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PART VII

**UNSTEADY AERODYNAMICS FOR  
ADVANCED CONFIGURATIONS**

**PART VII — VELOCITY POTENTIALS IN NON-UNIFORM  
TRANSONIC FLOW OVER A THIN WING**

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## FOREWORD

This report covers a portion of the research conducted by the Los Angeles Division of North American Rockwell Corporation, Los Angeles, California, for the Aerospace Dynamics Branch, Vehicle Dynamics Division, Air Force Flight Dynamics Laboratory, Wright-Patterson Air Force Base, Ohio, under Contract No. AF33(615)-2896.

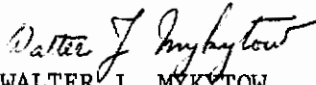
The work was performed to advance the state-of-the-art of flutter prediction for flight vehicles as part of the Air Force Systems Command exploratory development program. The research was conducted under Project No. 1370 "Dynamic Problems in Flight Vehicles", Task No. 137003 "Prediction and Prevention of Aero-thermoelastic Problems". Messrs. James J. Olsen and Samuel J. Pollock of the Aerospace Dynamics Branch were Project Engineers.

Mr. H. Hoge was the Program Manager for North American Rockwell. Mr. L. V. Andrew and Mr. T. E. Stenton were Principal Investigators. The basic approach was outlined by Dr. M. T. Landahl of the Massachusetts Institute of Technology. The calculus of variations approach was suggested by Mr. James Olsen.

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## ABSTRACT

Two methods have been outlined in detail, and one of them has been mechanized, for calculating acoustic ray paths emanating from any point in a non-uniform transonic flow field surrounding a wing. It gives the ray path, and the time, for the minimum time of travel from the acoustic source point to the field point. The resulting velocity potential is also computed.

It was necessary to establish an accurate representation of the flow characteristics in the field surrounding the wing. Some ray lines travel over the planform and into the surrounding flow field. It was established that once off the planform they do not return.

Available methods predict phase lags based on the assumption that acoustic rays travel in straight lines. The results of this study show this to be a very poor approximation at transonic speeds. Therefore, it is recommended that the method presented in this report be fully developed for the purpose of calculating generalized forces on wings in harmonic motion at transonic speeds. A computer program that would predict these phase lags with reasonable accuracy, and the corresponding flutter characteristics and unsteady aerodynamic loads on a wing responding to externally applied forces, such as gusts, would fill an important gap in the available technology.

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## SYMBOLS

|                                |   |
|--------------------------------|---|
| c                              | chord   |
| C                              | speed of sound  |
| g                              | time of travel of an acoustic signal  |
| M                              | Mach number   |
| r                              | Slope in the y-direction, dx/dy   |
| $\delta R$                     | Increment in radius vector  |
| s                              | Distance along a ray path, span   |
| t                              | Time  |
| U                              | Free-stream velocity  |
| V                              | Velocity  |
| x, y, z                        | Location of a field point   |
| $x_0, y_0, z_0$                | Location of a source or doublet point   |
| X, Y                           | $x/\beta s, y/s$  |
| X*, Y*                         | Linear transformation of coordinates X, Y   |
| $\hat{i}', \hat{j}', \hat{k}'$ | Unit vectors along x', y', z' axes  |
| $R$                            | Radius vector   |
| $\nabla$                       | Vector gradient operator, $\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z}$ |
| $\beta$                        | $\sqrt{1-M^2}, \sqrt{r^2 + 1-M^2}$  |
| $\delta$                       | $\sqrt{5 + M^2},$ an increment  |
| $\phi$                         | Velocity potential  |
| $\Lambda$                      | Ray Angle   |
| $\tau$                         | Thickness ratio   |
| Subscripts                     |   |
| a                              | Advancing   |
| l                              | Local, lower  |

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## SYMBOLS (Continued)

|              |   |
|--------------|---|
| r            | Receding  |
| u            | Upper   |
| x, y         | Partial derivatives with respect to x, y            |
| $\sigma$     | Sonic line  |
| $\infty$     | Infinity  |
| Superscripts |   |
| $\circ$      | Derivative with respect to time                     |
| '            | Derivative with respect to the independent variable |

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## INTRODUCTION

When an airfoil travels through the air at speeds near the speed of sound, the local speed of flow varies from subsonic near the forward edges to supersonic near the trailing edges. These wide variations of speed from that of the free-stream characterize the non-uniform transonic flow. This non-uniformity of the flow field must be accounted for in accurate calculations of unsteady pressures and forces; particularly their phase lags.

In order to determine an unsteady transonic flow field one requires solutions for singularities immersed in a non-uniform steady flow, (Reference 1). Source solutions for a mean flow that varied in the x-direction only were given in the high-frequency limit by Landahl (Reference 2). Rodemich (Reference 3) presented a "box" solution, based on pulsating doublets, which assumes a uniform mean flow at Mach number 1.0. No exact solutions for the case of a mean flow with arbitrary spatial variations have been found, thus far, but Landahl proposed the basic form of a solution which removes most of the limitations and restrictions of these approximate solutions. The method focuses attention on the time of transmission of an acoustic signal from a pulsating sending source to a distant receiving point. The signal travels through a nearly sonic flow field where the Mach number varies in a prescribed manner.

This report contains a difference equation approach, and differential equation approach to computing the paths and the transmission times for acoustic signals. The independent variable in the latter approach is a spatial rather than a time variable. A procedure that could be used to calculate the velocity potentials and generalized forces on an oscillating surface is described.

POTENTIAL OF A UNIT SOURCE

The basic expressions proposed by Landahl for the velocity potential at the point  $(x,y,z)$  due to a pulsating source at  $(x_0,y_0,z_0)$  are:

(a) for a source in a locally subsonic flow region

$$\phi = \frac{-1}{4\pi\bar{R}} \exp \{i\omega[t-g(x,y,z,x_0,y_0,z_0)]\} \quad (1)$$

where

$$\bar{R} = \sqrt{(x-x_0)^2 + [1-M^2(x,y,z)][(y-y_0)^2 + (z-z_0)^2]}$$

$M$  = Local Mach Number

$x_0,y_0,z_0$  = Location of source point

$g(x,y,z,x_0,y_0,z_0)$  = Time required for a disturbance to travel from  $(x_0,y_0,z_0)$  to  $(x,y,z)$ .

(b) for a source in a locally supersonic flow region

$$\phi = \frac{-1}{4\pi\bar{R}} \{ \exp[i\omega(t-g_a)] + \exp[i\omega(t-g_r)] \} \quad (2)$$

where

$g_{a,r} = g_{a,r}(x,y,z,x_0,y_0,z_0)$  = Time required for the advancing, receding wave to travel from  $(x_0,y_0,z_0)$  to  $(x,y,z)$

It is likely that good accuracy may be obtained with use of the value of  $g_a$  for uniform flow (in the supersonic case, and also for the advancing wave portion in the subsonic case). However, our purpose is to produce a general solution for  $g$  which applies to both the advancing and the receding portions of the wave and compare values with those for uniform flow.

Since the primary interest is in wing flows, we consider that both the source and receiver points lie in the  $x, y$ -plane, so that  $z = z_0 = 0$ . Furthermore, we consider that signals do not return to the plane once they leave. The problem is thus simplified to one in two spatial dimensions. Its solution should be applicable to a wide variety of nearly planar lifting surfaces.

Consider a signal emanating from a source at the point  $(x_0,y_0)$  on a wing. A second point past which the signal travels is located an incremental distance  $(dx, dy)$  away. There are two components of velocity of the signal, a radial component,  $C$ , where  $C$  is the local speed of sound and an  $x$ -component,  $U$ , where  $U$  is the local speed of flow over the wing.  $\Lambda$  is the angle the radial component makes with the negative extension of the  $x$ -axis. The path of this wavefront point will be referred to as a "ray". The shape of any ray depends on the initial choice of  $\Lambda$ ; for a given  $\Lambda$ ,  $dx$  and  $dy$  are components of the first element of this particular ray emanating from  $(x_0,y_0)$ . The situation depicted is general in that it applies not only at the source, but at any point on the ray path. Thus, the

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velocity at any point on the path is a function of three spatial parameters which vary with position,  $U$ ,  $C$ , and  $\Lambda$ . From the sketch, it is clear that

$$dx = [U(x, y) - C(x, y) \cos \Lambda] dt \quad (3)$$

$$dy = C(x, y) dt \sin \Lambda$$

Equations were developed for two methods of tracing the ray path to establish the magnitude and the phase relationship at field points to a unit source. These methods are: (1) a difference equation method, and (2) a non-linear differential equation method.

## Difference Equation Method

In this method, time is the independent variable. Equations (3) are two of the three equations needed to establish the variation of  $x$ ,  $y$ , and  $\Lambda$  with time. The third equation is obtained by considering the acceleration of the ray in the non-uniform flow field (see Figure 1).

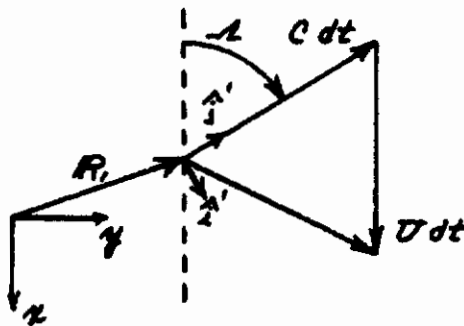


Figure 1. Velocity Components of a Sonic Ray Line In A Moving Airstream

In terms of components in the directions of the rotating unit vectors  $\hat{i}'$  and  $\hat{j}'$

$$\dot{R}_1 = (U \sin \Lambda) \hat{i}' + (C - U \cos \Lambda) \hat{j}' \quad (4)$$

and

$$\ddot{R}_1 = (\dot{U} \sin \Lambda + C \dot{\Lambda}) \hat{i}' + (\dot{C} - \dot{U} \cos \Lambda) \hat{j}'$$

It is necessary to express the angular velocity  $\dot{\Lambda}$  in terms of space variables. To do this, consider that at time  $t$  a second ray point is located at  $R_2 = R_1 + \delta R \hat{i}'$ , where  $\delta R$  is small, and it's direction of travel is  $\dot{R}_2 = \dot{R}_1 + \delta \dot{R}$ . Let the superscripts (0) and (1) denote times  $t_0$  and  $t_1 (= t_0 + \Delta t)$ . Then at time  $t_1$

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$$R_1^{(1)} = R_1^{(0)} + \dot{R}_1^{(0)} \Delta t$$

and 
$$R_2^{(1)} = R_2^{(0)} + \dot{R}_2^{(0)} \Delta t$$

Subtracting the first equation from the second

$$\delta R^{(1)} = \delta R^{(0)} + \delta \dot{R}^{(0)} \Delta t \quad (5)$$

where 
$$\delta R = R_2 - R_1$$

Recalling that the cross product of two vectors is a vector normal to the plane defined by the two vectors, and has a magnitude equal to the product of the two magnitudes times the sine of the angle between them, then

$$\delta R^{(0)} \times \delta R^{(1)} = \hat{k}' (-\delta R^{(0)} \delta R^{(1)} \sin \Delta \mathcal{A}) \quad (6)$$

which has the correct sense. When  $\Delta \mathcal{A}$  is small, and when Equation (5) is substituted into the left side of Equation (6), we get

$$\delta R^{(0)} \times \delta \dot{R}^{(0)} \Delta t = \hat{k}' (-\delta R^{(0)} \delta R^{(1)} \Delta \mathcal{A})$$

This may be rewritten as

$$\frac{\Delta \mathcal{A}}{\Delta t} = - \frac{\delta (C - U \cos \mathcal{A})}{\delta R^{(1)}}$$

and in the limit as  $\Delta t \rightarrow 0$

$$\dot{\mathcal{A}} = - \hat{k}' \cdot \nabla (C - U \cos \mathcal{A}) \quad (7)$$

where the operator  $\hat{k}' \cdot \nabla$  is

$$\hat{k}' \cdot \nabla = \left( \sin \mathcal{A} \frac{\partial}{\partial x} + \cos \mathcal{A} \frac{\partial}{\partial y} \right)$$

and operates only on C and U.

Equation (7) has a revealing physical interpretation. From Figure 4 we see that the gradient of the speed of sound C, on forward portions of the wing, is a vector pointing forward and slightly outward from the center-line; whereas, from Figure 3 we see that the gradient of the local flow speed U is nearly in the opposite direction. Although it is not apparent from the figures because they are plotted to different scales, the magnitude of the gradient of U is about five times that of the gradient of C. From the energy equation  $C^2 + \frac{\gamma-1}{2} U^2 = \text{constant}$ ,  $\nabla U = -5.0 \nabla C$ . The local Mach number is increasing in the downstream direction. Figure 2 shows that, under these conditions there are only two stable ray angles; those for which the gradient of  $C - U \cos \mathcal{A}$  is zero. As the ray propagates through the flow field it will always tend towards one of these two orientations.

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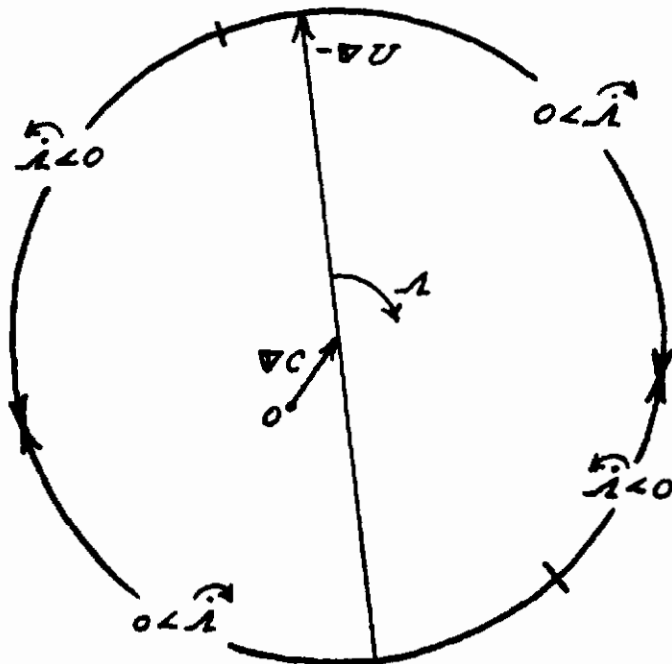


Figure 2. Stability of Ray Angles When The Gradient of Local Flow Speed Exceeds the Gradient of Local Speed of Sound.

We now write Equations (3) and (7) in difference form

$$\Delta x = [U - C \cos \Lambda] \Delta g \quad (8-a)$$

$$\Delta y = [C \sin \Lambda] \Delta g \quad (8-b)$$

and

$$\Delta \Lambda = - \left[ \sin \Lambda \left( \frac{\partial C}{\partial x} - \cos \Lambda \frac{\partial U}{\partial x} \right) + \cos \Lambda \left( \frac{\partial C}{\partial y} - \cos \Lambda \frac{\partial U}{\partial y} \right) \right] \Delta g \quad (8-c)$$

where  $\Delta g$  represents an increment in disturbance travel time  $g$ , defined previously. To determine  $\varphi(x, y, 0, x_0, y_0, 0)$  it is necessary to know a steady state distribution of  $C(x, y)$ ,  $U(x, y)$ , and their derivatives at any point in the flow field over the wing and in the surrounding flow field in the plane of the wing. A means for establishing these is given in Section 5. Assume they are known. Then the procedure used is as follows:

1. Select any source point, on or off the wing,  $(x_0, y_0)$ .

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2. Select a series of initial ray angles,  $\Lambda_1$ ,  $i = 1, 2, \dots$ .
3. Select an initial increment in disturbance travel time,  $\Delta g_0$ .
4. For each of the ray angles store  $x^{(1)}$ ,  $y^{(1)}$ ,  $\sin \Lambda^{(1)}$ ,  $\cos \Lambda^{(1)}$ , and  $\Delta g^{(1)}$ ,  $i = 1, 2, \dots$ .
  - a. At  $x^{(1)}$ ,  $y^{(1)}$  compute and store  $x^{(1)} = x^{(1)} + \Delta x^{(1)}/2$  and  $y^{(1)} = y^{(1)} + \Delta y^{(1)}/2$ , holding  $\Lambda$  constant.
  - b. Iterate on  $x_2^{(1)} = x^{(1)} + \Delta x^{(1)}/2$ ,  $y_2^{(1)} = y^{(1)} + \Delta y^{(1)}/2$ , and  $\Delta \Lambda(x_2^{(1)}, y_2^{(1)})$  until they converge or exceed ten trials. In the latter case replace  $\Delta g^{(1)}$  by  $\Delta g^{(1)}/2$  and repeat the iteration. If they converge in three trials or less, replace  $\Delta g^{(1)}$  by  $2\Delta g^{(1)}$ .
  - c. Replace  $x^{(1)}$  by  $x_2^{(1)}$ ,  $y^{(1)}$  by  $y_2^{(1)}$ , and return to a.

