -2.762 | 2.802 | -2.153 | 1.643 | -2.157 | 1.643 | -2.157 | 1.602 | -2.285 | $1.981^{\circ}$ |
|  | +Y Solar Panel Boom | \# $84 \times$ | -1.114 | 1.191 | -0.996 | 1.193 | -0.993 | 1.202 | -0.991 | 1.204 | -1.118 | 1.269 |
|  | +Y Solar Panel foom | \# 84 Y | -0.630 | 0.632 | -0.514 | 0.540 | -11.504 | 0.534 | -0.503 | 0.537 | -0.550 | 0.577 |
|  | +X SA Antenna Ribs | \#114 X | -15.354 | 16.508 | -8.434 | 10.6015 | -8.446 | 110.687 | -8.503 | 10.685 | -8.686 | 11.969 |
|  | +X SA Antenna Ribs | *114 Y | -10.779 | 11.661 | -13.079 | 11.731 | -13.151 | 11.865 | -13.198 | 11.948 | -20.363 | 20.347 |
|  | +X SA Antenna Ribs | \#115 X | -25.678 | 27.177 | -6.820 | 7.402 | -7.152 | 7.748 | . 7.010 | 7.303 | -9.588 | 9.243 |
|  | +X SA Antenna Ribs | \#115 Y | -11.695 | 12.841 | -12.517 | 14.1087 | -12.565 | 14.248 | -12.590 | 14.302 | -15.463 | 15.945 |

TABLE 2 sTg-26 TDRS MEMBER FORCES (lbs)