The solutions presented above are believed to be good approximations to the exact solutions for the following reasons:

1. For the case of a uniform flow they reduce to the proper linearized expressions.
2. The phase of the disturbance will be exact, although the amplitude may be slightly in error.
3. In an inner region in the immediate neighborhood of the source location  $(x_0, y_0, z_0)$  they approach the correct solution.
4. For a one-dimensional mean flow with  $M_\infty$  approaching unity they reduce to Landahl's earlier solution (Reference 2).
5. In the limit of steady flow ( $\omega = 0$ ), the solutions give results equivalent to the local linearization method of Spreiter and Alksne (Reference 4). This has been demonstrated by Rubbert (Reference 5).
6. Inasmuch as the proposed approximation only affects the receding part of the solution, the proper limiting solution for high frequencies (Reference 1), should always be obtained since then receding-wave effects are largely cancelled out due to the rapid phase variations.

This method gives reasonable results, i.e., reasonable based on a comparison with results obtained from the differential equation method. However, the ray paths did not conclusively show the existence of the focal point that the second method revealed.

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## Non-Linear Differential Equation Method

From Equations (3) we may write the slope of the ray path

$$\frac{dx}{dy} = \frac{M - \cos \Lambda}{\sin \Lambda} \quad (9)$$

and solving this equation for  $\cos \Lambda$ , we get

$$\cos \Lambda = \frac{M \pm r \sqrt{r^2 + 1 - M^2}}{1 + r^2} \quad (10)$$

where  $r = \frac{dx}{dy}$

The transmission time from source to receiving point is given by

$$g = \int \frac{ds}{V} \quad (11)$$

where the integration is taken along the path and

$$ds = \sqrt{1 + r^2} \quad dy \quad (12)$$

The velocity along the path is obtained from the vector sum of the two velocity components

$$V = C \sqrt{M^2 + 1 - 2M \cos \Lambda} \quad (13)$$

Substituting equations (12), (13), and (10) into equation (11) we have:

$$g = \int \frac{\sqrt{1+r^2} dy}{C \sqrt{M^2 + 1 - 2M \left[ \frac{M \pm r \sqrt{r^2 + 1 - M^2}}{1 + r^2} \right]}}$$

which reduces to

$$g = \int \frac{(1+r^2) dy}{C \sqrt{M^2 r^2 \mp 2Mr \sqrt{r^2 + 1 - M^2} + r^2 + 1 - M^2}} \quad (14)$$

The radicand in the denominator is a perfect square. Thus,

$$g = \int \frac{(1+r^2) dy}{C [Mr \mp \sqrt{r^2 + 1 - M^2}]}$$

which reduces to

$$g = \int \frac{Mr \pm \sqrt{r^2 + 1 - M^2}}{C (M^2 - 1)} dy \quad (15)$$

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At this point we relate the local acoustic velocity,  $C = C(x, y)$ , to the local Mach number by imposing the condition of conservation of energy. For non-viscous flow, the total temperature is conserved. It is easily verified, that under this condition

$$\frac{C^2}{C_\infty^2} = \frac{5+M^2}{5+M_\infty^2} \quad (16)$$

where  $\gamma = 1.4$ , for a diatomic gas, has been used. Substituting Equation (16) into Equation (15), we get

$$g = \frac{1}{C_\infty \sqrt{5+M_\infty^2}} \int \frac{\sqrt{5+M^2} [Mr \pm \sqrt{r^2+1-M^2}]}{(M^2-1)} dy \quad (17)$$

where the upper sign applies to receding waves and the lower sign to advancing waves. Equation (17) contains all the elements for the solution. However, the integrand is a function of  $x$ ,  $y$ , and  $dx/dy$ . This equation may be written in symbolic form

$$g = \int_{y_0}^{y_1} F(x, y, \frac{dx}{dy}) dy$$

which suggests the use of Euler's equation to find the minimum time  $g$ , for the disturbance to travel to a field point  $(x_1, y_1)$

$$\frac{d}{dy} \frac{\partial F}{\partial r} - \frac{\partial F}{\partial x} = 0 \quad (18)$$

In order to simplify the notation, we set

$$F = \frac{\delta (Mr \pm \beta)}{M^2-1}$$

where  $\delta = \delta(x, y) = \sqrt{5+M^2}$

$$\beta = \beta(x, y, r) = \sqrt{r^2+1-M^2}$$

and  $r$  has been previously defined. We will need

$$\frac{\partial F}{\partial x} = \frac{\delta}{M^2-1} \left[ r M_x \mp \frac{M M_x}{\beta} \right] + \frac{M r \pm \beta}{(M^2-1)^2} \left[ (M^2-1) \frac{M M_x}{\delta} - 2 \delta M M_x \right]$$

$$\begin{aligned} \frac{d}{dy} \left( \frac{\partial F}{\partial r} \right) &= \frac{\delta}{M^2-1} \left[ \frac{dM}{dy} \pm \beta \frac{dr}{dy} - r \frac{d\beta}{dy} \right] \\ &+ \left( M \pm \frac{r}{\beta} \right) \left[ \frac{(M^2-1) \frac{d\delta}{dy} - 2 \delta M \frac{dM}{dy}}{(M^2-1)^2} \right] \end{aligned}$$



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Then, making use of the relationships

$$\frac{dM}{dy} = rM_x + M_y$$

$$\frac{d\beta}{dy} = \frac{1}{\beta} \left[ r \frac{dr}{dy} - rMM_x - MM_y \right]$$

$$\frac{d\delta}{dy} = \frac{1}{\delta} \left[ rMM_x + MM_y \right],$$

solving for  $dr/dy$ , and combining terms, we get

$$\frac{dr}{dy} = \frac{1}{\delta^2(M^2-1)} \left\{ \left[ \frac{-M(M^2+1)r}{M^2-1} \mp \frac{\beta(7M^2+5)}{M^2-1} \right] r^2 + [2M(M^2+8)r \pm \beta(7M^2+5)] M_y + \left\{ \frac{M}{\delta^2}(r^2+6) \right\} M_x \right\} \quad (19)$$

Equation (19) is a second order, second degree differential equation of the form

$$\frac{d^2x}{dy^2} = f(x, y, \frac{dx}{dy})$$

It is second degree because  $\beta$  represents a radical. However, it can be solved numerically by any of the standard repetitive processes. We employed a fourth order Runge-Kutta procedure.

There are certain difficulties that arise in the numerical evaluation of Equation (19). These are first listed and interpreted and then equations used to surmount them are presented.

- (1) Along some ray paths  $dx/dy$  becomes infinite even when the Mach number is not equal to one.
- (2) Equation (19) is singular at Mach number = 1.0.
- (3) In the supersonic region, signals sometimes become trapped on the local Mach line. This happens when  $\cos \Lambda = 1/M$ . Signals tend to gravitate to this condition. Such trapped signals cannot then cross the sonic line. They approach the sonic line as a limit, and are cancelled out there.

To overcome the difficulty listed in Item (1), it is necessary to use  $x$  instead of  $y$  as the independent variable. This is done by applying the equation

$$\frac{d^2y}{dx^2} = \frac{-1}{\left(\frac{dx}{dy}\right)^3} \frac{d^2x}{dy^2} \quad (20)$$

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It is convenient here to introduce some new notation. Re-write equation (19) in the form

$$\begin{aligned} \kappa'' = \frac{1}{AB} \left\{ -\frac{M}{B} (M^2+11) \kappa'^3 + 2M(M^2+8) \kappa' \mp (7M^2+5) \frac{R_1^3}{B} \right\} M_y \\ + \frac{M}{A} \left\{ \kappa'^2 + 6 \right\} M_x \end{aligned} \quad (21)$$

where the new notation, together with some other notation which will be used later, is defined as follows

$$\begin{aligned} \kappa' &= \frac{dx}{dy} & A &= M^2+5 & M_x &= \frac{\partial M}{\partial x} \\ y' &= \frac{dy}{dx} & B &= M^2-1 & M_y &= \frac{\partial M}{\partial y} \\ R_1 &= \sqrt{\kappa'^2 - (M^2-1)} & \beta &= \sqrt{B} \\ R_2 &= \sqrt{1 - y'^2 (M^2-1)} & E &= C_D \sqrt{5+M^2} \end{aligned} \quad (22)$$

Substituting Equation (20) into Equation (21), we get

$$\begin{aligned} y'' = \frac{1}{AB} \left\{ \frac{M}{B} (M^2+11) - 2M(M^2+8) y'^2 \pm \frac{R_2^3}{B} (7M^2+5) \right\} M_y \\ - \frac{M y'}{A} (6y'^2+1) M_x \quad ; \quad 0 \leq y' \end{aligned} \quad (23-a)$$

$$\begin{aligned} y'' = \frac{1}{AB} \left\{ \frac{M}{B} (M^2+11) - 2M(M^2+8) y'^2 \mp \frac{R_2^3}{B} (7M^2+5) \right\} M_y \\ - \frac{M y'}{A} (6y'^2+1) M_x \quad ; \quad y' < 0 \end{aligned} \quad (23-b)$$

The limiting form of Equation (20) at  $M = 1$  is:

$$\kappa'' \Big|_{M=1.0} = \frac{1}{2A} \left\{ 2\kappa'^3 + \kappa' + \frac{9}{\kappa'} \right\} M_y + \frac{1}{A} (\kappa'^2 + 6) M_x$$

In the supersonic region, when the signal is trapped on the local Mach line, and

$$\cos \alpha = \frac{1}{M}, \quad \sin \alpha = \frac{\sqrt{1-M^2}}{M}, \quad \text{and} \quad |\kappa'| = \beta$$

equation (20) reduces to

$$\kappa'' = M \left( \frac{M_y}{\kappa'} + M_x \right)$$

# Contrails

A complete set of equations, together with their areas of applicability, will now be outlined.

Complete Set of Equations where Y is the Independent Variable

$$\kappa'' = \frac{1}{AB} \left\{ \frac{-M}{B} (M^2 + 11) \kappa'^3 + 2M(M^2 + 8) \kappa' \mp (7M^2 + 5) \frac{R_1^3}{B} \right\} M_y + \frac{M}{A} (\kappa'^2 + 6) M_x \quad (24)$$

$$\frac{dt}{dy} = \frac{1}{E} \frac{\sqrt{5+M^2} (M\kappa' \pm R_1)}{M^2 - 1} \quad (25)$$

$$\kappa'' \Big|_{M=1.0} = \frac{1}{2A} \left\{ 2\kappa'^3 + \kappa' + \frac{9}{\kappa'} \right\} M_y + \frac{M}{A} (\kappa'^2 + 6) \quad (26)$$

$$\frac{dt}{dy} \Big|_{M=1.0} = \frac{\sqrt{6}}{2E} \left( \kappa' + \frac{1}{\kappa'} \right) \quad (27)$$

$$\kappa'' \Big|_{|\kappa'|=\beta} = M \left( \frac{M_y}{\kappa'} + M_x \right) \quad (28)$$

$$\frac{dt}{dy} \Big|_{|\kappa'|=\beta} = \frac{M\sqrt{5+M^2}}{E\kappa'} \quad (29)$$

A complete set of equations were also developed using x as the independent variable. However, for the sake of brevity, and since they are obtained by a simple change of variable, they will not be listed here. Equations (26) and (27) apply where an advancing ray path crosses the sonic line, and equations (28), (29) apply where a ray path, in the supersonic region, becomes trapped on the local Mach line. It remains to describe the regions of applicability of the upper and lower signs of equations (24) and (25). In what follows, "right branch" will be specified where  $(0 < \mathcal{A} < \pi)$  and left branch will be specified if  $(-\pi < \mathcal{A} < 0)$ . Here  $\mathcal{A}$  is the local value along the ray path. The end points are not specified because for these points we use x as the independent variable.

The upper sign is used for

- (1) Subsonic, left branch
- (2) Supersonic, receding, right branch
- (3) Supersonic, advancing, left branch

# *Contrails*

The lower sign is used for

- (1) Subsonic, right branch
- (2) Supersonic, receding, left branch
- (3) Supersonic, advancing, right branch

## THE NON-UNIFORM FLOW FIELD

In the application of each of the methods contained in this report, it is necessary to know certain of the properties of the transonic flow field on, and in the neighborhood of, the wing. Figures 3 and 4 show the distributions of local flow speeds and sonic speeds over a  $65^\circ$  delta wing model in a wind tunnel in which the Mach number was 1.04 (taken from Reference 6). Speeds were computed from steady state pressure data at 27 points on the wing. The figures are intended only to show the general characteristics of the flow, such as: (1) The local sonic line shifts aft with distance from the centerline but crosses the leading edge inboard of the tip, (2) Mach number variations in both the streamwise and spanwise directions must be considered and cannot be considered to be linear, and (3) Separated flow is indicated over the aft and inboard portion of the wing. To consider the last of these characteristics is beyond the scope of this study. However, the first two are amenable to analysis using available theories and techniques.

$$M_\infty = 1.04$$

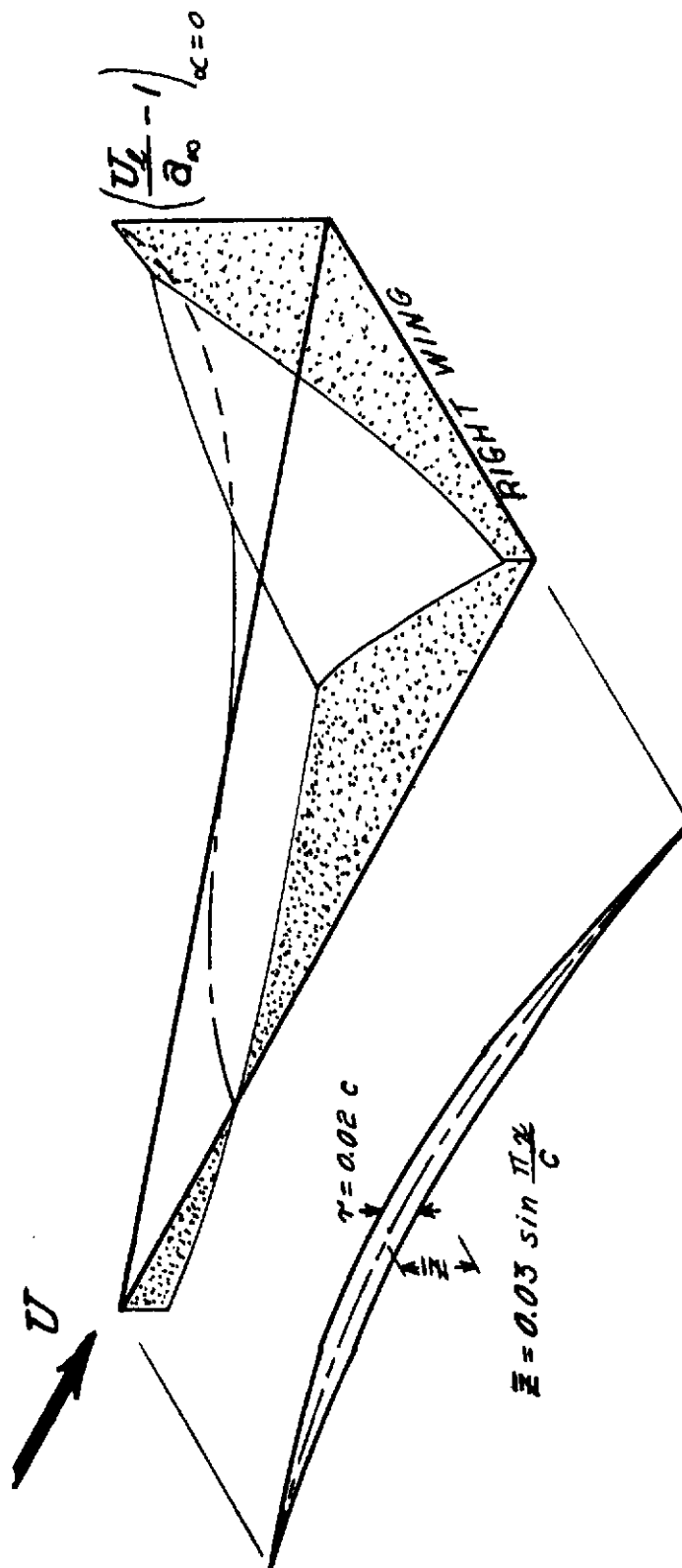


Figure 3. Local Flow Distribution on a  $65^\circ \Delta$  at a Transonic Speed

$$M_\infty = 1.04$$

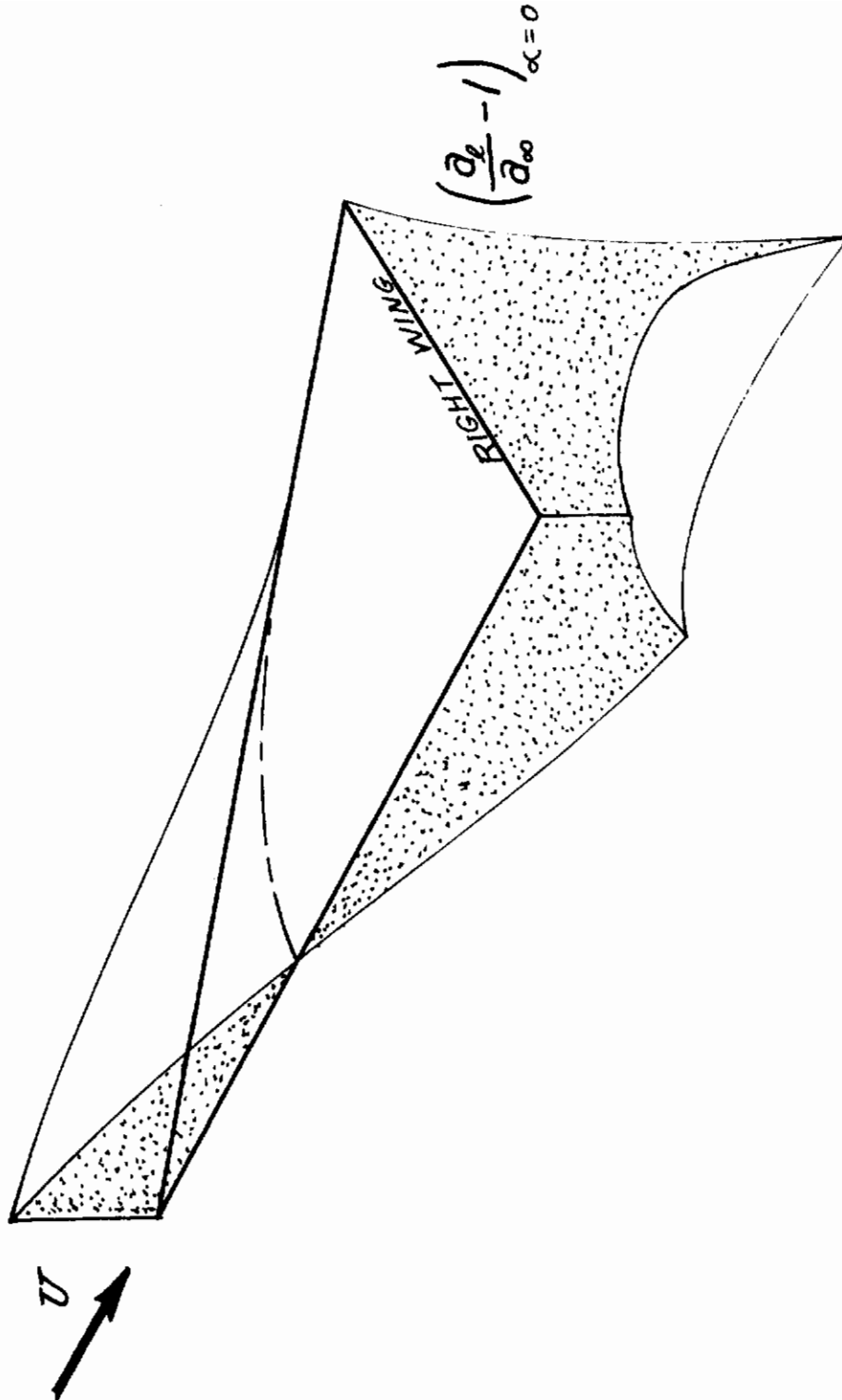


Figure 4. Sonic Speed Distribution on a  $65^\circ \Delta$  at a Transonic Speed

# Contrails

Mach number distributions over areas off the wing were computed from an approximate theoretical solution of the flow field that matched pressure distributions on the wing. In order to avoid a discontinuity at the juncture of the two regions, a small transition region was defined over which the two functions were joined by a numerical smoothing technique.

Let:

$$M_L = M_L(x,y) = \text{Mach number}$$

$$\Phi = \Phi(x,y) = \text{Perturbation potential}$$

$$\tilde{\tau} = \tau/c(x,y) = \text{Thickness ratio}$$

For a steady-state, non-lifting flow

$$(1 - M_L^2) \Phi_{xx} + \Phi_{yy} + \Phi_{zz} = 0 \quad (30)$$

and  $\Phi_z(x,y,0^+) = \pm \tilde{\tau} f_x(x,y) \quad (31)$

Where  $f(x,y)$  is a function describing the variation of the surface from the mean.

Using parametric differentiation with respect to  $\tilde{\tau}$ , (Reference 5),

$$g = g(x,y) = \frac{\partial \Phi}{\partial \tilde{\tau}}$$

Equation (30) becomes:

$$\frac{\partial}{\partial x} \{ [1 - M_L^2] g_x \} + g_{yy} + g_{zz} = 0 \quad (32)$$

$$g_z(x,y,0^+) = \pm f_x(x,y)$$

After having obtained the solution of equation (32), the local Mach number distribution is obtained by relating local Mach number to the coefficient of pressure, ( $C_p$ ). Starting with the following basic relations:

Let  $u = \frac{V_L - U_\infty}{U_\infty}$

then  $u = \frac{1}{U_\infty} \frac{\partial \Phi}{\partial x} = -C_p/2 \quad (33)$

$$a^2 + \frac{1}{2}(\gamma - 1) \beta^2 = \text{constant} \quad (34)$$

where  $q = U_\infty$  at infinity

$q = U_\infty(1 + u)$  elsewhere

$a = \text{speed of sound}$



# Contrails

We have: 
$$a_{\infty}^2 + \frac{1}{2}(\gamma-1)U_{\infty}^2 = a_L^2 + \frac{1}{2}(\gamma-1)U_{\infty}^2(1+\alpha)^2$$

$$U_{\infty}^2(1+\alpha)^2 \cong U_{\infty}^2(1+2\alpha)$$

$$a_L^2 \cong a_{\infty}^2 - (\gamma-1)U_{\infty}^2\alpha$$

using equation (33)

$$a_L^2 \cong a_{\infty}^2 \left[ 1 + \frac{1}{2}(\gamma-1)M_{\infty}^2 C_p \right]$$

The coefficient of pressure,  $C_p$  is of order (.1), and  $M$  is  $O(1.)$ . Therefore, to sufficient accuracy.

$$a_L \cong a_{\infty} \left[ 1 + \frac{1}{4}(\gamma-1)M_{\infty}^2 C_p \right]$$

$$U_L = U_{\infty}(1+\alpha) = U_{\infty} \left( 1 - \frac{1}{2}C_p \right)$$

and from these relations:

$$M_L \cong \frac{M_{\infty} \left( 1 - \frac{1}{2}C_p \right)}{1 + \frac{1}{4}(\gamma-1)M_{\infty}^2 C_p}$$

Noting again the order of  $M_{\infty}$  and  $C_p$ , to sufficient accuracy.

$$M_L \cong M_{\infty} \left( 1 - \frac{1}{2}C_p \right) \left[ 1 - \frac{1}{4}(\gamma-1)M_{\infty}^2 C_p \right]$$

or 
$$M_L \cong M_{\infty} \left[ 1 - \frac{\gamma+1}{4} C_p \right] \tag{35}$$

Equation (35) is the expression that was used to relate local Mach number to  $C_p$  on regions off the wing.

A solution of equation (30), using the results of equation (35) was worked out for a special configuration. The special wing configuration is depicted in figure (5).

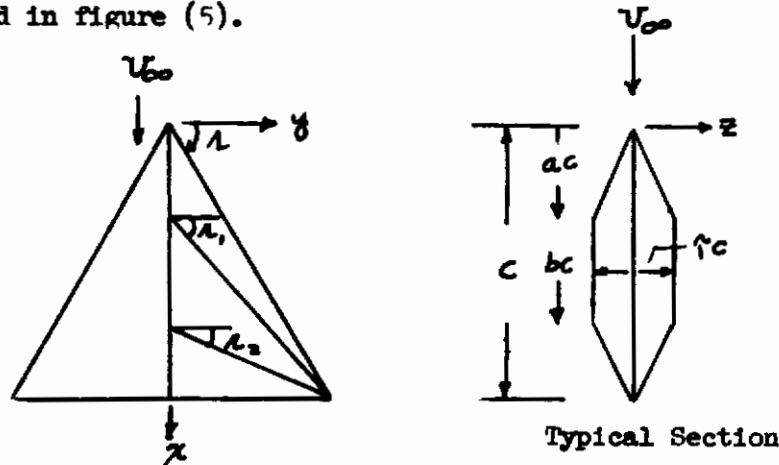


Fig. 5. A Thin Wing In Rectilinear Flight

The solution is:

$$C_p(x, y) - C_p(x, 0) = -2f_1 \left[ |y-s_1|^{\epsilon_1} + |y+s_1|^{\epsilon_1} - 2|s_1|^{\epsilon_1} \right] \\ - 2f_2 \left[ |y-s_2|^{\epsilon_2} + |y+s_2|^{\epsilon_2} - 2|s_2|^{\epsilon_2} \right] H(x-a) \\ - 2f_2 \left[ |y-s_2|^{\epsilon_2} + |y+s_2|^{\epsilon_2} - 2|s_2|^{\epsilon_2} \right] H(x-b) \tag{36}$$

where  $H(x)$  is a step function.

# Contrails

$$\begin{aligned} \frac{\partial \Delta C_p}{\partial x} = & -\frac{2f\epsilon}{\tan \Lambda} \left[ -|y-s|^{\epsilon-1} + |y+s|^{\epsilon-1} - 2|s|^{\epsilon-1} \right. \\ & + \frac{2f_1\epsilon_1}{\tan \Lambda_1} \left[ -|y-s_1|^{-\epsilon_1-1} + |y+s_1|^{-\epsilon_1-1} - 2|s_1|^{-\epsilon_1-1} \right] H(x-a) \\ & + \frac{2f_2\epsilon_2}{\tan \Lambda_2} \left[ -|y-s_2|^{-\epsilon_2-1} + |y+s_2|^{-\epsilon_2-1} - 2|s_2|^{-\epsilon_2-1} \right] H(x-b) \end{aligned} \quad (37)$$

$$\begin{aligned} \frac{\partial \Delta C_p}{\partial y} = & -2f\epsilon \left[ |y-s|^{\epsilon-1} + |y+s|^{\epsilon-1} \right] \\ & + 2f_1\epsilon_1 \left[ |y-s_1|^{-\epsilon_1-1} + |y+s_1|^{-\epsilon_1-1} \right] H(x-a) \\ & + 2f_2\epsilon_2 \left[ |y-s_2|^{-\epsilon_2-1} + |y+s_2|^{-\epsilon_2-1} \right] H(x-b) \end{aligned} \quad (38)$$

WHERE:

$$f = \frac{\cos^2 \Lambda}{\sin \Lambda}, \quad \epsilon = \frac{\tau}{2\pi a \cos \Lambda}, \quad s = \tau / \tan \Lambda$$

$$f_1 = \frac{\cos^2 \Lambda_1}{\sin \Lambda_1}, \quad \epsilon_1 = \frac{\tau}{2\pi a \cos \Lambda_1}, \quad s_1 = (x-a)/(1-a) \tan \Lambda_1$$

$$f_2 = \frac{\cos^2 \Lambda_2}{\sin \Lambda_2}, \quad \epsilon_2 = \frac{\tau}{2\pi a \cos \Lambda_2}, \quad s_2 = (x-b)/(1-b) \tan \Lambda_2$$

$$\tan \Lambda_1 = (1-a) \tan \Lambda, \quad \tan \Lambda_2 = (1-b) \tan \Lambda$$

After determining a distribution of  $C_p$  and its derivatives from equations (36), (37), and (38), the Mach number distribution, with its derivatives, is computed from equation (35).

# Contrails

## DESCRIPTION OF THE COMPUTER PROGRAM

The equations for the ray paths are solved in the following manner:  
Let the independent variable be  $y$  and

$$V_1 = \frac{dx}{dy}$$

$$V_2 = x$$

$$V_3 = t$$

Then

$$\frac{dV_1}{dy} = f_1(V_1, V_2, y)$$

$$\frac{dV_2}{dy} = V_1$$

$$\frac{dV_3}{dy} = f_3(V_1, V_2, y)$$

These three simultaneous differential equations are solved in a step-by-step manner by use of a standard "SHARE" subroutine which is based on the Runge Kutta method. When  $dx/dy$  becomes greater than one, a variable change takes place in the program, and  $x$  becomes the independent variable.

A signal (in the supersonic region) is considered "trapped" on the local Mach line when

$$|x'^2 - (M^2 - 1)| \leq E1$$

When, for this trapped signal,  $(M-1) < E2$ , the integration stops and a new ray line is started. This logical flow is shown in the chart on page 21.

The values of  $\mathcal{A}_0$  used in the program are determined by the parameter (NLA). If (NLA) is an odd integer, it will be rounded down in the program to an even integer. Values of  $\mathcal{A}_0$  vary from zero to  $\pi$  and from zero to  $-\pi$  in an arithmetic progression.

Computation of a ray path (other than for a "trapped signal") ceases under the following conditions:

$$x \leq 0$$

$$1 \leq x$$

$$|Y_{MAX}| \leq |y|$$

$$N_{MAX} \leq N_{CNT}$$

where  $N_{CNT}$  is the number of points on the ray path already computed. This logical flow is shown in the chart on page 22.

Subroutine DERIV computes the appropriate derivatives.

Subroutine CNTRL accomplishes variable changes, stores local values in appropriate locations for later printing, and performs exit tests.

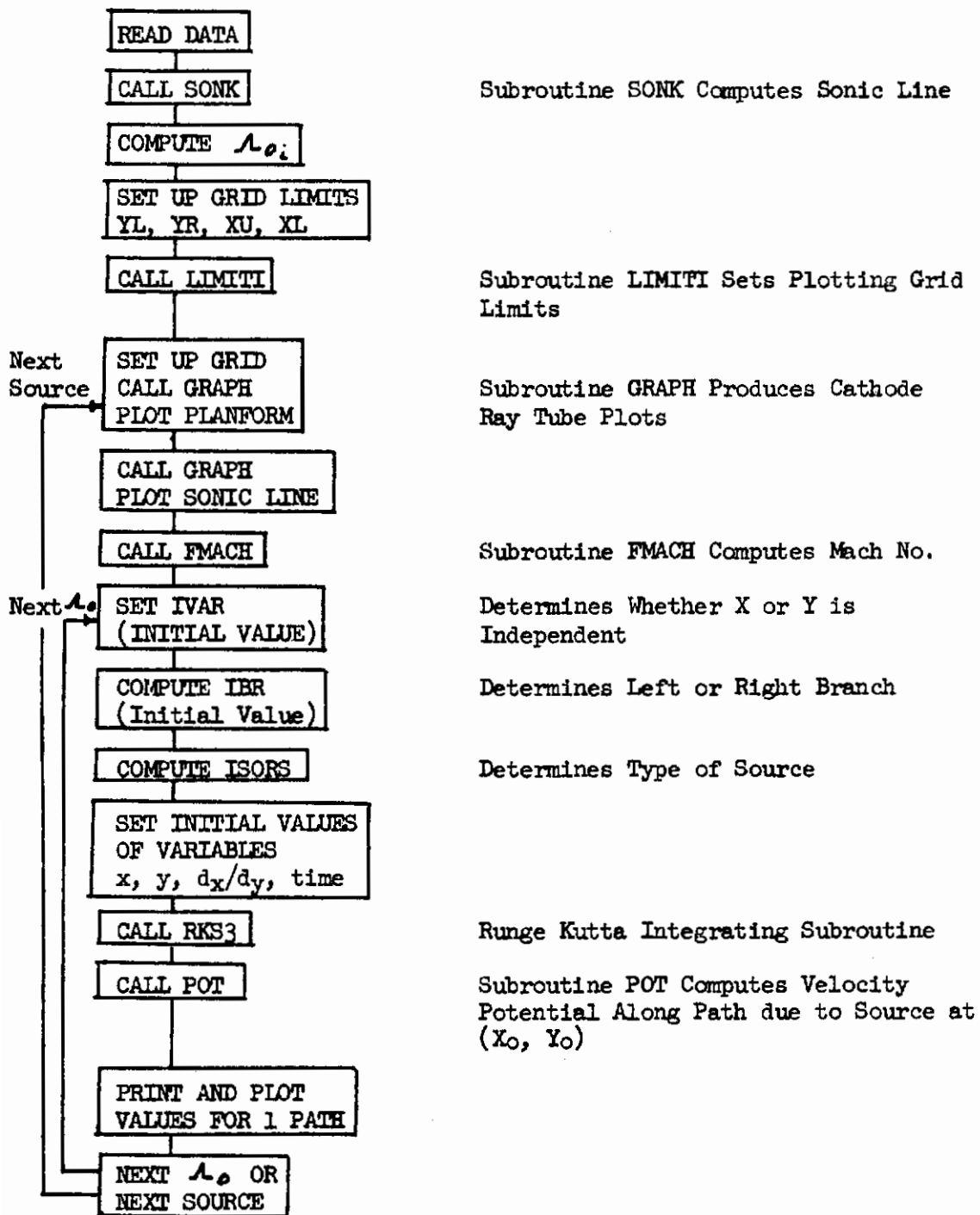
Subroutine FMACH computes the local Mach number and the partial derivatives of the Mach number.