|  | $1$ |  | 2 |  | 3 |  | 4 |  | 5 |  | 6 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | Dlagan | $1{ }^{\text {ThP }}$ | $1{ }^{11}$ | Ion | 211 | lons | Full 0 | -hatrix | 12 syatee | Omplns |
|  |  |  | Hin | Hax | Min | мак | Min | Hax | Min | Hax |  | Ma |
|  | Ltw row | 7 | -31.352 | 167.415 | 0.742 | 157.726 | 0.716 | 158,092 | 0.617 | 158.080 | -2.211 | 161.663 |
|  | LTM Row | 8 | -96.063 | 97,823 | -87.604 | ${ }^{88} .240$ | -87.553 | ${ }^{88.306}$ | -87.588 | ${ }^{86.383}$ | -86.971 | 88.046 |
|  | LTM Row | 9 | -239.302 | 24.8886 | -246.731 | 251.355 | -244. 525 | 250.255 | -244.563 | 250.720 | -248.046 | 247.068 |
|  | LtM Rour | 10 | -3779.852 | 4068.862 | -3714.292 | 3640.342 | -3695.735 | 3621.714 | -3698.991 | 3629.298 | -3790.617 | 3630.045 |
|  | Lim row | 11 | -15316.296 | 15094.682 | -15324.837 | 15111.364 | -15323.642 | 15100.433 | -15319.220 | 15105.14 | -15111.047 | 14984.546 |
|  | Lim row | 12 | -6588.836 | 573.225 | -6561.235 | 5739.295 | -6563.184 | 5741.535 | -6563.380 | 5742.944 | -6565.495 | 5721.129 |
|  | Ltm row | 13 | -18.031 | 180.071 | -16.874 | 172.224 | -16.880 | 172.217 | -17.039 | 172.166 | -17.373 | 173.139 |
|  | Lth row | 14 | -138.362 | 133.735 | -139.851 | 131.348 | -140.204 | 131.659 | -140.250 | 131.769 | -138.204 | 132.125 |
|  | Lth row | 15 | -100.012 | 141.493 | -99.980 | 150,813 | -100.441 | 151.302 | -100.627 | 151.446 | -95.802 | 149.997 |
|  | Ltw row | 16 | -2302.582 | 2691.156 | -2302.934 | 2720.638 | -2306.025 | 2709.896 | -2306.691 | 2710.976 | -2266.597 | 2597.096 |
|  | LTM Row | 17 | -9965.971 | 6880.373 | -10251.104 | 6822.615 | -12267. 150 | 6790.426 | -10277.a88 | 68001.634 | -10180.191 | 6558.322 |
|  | Lth Row | 18 | -12220.455 | 10234.987 | -12191.952 | 10223.477 | -12203.518 | 10212.594 | -12299.673 | 10212.579 | -12232.624 | 10286,099 |
|  | Lim Row | 49 | -910.043 | 1097.881 | -912.467 | 1100.754 | -912.520 | 1101.126 | -912.640 | 1101.521 | -903.279 | 1093.286 |
|  | LTM Row | 50 | -26.328 | 81.262 | -24.331 | 80.765 | -24.317 | ${ }^{79.826}$ | -24.691 | 79.748 | -29.942 | ${ }^{81.354}$ |
|  | LTM Row | 51 | -80,621 | 92.894 | -79.501 | 93.739 | -79.559 | 93.806 | -79.628 | 94.043 | -81.248 | 94.451 |
|  | Ltw row | 52 | -135.642 | 149.003 | -136.605 | 144.737 | $-136.593$ | 145.011 | -137.038 | 145.010 | -139.581 | 149.702 |
| ¢ | LTH Rou | 53 | -201.294 | 397.658 | -190.031 | 403.868 | -190.372 | 406.094 | -190.581 | 406.592 | -194.017. | ${ }^{397.817}$ |
| $\stackrel{\rightharpoonup}{\sim}$ | Lth rou | 54 | -160.506 | 126,025 | -132.522 | 120.977 | -132.493 | 121.681 | -131.217 | 121.498 | -146.011 | 133.610 |
|  | Ltw row | 55 | -123.741 | 77.939 | -136.341 | 67.925 | -137.451 | 68.019 | -137.441 | 68.392 | -147.906 | 71.118 |
|  | LTM Row | 56 | -6.905 | 9.548 | -6.711 | 9.667 | -6.701 | 9.695 | -6.668 | 9.694 | -7.013 | 9.579 |
|  | LIM row | 57 | -959.580 | 1150.629 | -967.793 | 1163.026 | -967.470 | 1162.230 | -967.944 | 1162.744 | -955.373 | 1152.610 |
|  | LTM \%ow | 58 | -129.990 | 107.954 | -125.973 | 110.180 | -127.025 | 109.513 | -127.377 | 109.769 | -130.990 | 107.068 |
|  | LTM \%ow | 59 | -125.672 | 130.057 | -121.905 | 125.366 | -122.312 | 125.712 | -122.174 | 125.77 | -131.460 | 130.171 |
|  | LTM Row | 60 | -10.191 | 13.442 | -9.569 | 13.204 | -9.566 | 13.240 | -9.524 | ${ }^{13.222}$ | -9.052 | 13.088 |
|  | LIM rou | 61 | -82.391 | 106.257 | -78.836 | 104.746 | -78.594 | 104.871 | -78.735 | 104.794 | -70.305 | 103.591 |
|  | LTM Row | 62 | -24.076 | 121.615 | -23.033 | 112.885 | -23.030 | 112.038 | -23.201 | 112.810 | -23.362 | 113.540 |
|  | Lim row | 63 | -72.986 | 80.180 | -74.006 | 79.146 | -74.254 | 79.043 | -74.308 | 79.096 | -73.436 | 77.057 |
|  | Lth row | ${ }_{6}$ | -2733.788 | 1981.400 | -2768,611 | 1955.748 | -2778.468 | 1962.435 | -2779.020 | 1969.346 | -2696.461 | 1998.610 |
|  | Ltm row | 05 | -1986.946 | 2390.924 | -1993.600 | 2403.396 | -1989.133 | 2392.075 | -1988.896 | 2393.278 | -1955.770 | ${ }^{2286.105}$ |
|  | Ltw row | 66 | -997.770 | 1339.519 | -892.090 | 1186.780 | -890.659 | 1199.413 | -896.419 | 1204.604 | -809.017 | 1233.651 |
|  | LTm Row |  | -20.023 | 21.690 | -9.954 | 11.891 | -9.766 | 11.829 | -9.916 | 11.700 | -13.832 | 14.085 |
|  | lith row |  | -83.035 | 81.520 | -61.647 | 52.603 | -62.227 | 57.756 | -61.261 | 59.241 | -81.855 | 71.202 |
|  | LTM Row |  | -16.768 | 13.490 | -9.358 | 6.927 | -9.768 | 7.056 | -9.832 | 6.998 | -12.200 | 10.020 |
|  | Lim row |  | -194.704 | 230.869 | -160.975 | 157.300 | -159.737 | 153.479 | -159.591 | 148.885 | -175.957 | 172.856 |
|  | LTM Row |  | -91.64 | 131.868 | -94.615 | 115.442 | .96.752 | 115.962 | -97.676 | 116.795 | -114.314 | 137.597 |
|  | LTM Row |  | -101.074 | 77.361 | -63.852 | 61.391 | -69.744 | 62.328 | -71.259 | 61.854 | -85.486 | ${ }^{86} .363$ |
|  | LTM Row |  | -37.265 | 49.875 | -27. 195 | 37.518 | -26.767 | 37.609 | -26.729 | 37.616 | -36.476 | 45.999 |
|  | lth row | 156 | -17.823 | 24.791 | -16.363 | 19.031 | -16.361 | 19.043 | -16.262 | 18.948 | -21.151 | 23.945 |