# *Contrails*

Subroutine SONK computes coordinates on the planform where  $M = 1$ .

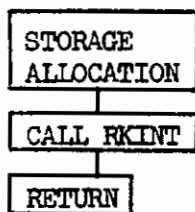
Sample data sheets with numbers which have been used in a computer run are in Appendix II. The output sheets are included. The output format is self-explanatory, with the exceptions of certain test words that are printed out at the beginning of the plots for each ray-path. Definitions for these words can be found in the comment statements at the beginning of the listing in Appendix I. The values listed for these test words apply to the last point plotted for the ray-path.

## MAIN PROGRAM

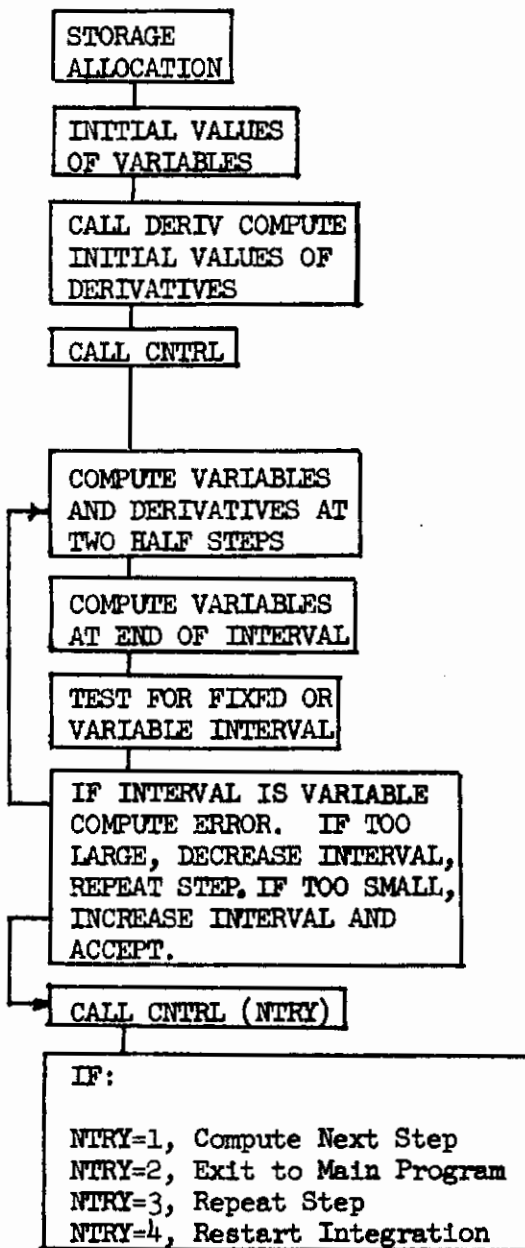


# Contrails

## SUBROUTINE RKS3



## SUBROUTINE RKINT



Subroutine DERIV Computes Derivatives

Subroutine CNTRL Executes Variable Changes, Stores Current Values, Executes Exit Tests

This Loop Calls DERIV 8 Times

NTRY is Re-set in CNTRL

## DISCUSSION OF RESULTS

This report contains two methods for calculating the velocity potential along sonic ray lines emanating from any point in a non-uniform flow field, i.e., one that varies from locally subsonic to supersonic speeds. Both methods apply to pulses emitted by sources or doublets. It has been demonstrated that both methods yield nearly identical ray paths and times of transmission. Those presented were obtained using the second method.

Figures 6 through 13 show ray paths of acoustic signals emanating from various points in a non-uniform transonic flow field. The reader may want to try his hand at tracing one of the ray paths in a region of interest such as near a leading edge. If so, it should be helpful to recall the discussion starting with Equation (7), through the difference equations of the path, Equation (8), and to the end of that section. An analysis of the differential equation of the path, Equation (24) should also be helpful. These show, for instance, that where the Mach number is constant the curvature of the ray path is zero; for a given Mach number and slope of ray path the curvature is proportional to rate of change of Mach number along the path. Figures 6, 7, 9, and 10 conclusively show that when the variation in Mach number is parabolic in the chordwise and spanwise directions focal points exist, both in subsonic and supersonic portions of the flow. None of the present theories accounts for the corresponding multiple crossings of the acoustic wave front. Figures 9 and 12 show acoustic signals traveling from regions of supersonic flow to regions of subsonic flow. This can occur, of course, only when the sonic line is swept downstream. Figures 9 and 12 also show rays that have been trapped on the Mach wave, travel outward to the sonic line where the spanwise slope of the ray path becomes zero, and are cancelled there. A study of the ray paths that cross the leading edge shows that in practical applications it is correct to assume they do not return.

These results permit the formulation of a numerical procedure. A box method is outlined in Appendix III. It establishes velocity potentials at all box centers on an aerodynamic surface and the corresponding generalized forces.

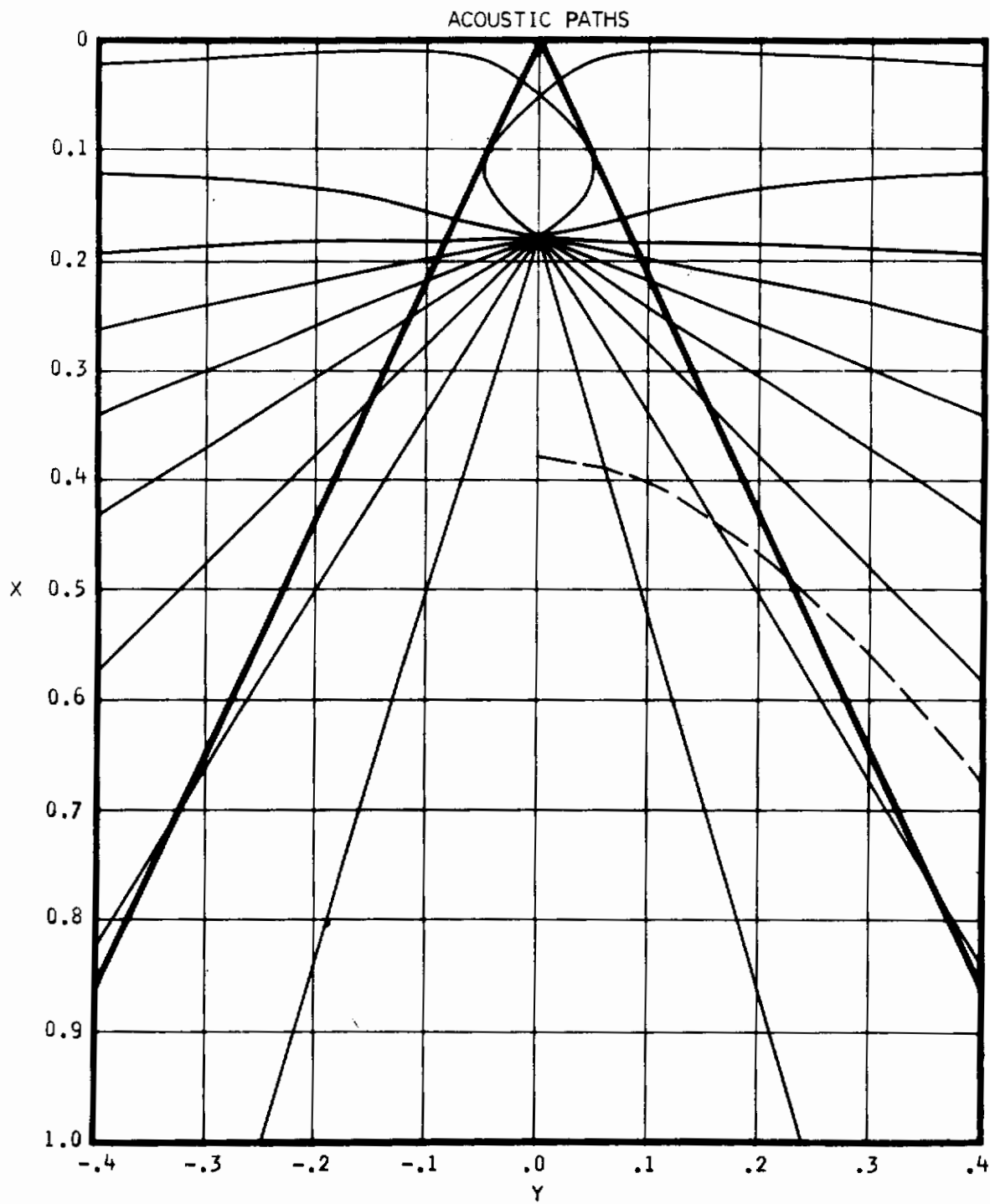


Figure 6. Ray Paths for a Source or Doublet at  $(0.18c, 0.0)$



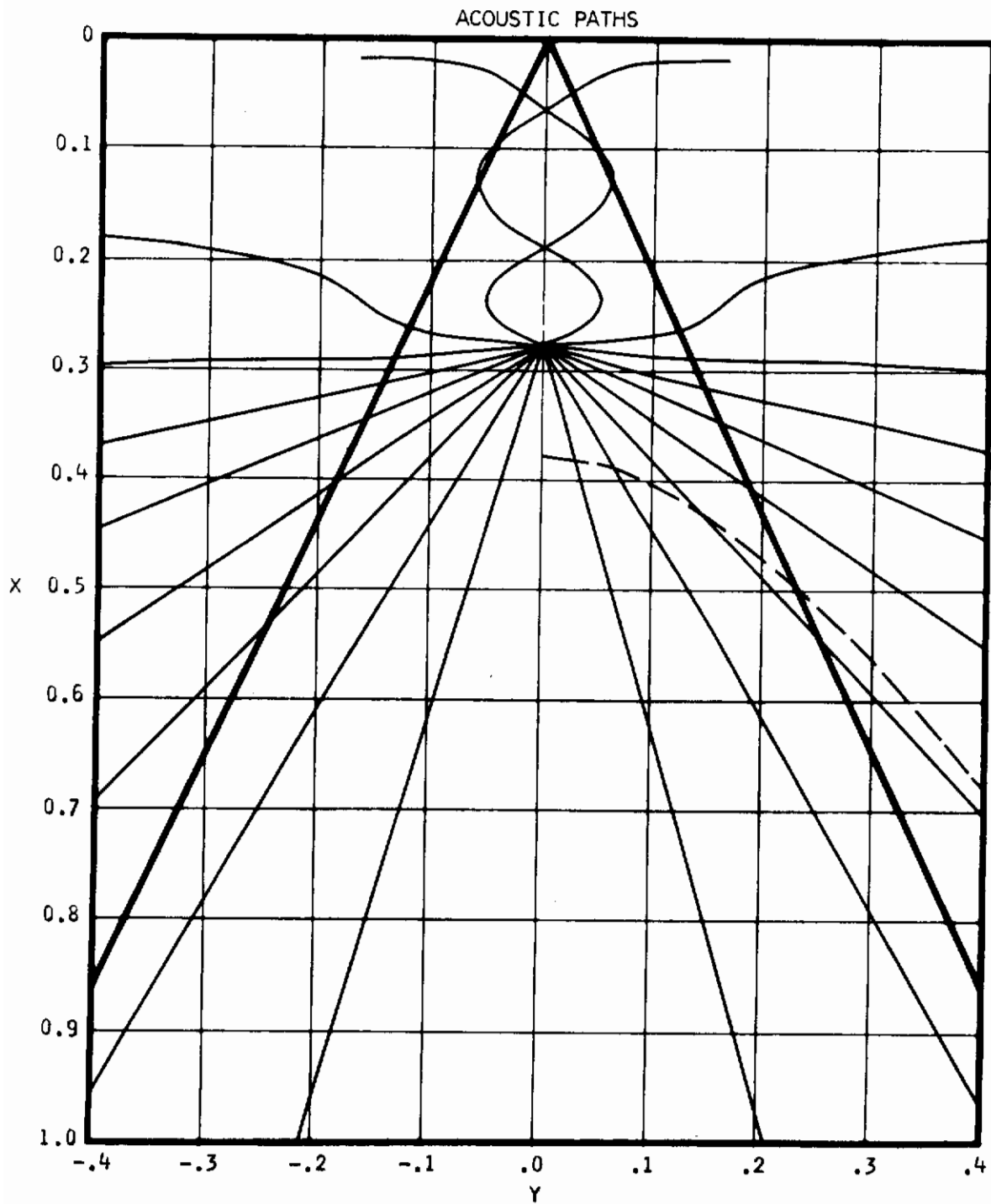


Figure 7. Ray Paths for a Source or Doublet at  $(0.28c, 0.0)$

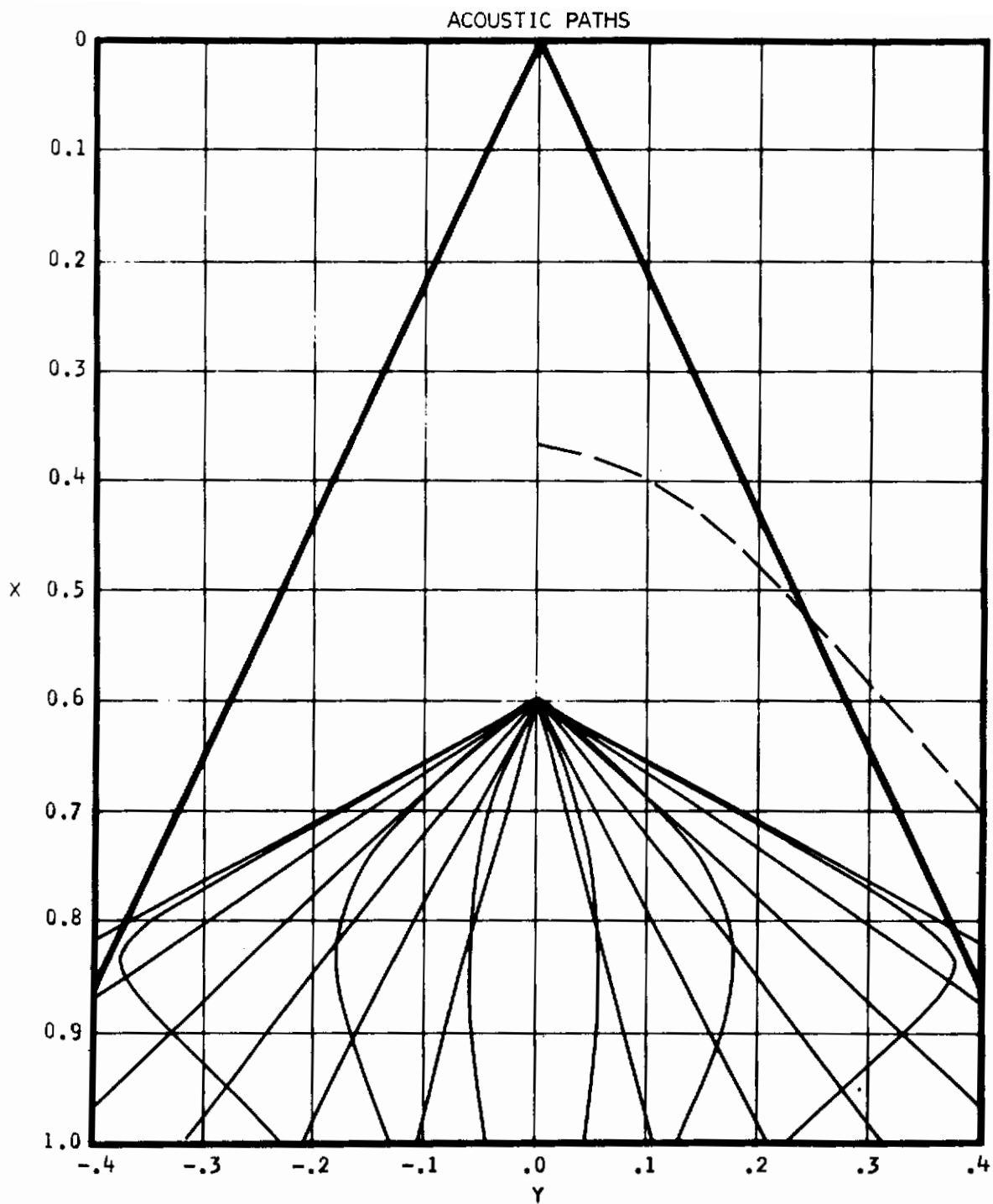


Figure 8. Ray Paths for a Source or Doublet at  $(0.6c, 0)$

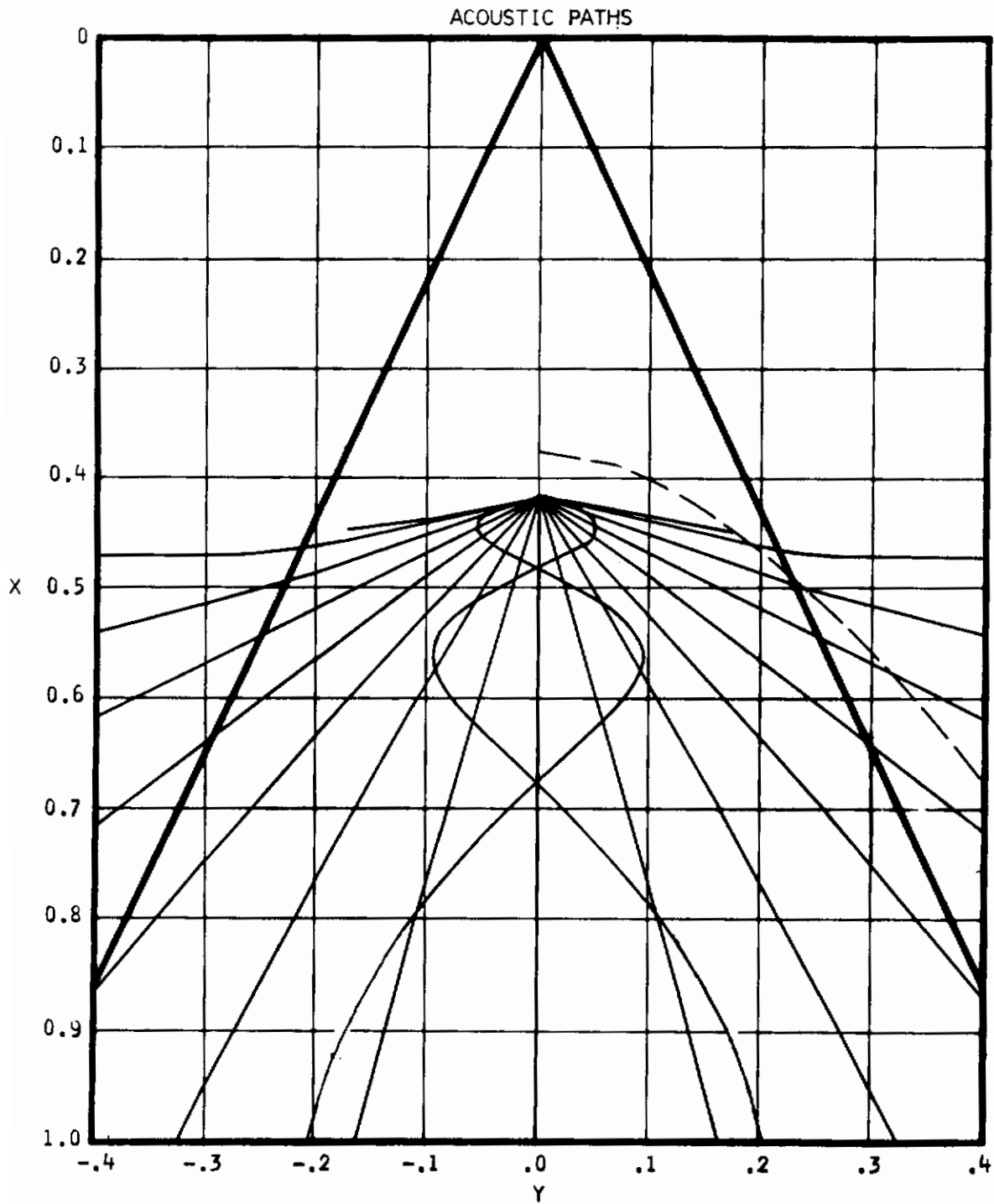


Figure 9. Ray Paths for a Source or Doublet at  $(0.42c, 0.0)$

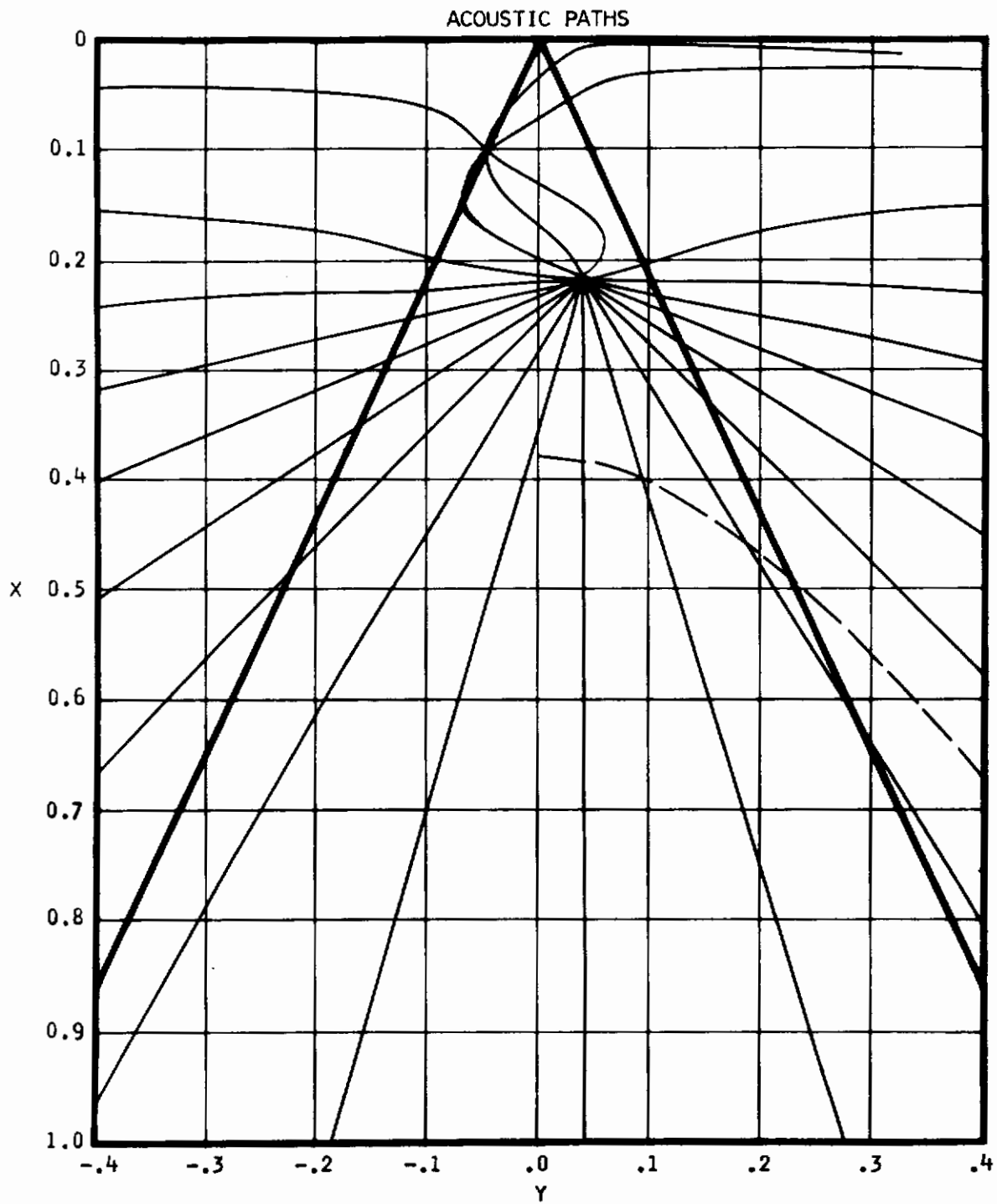


Figure 10. Ray Paths for a Source or Doublet at  $(0.22c, 0.04c)$

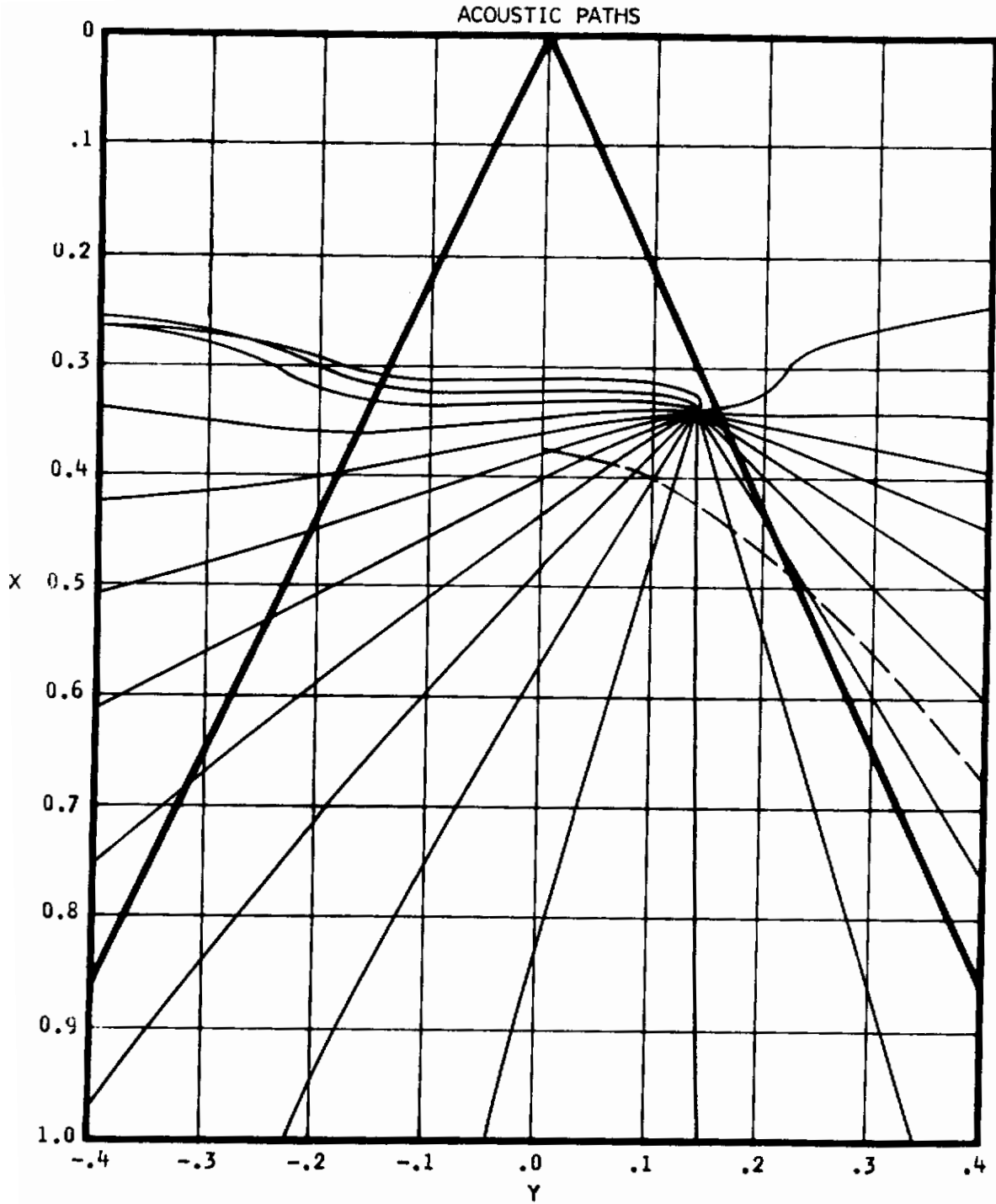


Figure 11. Ray Paths for a Source or Doublet at (0.34c, 0.14c)