TABLE 3 BTS-26 PAYLOAD/ORBITER INTERFACE LOADS (1bs)


|  |  | TABLE | ST: | 26 I | MOTO | ncce | RATI | NS (9) |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 1 | 2 |  |  |  |  |  |  |  |  |  |
|  |  | Dilagonal | TMP | 1 iter |  | 2 Iter | ons | full Dem | atrix | 1x syatem | mping |
|  | Items | Min | Max | Min | Max | Min | Max | Min | Max | Min | Max |
| c | Small ius Motor $33657 \times$ | 0.101 | 3.091 | 0.119 | 3.123 | 0.118 | 3.124 | 0.119 | 3.125 | 0.060 . | 3.156 |
| \% | Small IUS Motor \#3457 Y | -0.362 | 0.617 | -0.188 | 0.228 | -0.188 | 0.235 | -0.181 | 0.238 | -0.267 | 0.243 |
| $\stackrel{\sim}{*}$ | Small JUS Motor \#3457 2 | -0.450 | 0.629 | -0.670 | 0.610 | -0.669 | 0.610 | -0.470 | 0.611 | -0.473 | 0.621 |
|  | Small ins Motor 33457 Rx | -1.916 | 1.726 | -1.201 | 1.206 | -1.192 | 1.267 | -1.201 | 1.277 | -1.735 | 1.590 |
|  | Small IUS Motor \#3457 Ry | -5.630 | 5.509 | -5.889 | 4.998 | -5.905 | 5.022 | -5.903 | 5.012 | -6.454 | 5.528 |
|  | Small IUS Motor \#3457 Rz | -1.849 | 1.439 | -1.706 | 1.310 | $-1.753$ | 1.420 | -1.756 | 1.644 | -2.410 | 1.950 |
|  | Large lus Motor \#1331 $x$ | 0.423 | 2.735 | 0.418 | 2.754 | 0.416 | 2.754 | 0.416 | 2.754 | 0.418 | 2.727 |
|  | Large IUS Motor \#1331 Y | -0.183 | 0.319 | -0.147 | 0.246 | -0.149 | 0.244 | -0.146 | 0.242 | -0.195 | 0.281 |
|  | Large IUS Motor \#1331 2 | -0.510 | 0.506 | -0.518 | 0.510 | -0.518 | 0.510 | -0.518 | 0.510 | -0.514 | 0.500 |
|  | Large iUs Motor \#1331 Rx | -1.255 | 1.236 | -1.269 | 1.433 | -1.261 | 1.456 | -1.258 | 1.458 | -1.345 | 1.762 |
|  | Lerge IUS Motor \#1331 Ry | -2.246 | 2.395 | -2.332 | 2.421 | -2.321 | 2.416 | -2.319 | 2.427 | -2.596 | 2.714 |
|  | Large IUS Hotor \#1331 Rz | -4.069 | 4.315 | -2.261 | 2.384 | -2.270 | 2.423 | -2.179 | 2.313 | -4.045 | 4.218 |

TABLE 5 STS-26 PAYLOAD/OREITER INTERFACE ACCELERATIONS (g)



Figure 1 MODAL CONTRIBUTIONS FOR I-ACCELERATION (in/sec ${ }^{2}$ ) OF STS-26 SA ANTENNA RIBS USING TMP METHOD, $t=8.3$ sec.


Figure 2 MODAL CONTRIBUTIONS FOR X-ACCELERATION (in/sec ${ }^{2}$ ) OF STS-26 SA ANTEMNA RIBS USING ITERATIVE METHOD $t=8.3 \mathrm{sec}$.