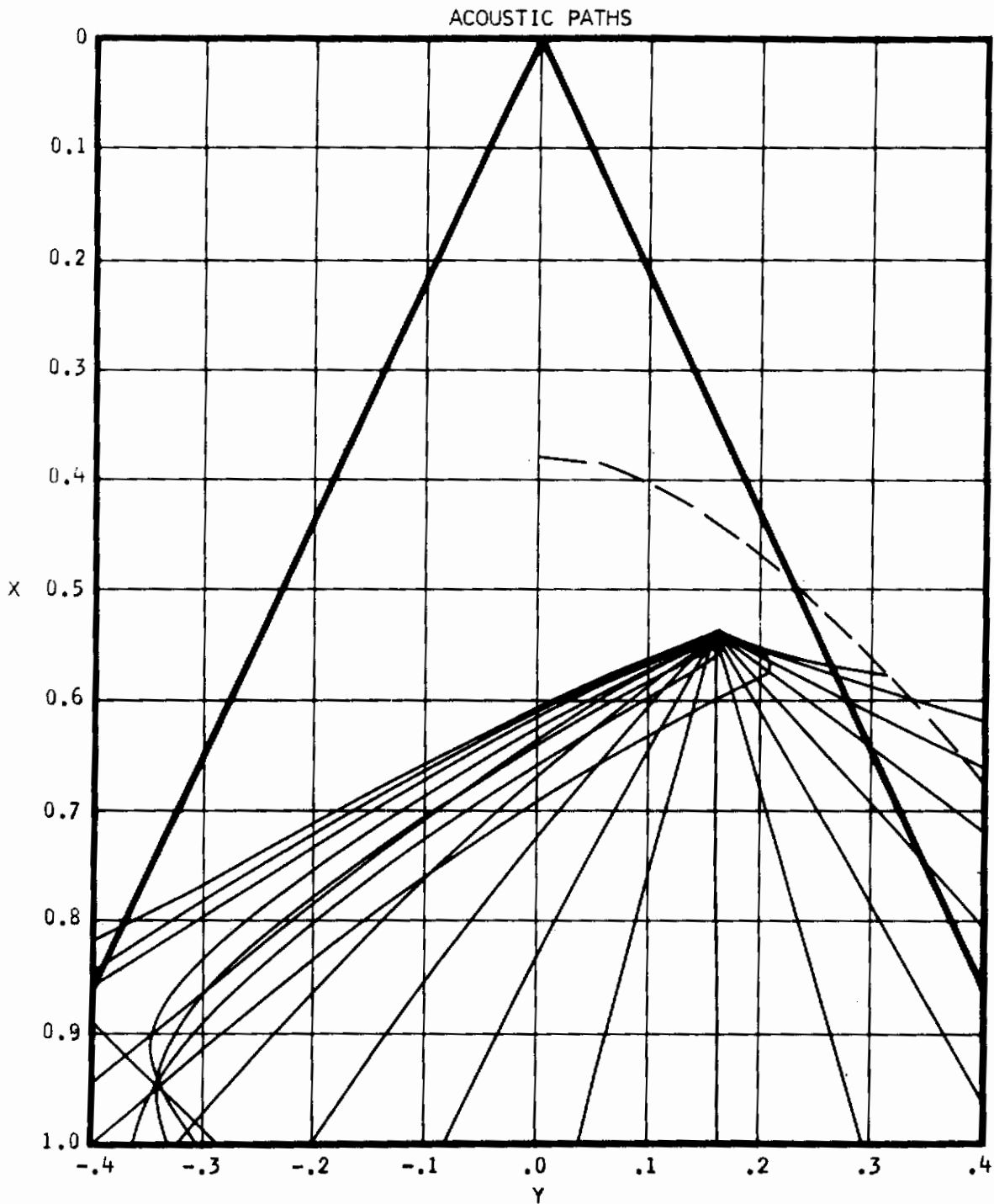


Figure 12. Ray Paths for a Source or Doublet at  $(0.54c, 0.16c)$

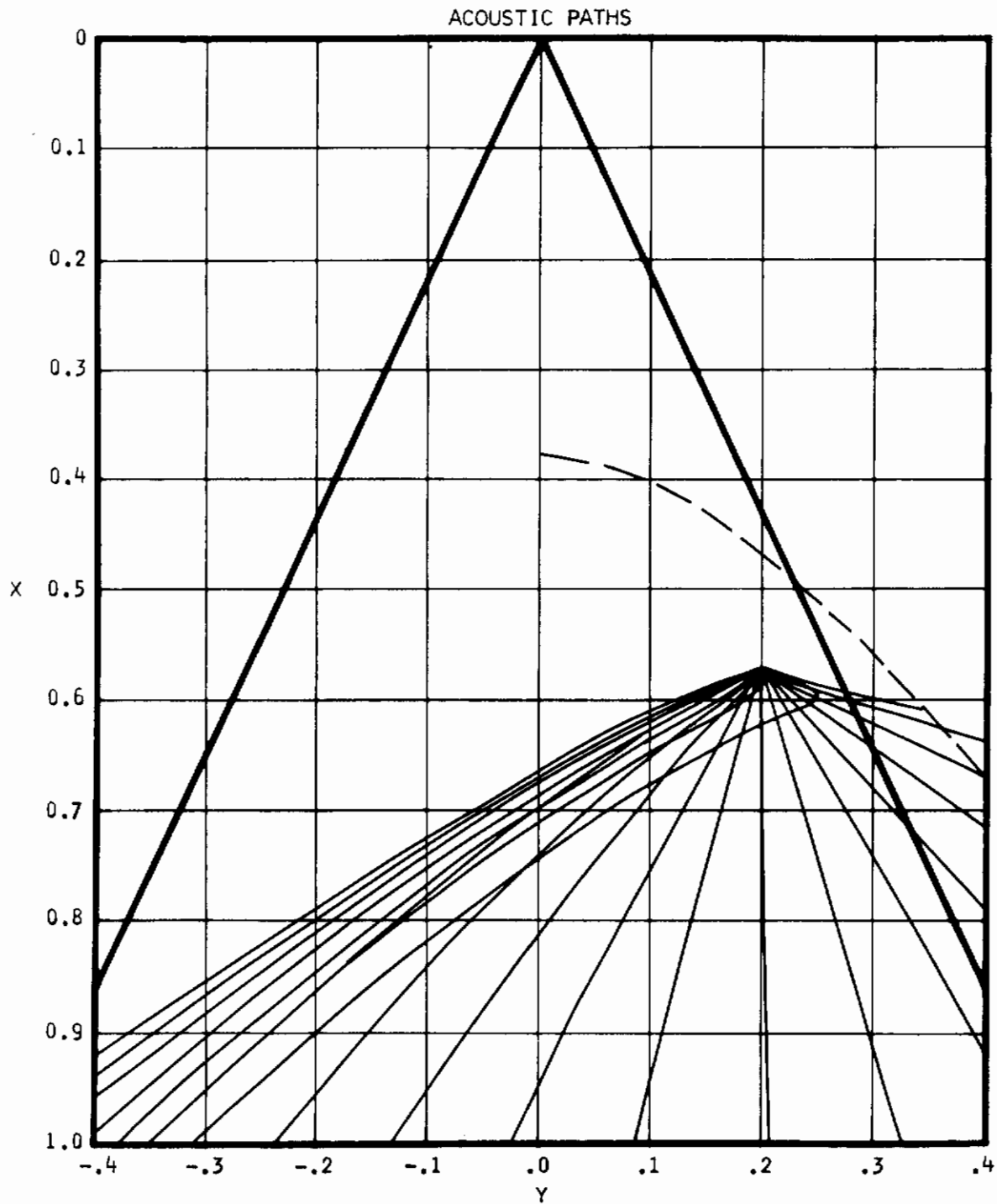


Figure 13. Ray Paths for a Source or Doublet at  $(0.57c, 0.20c)$

## CONCLUSIONS AND RECOMMENDATIONS

Two methods have been outlined in detail, and one of them has been completely mechanized for calculating the velocity potentials along acoustic ray paths emanating from any point in a non-uniform transonic flow field over a lifting surface. The one mechanized gives the ray path and velocity potential for the minimum time of travel from the source point to the field point.

To calculate pressures over the planform and generalized forces, it will be necessary to develop a procedure for calculating the velocity potential at an arbitrary point due to a sheet of sources, covering the wing surface, and the flow field in the plane of the wing out to a distance of several wing spans in the y-direction, or due to a sheet of doublets covering the wing surface. The latter is recommended for economy reasons.

The computer program in this report may be used to refine the doublet box method of Rodemich (3) in such a way as to include the (possibly very important) influence of wing thickness distribution on transonic airloads. A doublet box method similar to the one Rodemich developed (Reference 3) is recommended. The procedure is heuristically described in Appendix III. For each of a selected set of points in a sending box, the distribution of velocity potentials along ray lines throughout the zone of influence can be determined. An interpolation scheme will yield from these the velocity potentials at box centers and a numerical integration procedure will yield a velocity potential influence coefficient for each of the box centers. It will be necessary to solve a set of simultaneous equations to establish the strengths of doublets required to satisfy the tangential flow condition in the subsonic flow region. The order of the set will be equal to the number of box centers in the subsonic region on the wing. In the supersonic region the doublet strengths can be established sequentially. The use of doublets to solve unsteady supersonic flow problems has been outlined by Ashley in Reference 7.

It is recommended that this method be fully developed for the purpose of calculating generalized forces on wings in harmonic motion at transonic speeds. A computer program that would predict, with reasonable accuracy, the flutter characteristics and unsteady aerodynamic loads on a wing responding to externally applied forces, such as gusts, would fill an important gap in available technology.



## REFERENCES

1. Landahl, M. Uniformly Valid Approximate Solution for an Oscillating Source in a Transonic Flow Field. Unpublished paper, (1965).
2. Landahl, M. Approximate Solution for an Oscillating Source in a Non Uniform Transonic Stream. NAA, SID 63-1194 (August 1963)
3. Rodesich, E. R. and L. V. Andrew Unsteady Aerodynamics for Advanced Configurations, Part II - A Transonic Box Method for Planar, Lifting Surfaces. FIL-TDR-64-152, (May 1965)
4. Spreiter, J. R., and A. Y. Alksne, Thin Airfoil Theory Based on Approximate Solution of the Transonic Flow Equation, NACA Report 1359.
5. Rubbert, Paul E., Analysis of Transonic Flow by Means of Parametric Differentiation AFOSR 64-1932, MIT Fluid Dynamics Research Laboratory Report No. 65-2 (1965).
6. Wind Tunnel Tests of Four Reflection Plane Mounted .024 Scale Models Simulating The YB-70 Wing to Investigate the Effect of Camber on the Chordwise Pressure Distribution at Mach Numbers from 0.4 to 3.0. NA-61-55, TWT-50, (Unpublished)
7. Ashley, Holt, Further Studies on Aerodynamic Influence Coefficients For Supersonic Wings - Use of Doublet Sheets. NAA, SID 66-373 (July 1964).
8. Andrew, L. V. and Stenton, T. E., Unsteady Aerodynamic Forces on a Thin Wing Oscillating in Transonic Flow. AIAA Paper No. 67-16, January 23-26, 1967.

# Contrails

## APPENDIX I. Program Listings

```

$IBFTC MAIN      SDD                                          SNIC0005
C   FORTRAN PROGRAM TO COMPUTE (AND PLOT) THE PATHS OF ACOUSTIC SIG - SNIC0010
C   NALS (AND TRANSMISSION TIMES) ON AN AIRFOIL IN A SONIC FLOW FIELD, SNIC0015
C   ACCOUNTING FOR VARIATION IN LOCAL MACH NUMBER.              SNIC0020
C   CM = COEFFICIENTS OF MACH EQUATION. (SEE SUBROUTINE FMACH )  SNIC0025
C   PLX AND PLY ARE CONSTANTS DESCRIBING THE PLANFORM GEOMETRY. SNIC0030
C   THE PROGRAM ALLOWS FOR EITHER X OR Y TO BE THE INDEPENDENT VARIA- SNIC0035
C   BLE, DEPENDING ON THE CURRENT VALUE OF X-PRIME, WHICH SETS IVAR. SNIC0040
C       IF IVAR = 1,                IF IVAR = 2,              SNIC0045
C   YY = CURRENT VALUE OF X          YY = CURRENT VALUE OF Y  SNIC0050
C   DYY= CURRENT VALUE OF DX         DYY= CURRENT VALUE OF DY  SNIC0055
C   XX(1)= CURRENT VALUE OF Y-PRIME  XX(1)= CURRENT VALUE OF X-PRIME SNIC0060
C   XX(2) = CURRENT VALUE OF Y       XX(2) = CURRENT VALUE OF X  SNIC0065
C   XX(3) = CURRENT VALUE OF TIME    XX(3) = CURRENT VALUE OF TIME  SNIC0070
C   XX(4) = CURRENT VALUE OF R-BAR   XX(4) = CURRENT VALUE OF R-BAR  SNIC0075
C   DXX(1)= Y-DOUBLE PRIME           DXX(1)= X-DOUBLE PRIME       SNIC0080
C   DXX(2)= CURR. VALUE OF Y-PRIME   DXX(2)= CURR. VALUE OF X-PRIME  SNIC0085
C   DXX(3)= CURR. VALUE OF DT/DX     DXX(3)= CURR. VALUE OF DT/DY   SNIC0090
C   DXX(4)= CURRENT VALUE OF DR/DX   DXX(4)= CURRENT VALUE OF DR/DY  SNIC0095
C
C   IVAR IS ORIGINALLY SET IN MAIN PROGRAM, AND THEN RESET ON EACH SNIC0105
C   PASS THROUGH SUBROUTINE CNTRL.  SNIC0110
C
C   WORK = WORKING AREA FOR SUBROUTINE RKSS .                    SNIC0120
C   IFVD = FALSE AND IDKP= TRUE FOR VARIABLE INTERVAL.          SNIC0125
C   IFVD = TRUE FOR FIXED INTERVAL.                              SNIC0130
C
C   SX = VECTOR CONTAINING COMPUTED X- VALUES.                 SNIC0140
C   SXP = VECTOR CONTAINING COMPUTED X-PRIME VALUES.           SNIC0145
C   SY CONTAINS COMPUTED Y VALUES                              SNIC0150
C   SYP CONTAINS COMPUTED R-BAR VALUES                         SNIC0155
C   TIM CONTAINS TRANSMISSION TIMES.                            SNIC0160
C   FM = CURRENT MACH NUMBER                                    SNIC0165
C   ISORS = -1 DEFINES A SUPERSONIC SOURCE, RECEDING PATH.     SNIC0170
C   ISROS = 0 DEFINES A SUPERSONIC SOURCE, ADVANCING PATH.     SNIC0175
C   ISORS = 1 DEFINES A SUBSONIC SOURCE.                       SNIC0180
C   IBR = 1 FOR RIGHT BRANCH, 2 FOR LEFT                       SNIC0185
C   NCNT IS THE COUNTER FOR THE VECTORS SX,SY,SXP,SYP,TIM. WHEN NCNTSNIC0190
C   = NMAX, INTEGRATION STOPS, AND THE FLOW PASSES TO NEXT PATH SNIC0195
C   ITRAP =1 INDICATES SIGNAL IS TRAPPED ON THE LOCAL MACH CONE. SNIC0200
C   DZ IS INITIAL VALUE OF INCREMENT.                          SNIC0205
C   CINF = REMOTE SPEED OF SOUND IN ROOT CHORDS PER SECOND.    SNIC0210
C   FMINF= REMOTE MACH NUMBER                                    SNIC0215
C   POTE - THE POTE MATRIX CONTAINS THE VELOCITY POTENTIALS ALONG A SNIC0220
C   RAY PATH, NORMALIZED ON BD .                                SNIC0225
C   FREQ =ASSUMED FREQUENCIES IN RADIAN$ PER SECOND.           SNIC0230
C
C   SNIC0235
C   EXTERNAL DERIV, CNTRL                                       SNIC0240
C   COMMON                                                         SNIC0245
C   */WORK/ WORK(SD)                                             SNIC0250

```

# Contrails

```
*/XYZ/ SX(101),SXP(101),SY(101),SYP(101),AL(41),TIM(101) SNIC0255
*/XDX/ XX(4),DXX(4),YY,DYY,DZ SNIC0260
*/CM/ CM(6) SNIC0265
*/TABLE/ ATABL(4),RTABL(4) SNIC0270
*/PL/ PLX(8),PLY(8) SNIC0275
*/ICNT/ IVAR,NCNT,ISORS,IBR,ITRAP,NMAX SNIC0280
*/SOURCE/ XO(20),YO(20) SNIC0285
*/EPS/ E1,E2,FM,YMAX SNIC0290
*/NNN/ NSS,NLCS,NLLS SNIC0295
*/ECM/ ECM SNIC0300
*/C4/ CM2(7) SNIC0305
C SNIC0310
1000 FORMAT(2L12 ) SNIC0315
1010 FORMAT(6E12.8) SNIC0320
1020 FORMAT(6I12 ) SNIC0325
3 READ (5,1020) NSORCE,NLA,NPL,NMAX,NF SNIC0330
READ (5,1000) IFVD,IBKP SNIC0335
READ (5,1010) (XO(I),YO(I),I=1,NSORCE) SNIC0340
READ (5,1010) (CM(I),I=1,6 ) SNIC0345
READ (5,1010) DZ,E1,E2,YMAX SNIC0350
READ(5,1010) (ATABL(I),I=1,4),(RTABL(I),I=1,4 ) SNIC0355
READ (5,1010) (PLX(I),PLY(I),I=1,NPL ) SNIC0360
READ (5,1010) CINF,FMINF,TAU,TSAA SNIC0365
C SNIC0370
C TAU=MAX. (T/C), TSAA = TANGENT OF SEMI-APEX ANGLE SNIC0375
C DIMENSION FREQ(10),POTE(101,2,10 ) SNIC0380
C READ (5,1010) (FREQ(I),I=1,NF ) SNIC0385
C SNIC0390
DIMENSION XSO(40),YSO(40) SNIC0395
ECM = CINF*SQRT(5.0+FMINF**2) SNIC0400
ECM=1.0/ECM SNIC0405
CALL SONK(40,NXY,YMAX,YSO,XSO,IER ) SNIC0410
2000 FORMAT(49HD ERROR IN SUBROUTINE SONIC. CHECK MACH CONSTANTS ) SNIC0415
GO TO (1,2), IER SNIC0420
2 WRITE (6,2000) SNIC0425
1 CONTINUE SNIC0430
NVAR =4 SNIC0435
C NVAR IS THE NUMBER OF VARIABLES SNIC0440
CM2(1) =0.3 SNIC0445
CM2(2) =0.7 SNIC0450
CM2(3) =ATAN(1./TSAA) SNIC0455
CM2(4) =TAU SNIC0460
CM2(5) =1.18*TSAA SNIC0465
CM2(6) =.04 SNIC0470
CM2(7)=FMINF SNIC0472
C DEVELOP LAMDAS SNIC0475
NL=2*(NLA/2) SNIC0480
C THERE WILL ACTUALLY BE NL VALUES. IF NLA IS EVEN,NL=NLA. BUT NL= SNIC0485
C NLA - 1 IF NLA IS ODD. SNIC0490
NL1=NL-1 SNIC0495
```

# Contrails

```
NL2 =NL/2
XN= NL2*(NL2+1)
DG= 6.28318/XN
AL(1)=0.
DO 10 J=3,NL1,2
XJ = (J-1)/2
J1 = J-1
AL(J)=AL(J-2)+XJ*DG
10 AL(J1)=-AL(J)
AL(NL)= 3.14159
C SET UP GRID LIMITS
XU=0.
XL=1.
YL=-YMAX
YR = YMAX
CALL LIMITI(YL,YR,XL,XU)
DO 600 NS=1,NSORCE
NSS = NS
CALL GRAPH(1,42,-NPL,PLY,PLX,2H Y,2H X,15H ACOUSTIC PATHS )
XOF=XO(NS)
YOF=YO(NS)
CALL GRAPH(0,42,-NXY,YSO,XSO )
NLLS = NL
DO 500 NLC=1,NL
NLCs = NLC
ITRAP = 0
CALL FMACH (XOF,YOF,FM,FMX,FMY )
TEST1 =FM - COS(AL(NLC))
TEST2 = SIN(AL(NLC))
IF(NLC .NE. NL ) GO TO 11
IF (YOF .GT. 0. ) GO TO 11
TEST2 = -TEST2
11 IF (NLC-1) 14,12,14
12 IVAR=1
GO TO 30
14 IF(NLC-NL) 18,12,18
18 IF(TEST1) 22,20,22
20 IVAR=2
GO TO 30
22 TEST = TEST1/TEST2
ART = ABS(TEST)
IF(ART-1.0) 20,12,12
30 CONTINUE
C SET IBR
FL=AL(NLC)
IF(NLC-1) 32,31,32
31 IF(YOF) 41,41,42
32 IF(NLC-NL) 36,34,36
34 IF(YOF) 42,41,41
36 IF(FL) 42,42,41
```

SNIC0500  
SNIC0505  
SNIC0510  
SNIC0515  
SNIC0520  
SNIC0525  
SNIC0530  
SNIC0535  
SNIC0540  
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SNIC0550  
SNIC0555  
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SNIC0565  
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SNIC0580  
SNIC0585  
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SNIC0600  
SNIC0605  
SNIC0610  
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SNIC0640  
SNIC0645  
SNIC0650  
SNIC0655  
SNIC0660  
SNIC0665  
SNIC0670  
SNIC0675  
SNIC0680  
SNIC0685  
SNIC0690  
SNIC0695  
SNIC0700  
SNIC0705  
SNIC0710  
SNIC0715  
SNIC0720  
SNIC0725  
SNIC0730  
SNIC0735  
SNIC0740  
SNIC0745

# Contrails

|   |          |
|---|----------|
| 41 IBR=1  | SNIC0750 |
| GO TO 50  | SNIC0755 |
| 42 IBR=2  | SNIC0760 |
| 50 CONTINUE   | SNIC0765 |
| C    SET ISORS  | SNIC0770 |
| CSL =COS(FL)  | SNIC0775 |
| RM=1.0/FM   | SNIC0780 |
| IF (FM-1.0) 60,51,51  | SNIC0785 |
| 51 IF ((FM-1.0)-E2) 52,52,58  | SNIC0790 |
| 52 GO TO (53,54),IVAR   | SNIC0795 |
| 53 YPR =TEST2/TEST1   | SNIC0800 |
| TST= 1.0-YPR**2*(FM**2 - 1.0)   | SNIC0805 |
| 55 IF (TST-E1) 500,500,58   | SNIC0810 |
| 54 XPR = TEST1/TEST2  | SNIC0815 |
| TST= XPR**2-(FM**2-1.0)   | SNIC0820 |
| GO TO 55  | SNIC0825 |
| 58 IF (CSL-RM) 68,68,64   | SNIC0830 |
| 60 ISORS=1  | SNIC0835 |
| GO TO 70  | SNIC0840 |
| 64 ISORS= -1  | SNIC0845 |
| GO TO 70  | SNIC0850 |
| 68 ISORS = 0  | SNIC0855 |
| C   | SNIC0860 |
| 70 NCNT=1   | SNIC0865 |
| GO TO (80,90),IVAR  | SNIC0870 |
| C    IF IVAR=1,X IS THE INDEPENDENT VARIABLE.                             | SNIC0875 |
| C   | SNIC0880 |
| 80 YY = XOF   | SNIC0885 |
| IF (TEST1) 81,81,82   | SNIC0890 |
| 81 DYY=-DZ  | SNIC0895 |
| GO TO 83  | SNIC0900 |
| 82 DYY = DZ   | SNIC0905 |
| 83 XX(1) = TEST2/TEST1  | SNIC0910 |
| XX(2)=YOF   | SNIC0915 |
| XX(3) = 0.  | SNIC0920 |
| XX(4) = 0.  | SNIC0925 |
| GO TO 100   | SNIC0930 |
| C    IF IVAR=2,Y IS THE INDEPENDENT VARIABLE.                             | SNIC0935 |
| 90 YY = YOF   | SNIC0940 |
| GO TO (91,92),IBR   | SNIC0945 |
| 91 DYY = DZ   | SNIC0950 |
| GO TO 93  | SNIC0955 |
| 92 DYY = -DZ  | SNIC0960 |
| 93 XX(1) = TEST1/TEST2  | SNIC0965 |
| XX(2)= XOF  | SNIC0970 |
| XX(3) = 0.  | SNIC0975 |
| XX(4) = 0.  | SNIC0980 |
| 100 CALL RK53 (DERIV,CNTRL,XX,DXX,ATABL,RTABL,WORK,YY,DYY,NVAR,IFVD,IBS   | SNIC0985 |
| IKP,NTRY,IERR )   | SNIC0990 |
| 1070 FORMAT(1H1,30X,43H IVAR    NCNT    ISORS    IBR    ITRAP    NLCS = ) | SNIC0995 |

# Contrails

```
1080 FORMAT(1H0,27X, 6I7 ) SNIC1000
      WRITE (6,1070) SNIC1005
      WRITE (6,1080) IVAR,NCNT,ISORS,IBR,ITRAP,NLCS SNIC1010
C SNIC1015
1060 FORMAT(22H ERROR IN RKS3, IERR = I4 ) SNIC1020
      IF (IERR) 103,140,103 SNIC1025
      103 WRITE (6,1060) IERR SNIC1030
      GO TO 500 SNIC1035
1050 FORMAT(1H-,42X,4HXO = E16.8/ 43X,4HYO = E16.8/ 43X,1DHMACH NO. = ESNIC1040
      116.8// 29X,31H ACOUSTIC RAY PATH FOR LAMBDA = E16.8///17X,1HX,17X, SNIC1045
      11HY,14X,7HX-PRIME,11X,7HR-BAR ,12X,4HTIME// ) SNIC1050
1040 FORMAT(1H 7X,5E18.8) SNIC1055
      140 WRITE (6,1050) XO(NS),YO(NS),FM,FL SNIC1060
      WRITE (6,1040) (SX(I),SY(I),SXP(I),SYP(I),TIM(I),I=1,NCNT ) SNIC1065
C SNIC1070
      CALL GRAPH (D,NLC,-NCNT,SY,SX ) SNIC1075
      CALL POT (NF,FREQ,POTE ) SNIC1080
C SNIC1085
1100 FORMAT(1H1,25X,54H VELOCITY POTENTIALS ALONG A RAY PATH FOR A SOURSNIC1090
      1CE AT ) SNIC1095
1110 FORMAT(1H-,42X,4HXO = E16.8/43X, 4HYO = E16.8/ 43X, 8HLAMBDA = E16SNIC1100
      1.8 //39X,30HALTERNATING REAL AND IMAGINARY ) SNIC1105
1120 FORMAT (1H-,6X,7HOMEGA =E16.8// ) SNIC1110
1090 FORMAT(1H 6X,6E16.6) SNIC1115
C SNIC1120
      DO 300 N=1,NF SNIC1125
      IF (N .NE. 1 ) GO TO 200 SNIC1130
      WRITE (6,1100) SNIC1135
      WRITE (6,1110) XO(NS),YO(NS),FL SNIC1140
      200 WRITE (6,1120 ) FREQ(N) SNIC1145
      WRITE (6,1090) (( POTE(I,K,N),K=1,2 ),I=1,NCNT ) SNIC1150
      300 CONTINUE SNIC1155
      500 CONTINUE SNIC1160
      600 CONTINUE SNIC1165
      GO TO 3 SNIC1170
      END SNIC1175
```

# Contrails

|              |  |          |
|--------------|--|----------|
| \$IBFTC DERI | SDD                                    | SNIC1180 |
|              | SUBROUTINE DERIV                       | SNIC1185 |
|              | COMMON                                 | SNIC1190 |
|              | */XDX/ XX(4),DXX(4),YY,DYY,DZ          | SNIC1195 |
|              | */CM/ CM(6)                            | SNIC1200 |
|              | */ICNT/ IVAR,NCNT,ISORS,IBR,ITRAP,NMAX | SNIC1205 |
|              | */EPS/ E1,E2,FM,YMAX                   | SNIC1210 |
|              | */NNN/ NSS,NLCS,NLLS                   | SNIC1215 |
|              | */ECM/ ECM                             | SNIC1220 |
| C            |  | SNIC1225 |
|              | GO TO(10,50),IVAR                      | SNIC1230 |
| C            | X IS THE INDEPENDENT VARIABLE          | SNIC1235 |
| 10           | CALL FMACH(YY,XX(2),FM,FMX,FMY)        | SNIC1240 |
|              | R=XX(1)                                | SNIC1245 |
|              | DXX(2)=R                               | SNIC1250 |
|              | B=FM*FM-1.0                            | SNIC1255 |
|              | TSI=1.0-R*R*B                          | SNIC1260 |
|              | A=FM*FM+5.0                            | SNIC1265 |
|              | SA=SQRT(A)                             | SNIC1270 |
|              | IF(B) 103,103,101                      | SNIC1275 |
| 101          | BETA=SQRT(B)                           | SNIC1280 |
|              | IF(ITRAP.EQ.1) GO TO 104               | SNIC1285 |
|              | IF(ISORS.EQ.1) GO TO 103               | SNIC1290 |
|              | IF(TSI.GT.E1) GO TO 103                | SNIC1295 |
|              | ITRAP=1                                | SNIC1300 |
|              | GO TO 104                              | SNIC1305 |
| 103          | IF(TSI.GE.0.) GO TO 215                | SNIC1310 |
|              | ITRAP=2                                | SNIC1315 |
|              | TSI=0.                                 | SNIC1320 |
| 215          | RAD=SQRT(TSI)                          | SNIC1325 |
|              | DXX(4)=RAD                             | SNIC1330 |
|              | RAB=1.0/(A*B)                          | SNIC1335 |
|              | TM1=FM*(FM**2+11.0)/B                  | SNIC1340 |
|              | TM2=2.0*FM*(FM**2+8.0)*R**2            | SNIC1345 |
|              | TM3=((RAD**3)/B)*(7.0*FM**2+5.0)       | SNIC1350 |
|              | TM4=(FM/A)*R*(6.0*R**2+1.0)            | SNIC1355 |
|              | GO TO 105                              | SNIC1360 |
| 104          | RDB=1.0/(B**2)                         | SNIC1365 |
|              | DXX(4)=0.                              | SNIC1370 |
| 105          | IF(ISORS) 11,15,18                     | SNIC1375 |
| 11           | GO TO(12,13),IBR                       | SNIC1380 |
| 12           | IF(ITRAP) 91,91,7                      | SNIC1385 |
| 13           | IF(ITRAP) 92,92,8                      | SNIC1390 |
| 15           | GO TO(16,17),IBR                       | SNIC1395 |
| 16           | IF(ITRAP) 92,92,7                      | SNIC1400 |
| 17           | IF(ITRAP) 91,91,8                      | SNIC1405 |
| 18           | GO TO(92,91),IBR                       | SNIC1410 |
| 91           | IF(R) 4,3,3                            | SNIC1415 |
| 92           | IF(R) 3,3,4                            | SNIC1420 |
| C            | Y IS THE INDEPENDENT VARIABLE          | SNIC1425 |

# Contrails

|   |          |
|---|----------|
| 50 CALL FMACH (XX(2),YY,FM,FMX,FMY )                          | SNIC1430 |
| R = XX(1)   | SNIC1435 |
| DXX(2) = R  | SNIC1440 |
| B = FM*FM-1.0   | SNIC1445 |
| TSI = R*R-B   | SNIC1450 |
| C   | SNIC1455 |
| A = 5.0+FM*FM   | SNIC1460 |
| SA = SQRT(A)  | SNIC1465 |
| IF(B .LT. 0.) GO TO 108                                       | SNIC1470 |
| 106 BETA = SQRT(B)  | SNIC1475 |
| IF(ITRAP .EQ. 1) GO TO 109                                    | SNIC1480 |
| IF(ISORS .EQ. 1) GO TO 108                                    | SNIC1485 |
| IF(TSI .GT. E1) GO TO 108                                     | SNIC1490 |
| ITRAP = 1   | SNIC1495 |
| GO TO 109   | SNIC1500 |
| 108 IF(TSI .GE. 0.) GO TO 107                                 | SNIC1505 |
| TSI=0.  | SNIC1510 |
| ITRAP = 2   | SNIC1515 |
| 107 RAD = SQRT(TSI)   | SNIC1520 |
| DXX(4) = RAD  | SNIC1525 |
| RAB = 1.0/(A*B)   | SNIC1530 |
| TM1=(FM/B)*(FM**2+11.0)*R**3                                  | SNIC1535 |
| TM2= 2.0*FM*(FM**2+8.0)*R                                     | SNIC1540 |
| TM3=(RAD**3/B)*(7.0*FM**2+5.0)                                | SNIC1545 |
| TM4 = (FM/A)*(R**2+6.0)                                       | SNIC1550 |
| GO TO 110   | SNIC1555 |
| 109 DXX(4) = 0.   | SNIC1560 |
| 110 IF(ISORS) 52,60,68  | SNIC1565 |
| 52 GO TO (54,56),IBR  | SNIC1570 |
| 54 IF(ITRAP ) 1,1,5   | SNIC1575 |
| 56 IF(ITRAP ) 2,2,6   | SNIC1580 |
| 60 GO TO (62,64),IBR  | SNIC1585 |
| 62 IF(ITRAP ) 2,2,5   | SNIC1590 |
| 64 IF(ITRAP ) 1,1,6   | SNIC1595 |
| 68 GO TO (2,1),IBR  | SNIC1600 |
| C   | SNIC1605 |
| FORMULAS FOR THE SECOND DERIVS FOLLOW                         | SNIC1610 |
| C   | SNIC1615 |
| 1 IF(ABS(B) .LE. 1.E-03) GO TO 220                            | SNIC1620 |
| DXX(1)=RAB*(-TM1 + TM2 - TM3)*FMY + TM4 *FMX                  | SNIC1625 |
| DXX(3)=(SA*ECM/B)*(FM*XX(1)+RAD)                              | SNIC1630 |
| GO TO 100   | SNIC1635 |
| 2 IF(ABS(B) .GT. 1.E-03 ) GO TO 209                           | SNIC1640 |
| 220 DXX(1)=(.5/A)*(2.*R**3+R+9./R)*FMY + (FM/A)*(R**2+6.)*FMX | SNIC1645 |
| DXX(3)=(1.22475*ECM)*(R+(1./R))                               | SNIC1650 |
| GO TO 100   | SNIC1655 |
| 209 DXX(1)= RAB*(-TM1+TM2+TM3)*FMY+TM4 * FMX                  | SNIC1660 |
| DXX(3)=(SA*ECM/B)*(FM*XX(1)-RAD)                              | SNIC1665 |
| GO TO 100   | SNIC1670 |
| 3 IF(NLCS .EQ. NLLS) GO TO 4                                  | SNIC1675 |
| DXX(1)=RAB*(TM1-TM2+TM3)*FMY -TM4* FMX                        |          |



# Contrails

|   |          |
|---|----------|
| DXX (3) = (SA*ECM/B) * (FM + RAD)   | SNIC1680 |
| GO TO 100   | SNIC1685 |
| 4 IF (ABS (B) .GT. 1.E-03 ) GO TO 205   | SNIC1690 |
| 204 DXX (1) = - (.5/A) * (9.*R**4 + R**2 + 2.) * FMY - (R/A) * (6.*R**2 + 1.) * FMX | SNIC1695 |
| DXX (3) = (1.22475*ECM) * (1. + R*R)  | SNIC1700 |
| GO TO 100   | SNIC1705 |
| 205 DXX (1) = RAB * (TM1 - TM2 - TM3) * FMY - TM4 * FMX                             | SNIC1710 |
| DXX (3) = (SA*ECM/B) * (FM - RAD)   | SNIC1715 |
| GO TO 100   | SNIC1720 |
| 5 DXX (1) = FM * ((FMY/BETA) + FMX)   | SNIC1725 |
| DXX (3) = (SA*ECM/B) * FM * XX (1)  | SNIC1730 |
| GO TO 100   | SNIC1735 |
| 6 DXX (1) = FM * ((-FMY/BETA) + FMX )   | SNIC1740 |
| DXX (3) = (SA*ECM/B) * FM * XX (1)  | SNIC1745 |
| GO TO 100   | SNIC1750 |
| 7 DXX (1) = - (FM * RBB) * (FMY + BETA * FMX )                                      | SNIC1755 |
| DXX (3) = (SA*ECM/B) * FM   | SNIC1760 |
| GO TO 100   | SNIC1765 |
| 8 DXX (1) = FM * RBB * (-FMY + BETA * FMX)  | SNIC1770 |
| DXX (3) = (SA*ECM/B) * FM   | SNIC1775 |
| 100 IF (DYY .LT. 0. ) GO TO 31  | SNIC1780 |
| DXX (3) = ABS (DXX (3))   | SNIC1785 |
| GO TO 32  | SNIC1790 |
| 31 DXX (3) = - 1.0 * (ABS (DXX (3)))  | SNIC1795 |
| DXX (4) = - 1.0 * (ABS (DXX (4)))   | SNIC1800 |
| 32 RETURN   | SNIC1805 |
| END   | SNIC1810 |

# Contrails

```
$IBFTC CONT   SDD
  SUBROUTINE CNTRL(NTRY)
  COMMON
  */XYZ/ SX(101),SXP(101),SY(101),SYP(101),AL(41),TIM(101)
  */XDX/ XX(4),DXX(4),YY,DYY,DZ
  */CM/ CM(6)
  */ICNT/ IVAR,NCNT,ISORS,IBR,ITRAP,NMAX
  */EPS/ E1,E2,FM,YMAX
  */NNN/ NSS,NLCS,NLLS
  IF(NCNT.NE.1) GO TO 6
  NCO = 1
  IF(NR.EQ.1) GO TO 6
  NR = 1
  IF(ABS(DXX(1)*DYY) .LE. .25) GO TO 6
  4 DYY = .5*DYY
  IF(ABS(DXX(1)*DYY) .LE. .25) GO TO 7
  GO TO 4
  7 NTRY = 4
  RETURN
  6 IF(ABS(XX(1)).LT.1.0) GO TO 20
  1 NTRY =4
  GO TO (2,3),IVAR
  2 IVAR=2
  GO TO 5
  3 IVAR=1
C   SWITCH VARIABLES,SET NEW INITIAL CONDITIONS
  5 SAV =YY
  DYY = DYY*XX(1)
  10 YY = XX(2)
  XX(1)=1.0/XX(1)
  XX(2)=SAV
  RETURN
  20 GO TO (25,35),IVAR
C   STORE CURRENT VALUES WHERE X IS INDEPENDENT VARIABLE.
  25 SX(NCNT) = YY
C   CHANGE IBR WHEN Y-PRIM PASSES THROUGH ZERO
  IF(ABS(XX(1)) .GT. 1.0E-02) GO TO 15
  IF((DXX(1)*DYY*XX(1)) .GE. 0.0) GO TO 15
  IF(NCO .EQ. 2) GO TO 15
  NCO = 2
  XX(1)=-XX(1)
  NTRY = 4
  GO TO (11,12), IBR
  11 IBR =2
  GO TO 19
  12 IBR = 1
  GO TO 19
  15 IF(NCO .NE. 2 ) GO TO 19
  IF(ABS(XX(1)) .LT. 1.0E-01) GO TO 19
  NCO = 1
```

SNIC1815  
SNIC1820  
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SNIC2045  
SNIC2050  
SNIC2055  
SNIC2060

# Contrails

|                                    |          |
|------------------------------------|----------|
| 19 IF (XX(1) .NE. 0.0 ) GO TO 27   | SNIC2065 |
| 26 SXP (NCNT) = UNDEF              | SNIC2070 |
| GO TO 28                           | SNIC2075 |
| 27 SXP (NCNT) = 1.0/XX(1)          | SNIC2080 |
| 28 SY (NCNT) = XX(2)               | SNIC2085 |
| SYP (NCNT) = XX(4)                 | SNIC2090 |
| TIM (NCNT) = XX(3)                 | SNIC2095 |
| GO TO 50                           | SNIC2100 |
| 35 SX (NCNT) = XX(2)               | SNIC2105 |
| TIM (NCNT) = XX(3)                 | SNIC2110 |
| SXP (NCNT) = XX(1)                 | SNIC2115 |
| SY (NCNT) = YY                     | SNIC2120 |
| SYP (NCNT) = XX(4)                 | SNIC2125 |
| 50 CONTINUE                        | SNIC2130 |
| C NOW TEST FOR EXIT CONDITIONS     | SNIC2135 |
| IF (ITRAP .NE. 2) GO TO 51         | SNIC2140 |
| ITRAP = 0                          | SNIC2145 |
| NCNT = NCNT - 1                    | SNIC2150 |
| GO TO 100                          | SNIC2155 |
| 51 IF (ITRAP) 60,60,52             | SNIC2160 |
| 52 TEST = FM - 1.0                 | SNIC2165 |
| IF (TEST) 100,100,53               | SNIC2170 |
| 53 IF (TEST - E2) 100,100,60       | SNIC2175 |
| 60 IF (SX (NCNT)) 100,70,70        | SNIC2180 |
| 70 IF (SX (NCNT) - 1.0) 80,100,100 | SNIC2185 |
| 80 AY = ABS (SY (NCNT))            | SNIC2190 |
| IF (AY - YMAX) 105,100,100         | SNIC2195 |
| 105 IF (NCNT - NMAX) 110,100,100   | SNIC2200 |
| 100 NTRY = 2                       | SNIC2205 |
| NR = 0                             | SNIC2210 |
| RETURN                             | SNIC2215 |
| 110 NCNT = NCNT + 1                | SNIC2220 |
| RETURN                             | SNIC2225 |
| END                                | SNIC2230 |

# Contrails

|   |          |
|---|----------|
| 3IBFTC MACH   | SNIC2235 |
| C MASTER SUBR., M, MX, MY   | SNIC2245 |
| SUBROUTINE FMACH(FX,FY,FMS,FMXS,FMYS)                             | SNIC2240 |
| COMMON  | SNIC2250 |
| */C4/ CM2(7)  | SNIC2255 |
| EQUIVALENCE (A,CM2(1)),(B,CM2(2)), (AL,CM2(3)), (TAU,CM2(4)),(AK, | SNIC2260 |
| * CM2(5)), (R1,CM2(6)) ,(FMINF,CM2(7))                            | SNIC2265 |
| AY=ABS(FY)  | SNIC2275 |
| AYY = ABS(AK*FX)  | SNIC2280 |
| IF (AY .LE. AYY) GO TO 200  | SNIC2285 |
| SK = 1./ (SQRT(1.+AK*AK))   | SNIC2290 |
| T =(AY-AYY)*SK  | SNIC2295 |
| 100 CALL FMAC1 (FX,AYY,FMS,FMXS,FMYS)                             | SNIC2300 |
| CALL FMAC2 (FX,AY,A,B,AL,TAU,D1FM,D1MX,D1MY )                     | SNIC2305 |
| CALL FMAC2 (FX,AYY,A,B,AL,TAU,D2FM,D2MX,D2MY)                     | SNIC2310 |
| C   | SNIC2315 |
| FMS = FMS -0.6*FMINF*(D1FM-D2FM)                                  | SNIC2320 |
| FMXS= FMXS+FMYS*AK-0.6*FMINF*(D1MX-D2MX-AK*D2MY)                  | SNIC2325 |
| FMYS = -0.6*FMINF*D1MY*(AY/FY)                                    | SNIC2330 |
| IF (T .GE. R1) GO TO 300  | SNIC2335 |
| 120 CALL FMAC1 (FX,FY,SM,SMX,SMY )                                | SNIC2340 |
| ARG =1.57079*T/R1   | SNIC2345 |
| SI = SIN(ARG)   | SNIC2350 |
| SMO = SI*SI   | SNIC2355 |
| FMS=(FMS-SM)*SMO + SM   | SNIC2360 |
| FMXS=(FMXS-SMX)*SMO + SMX   | SNIC2365 |
| FMYS=(FMYS-SMY)*SMO +SMY  | SNIC2370 |
| GO TO 300   | SNIC2375 |
| 200 CALL FMAC1 (FX,FY,FMS,FMXS,FMYS)                              | SNIC2380 |
| 300 CONTINUE  | SNIC2385 |
| RETURN  | SNIC2390 |
| END   | SNIC2395 |

# Contrails

|    |  |          |
|----|--|----------|
|    | \$IBFTC MAC2 SDD   | SNIC2400 |
|    | SUBROUTINE FMAC2(X,Y,A,B,AL,TAU,DELCP,DDXCP,DDYCP)                 | SNIC2405 |
| C  |  | SNIC2410 |
| C  | SUBROUTINE COMPUTES DELTA CP                                       | SNIC2415 |
| C  |  | SNIC2420 |
|    | CS= COS(AL)  | SNIC2425 |
|    | CS1=1./ (SQRT(1.+((1.-A)**2)*(CS**2)))                             | SNIC2430 |
|    | CS2=1./ (SQRT(1.+((1.-B)**2)*(CS**2)))                             | SNIC2435 |
|    | TA = SIN(AL)/CS  | SNIC2440 |
|    | TA1=(1.-A)*TA  | SNIC2445 |
|    | TA2=(1.-B)*TA  | SNIC2450 |
|    | EPS=TAU/(2.*3.1415927*A*CS)  | SNIC2455 |
|    | EPS1= EPS*CS/CS1   | SNIC2460 |
|    | EPS2 = EPS*A*CS/((1.-B)*CS2)                                       | SNIC2465 |
|    | EDS = 1.0 - EPS  | SNIC2470 |
|    | EDS1 = EPS1 + 1.0  | SNIC2475 |
|    | EDS2 = EPS2 + 1.0  | SNIC2480 |
|    | S = ABS(X/TA)  | SNIC2485 |
|    | S1=(X-A)/TA1   | SNIC2490 |
|    | S2=(X-B)/TA2   | SNIC2495 |
|    | Q1=ABS(Y-S)  | SNIC2500 |
|    | Q2=ABS(Y+S)  | SNIC2505 |
|    | Q3=ABS(Y-S1)   | SNIC2510 |
|    | Q4=ABS(Y+S1)   | SNIC2515 |
|    | Q5 =ABS(Y-S2)  | SNIC2520 |
|    | Q6 =ABS(Y+S2)  | SNIC2525 |
|    | FAC =2.*CS/TA  | SNIC2530 |
|    | FAC1=2.*CS1/TA1  | SNIC2535 |
|    | FAC2=2.*CS2/TA2  | SNIC2540 |
|    | DEL =-FAC*(Q1**EPS+Q2**EPS-2.*S**EPS)                              | SNIC2545 |
|    | DDX=-FAC*(-1./(Q1**EDS)+1./(Q2**EDS))-2./(S**EDS) * EPS /TA        | SNIC2550 |
|    | DDY =-FAC*(1./(Q1**EDS)+1./(Q2**EDS)) * EPS                        | SNIC2555 |
|    | IF (S1) 10,10,5  | SNIC2560 |
| 10 | DELCP= DEL   | SNIC2565 |
|    | DDXCP= DDX   | SNIC2570 |
|    | DDYCP= DDY   | SNIC2575 |
|    | GO TO 50   | SNIC2580 |
| 5  | DEL1=-FAC1*(1./(Q3**EPS1)+1./(Q4**EPS1)-2./(S1**EPS1))             | SNIC2585 |
|    | DDX1=FAC1*(-1./(Q3**EDS1)+1./(Q4**EDS1)-2./(S1**EDS1)) * EPS1 /TA1 | SNIC2590 |
|    | DDY1= FAC1*(1./(Q3**EDS1)+1./(Q4**EDS1)) * EPS1                    | SNIC2595 |
|    | IF (S2) 20,20,30   | SNIC2600 |
| 20 | DELCP= DEL+ DEL1   | SNIC2605 |
|    | DDXCP= DDX+ DDX1   | SNIC2610 |
|    | DDYCP= DDY+ DDY1   | SNIC2615 |
|    | GO TO 50   | SNIC2620 |
| 30 | DEL2=-FAC2*(1./(Q5**EPS2)+1./(Q6**EPS2)-2./(S2**EPS2))             | SNIC2625 |
|    | DDX2=FAC2*(-1./(Q5**EDS2)+1./(Q6**EDS2)-2./(S2**EDS2)) * EPS2 /TA2 | SNIC2630 |
|    | DDY2= FAC2*(1./(Q5**EDS2)+1./(Q6**EDS2)) * EPS2                    | SNIC2635 |
|    | DELCP = DEL+ DEL1 + DEL2   | SNIC2640 |
|    | DDXCP= DDX+ DDX1 + DDX2  | SNIC2645 |

# Contrails

```
DDYCP= DDY+ DDY1 + DDY2  
50 RETURN  
END
```

```
SNIC2650  
SNIC2655  
SNIC2660
```

# Contrails

```
$IBFTC MAC1      SDD                                SNIC2665
      SUBROUTINE FMAC1 (FX,FY,FMS,FMXS,FMYS)          SNIC2670
C                                                    SNIC2675
C      SUBROUTINE COMPUTES MACH NO, MX, MY.          SNIC2680
C      FX = X                                        FY = Y                                SNIC2685
C      FMS = MACH NO.                               FMXS= PARTIAL M W/RESP TO X SNIC2690
C      FMYS= PARTIAL M W/RESP TO Y                  SNIC2695
C      EQ. FOR MACH IS M=CM(2)+EXP(-CM(1)*Y**2/X)*(CM(3)*X+CM(4)*X**2+ SNIC2700
C      CM(5)*Y**2+CM(6)*Y**4 )                       SNIC2705
C      COMMON                                        SNIC2710
C      */CM/ CM(6)                                   SNIC2715
C                                                    SNIC2720
C      EQUIVALENCE                                   SNIC2725
C      1      ( C ,CM(1)), ( FMO,CM(2)), ( A1 ,CM(3)), ( A2 ,CM(4)), SNIC2730
C      2      ( A3 ,CM(5)), ( A4 ,CM(6))              SNIC2735
C      IF (FX .EQ. 0.) GO TO 5                        SNIC2740
C      ARG1 = (-C*FY**2)/FX                           SNIC2745
C      ARG1 = - ABS(ARG1)                             SNIC2750
C      IF (ABS(ARG1) .GE. 50.) GO TO 5                SNIC2755
C      ARG2 = A1*FX+A2*FX**2 +A3*FY**2 +A4*FY**4      SNIC2760
C                                                    SNIC2765
C      ARG3 = A1+ 2. * A2*FX                          SNIC2770
C      ARG4 = 2.*A3*FY +4.*A4*FY**3                   SNIC2775
C      EX = EXP (ARG1)                                SNIC2780
C      GO TO 10                                       SNIC2785
C      5 FMS = FMO                                     SNIC2790
C      FMXS = 0.                                       SNIC2795
C      FMYS = 0.                                       SNIC2800
C      RETURN                                         SNIC2805
C      10 FMS = FMO +EX* ARG2                          SNIC2810
C      FMXS= EX*((-ARG1/FX)* ARG2 +ARG3)              SNIC2815
C      PAUL= -2.*C*FY/FX                               SNIC2820
C      FMYS= EX*( PAUL*ARG2 + ARG4)                   SNIC2825
C      RETURN                                         SNIC2830
C      END                                           SNIC2835
```

# Contrails

```
$IBFTC SONI    SDD                                SNIC2840
      SUBROUTINE SONK(NM,NCR,YM,FY,FX,IER )        SNIC2845
C                                                    SNIC2850
C   NM = MAX NO OF X,Y ALLOWED. MUST EQUAL DIMENSION OF X,Y, IN MAIN SNIC2855
C   NCR = NO OF X,Y ACTUALLY COMPUTED             SNIC2860
C   YM = MAX. ALLOWABLE VALUE OF Y                SNIC2865
C   FX = X-VALUES                                 SNIC2870
C   FY = Y-VALUES                                 SNIC2875
C   IER = 1 IS NORMAL RETURN                      SNIC2880
C   IER = 2 INDICATES AN ERROR                    SNIC2885
C   CM= MACH. CONSTANTS IN THE EQUATION M=EXP(-CM(1)*Y**2/X)*(CM(3)*X SNIC2890
C   +CM(4)*X*X+CM(5)*Y*Y+CM(6)*Y**4) +CM(2).      SNIC2895
C   THE SUBROUTINE COMPUTES A SET OF X AND Y VALUES ON THE WING WHERE SNIC2900
C   M= 1                                           SNIC2905
C                                                    SNIC2910
C                                                    SNIC2915
C   COMMON                                         SNIC2920
C   */CM/ CH(6)                                    SNIC2925
C   DIMENSION FX(1),FY(1)                          SNIC2930
C   IER =1                                          SNIC2935
C   C=CH(1)                                         SNIC2940
C   FMO=CH(2)                                       SNIC2945
C   A1 =CH(3)                                       SNIC2950
C   A2 =CH(4)                                       SNIC2955
C   A3 =CH(5)                                       SNIC2960
C   A4 =CH(6)                                       SNIC2965
C   FIRST COMPUTE X WHEN Y=0                       SNIC2970
C   ARG = A1**2 -4.*A2*(FMO-1.)                    SNIC2975
C   IF(ARG .GE. 0.0) GO TO 2                        SNIC2980
1  IER = 2                                          SNIC2985
   RETURN                                          SNIC2990
2  FX(1) = (.5/A2)*(-A1+SQRT(ARG))                 SNIC2995
   FY(1) = 0.                                       SNIC3000
   IF(FX(1) .LT. 0.0) GO TO 1                       SNIC3005
   IF(FX(1) .LT. 1.0) GO TO 4                       SNIC3010
   FX(1)=(.5/A2)*(-A1-SQRT(ARG))                   SNIC3015
   IF(FX(1) .LT. 0.0) GO TO 1                       SNIC3020
   IF(FX(1) .GE. 1.0) GO TO 1                       SNIC3025
4  NCR = 2                                          SNIC3030
10  NC1= NCR - 1                                    SNIC3035
   FX(NCR)= FX(NC1)+.01                             SNIC3040
   X= FX(NCR)                                       SNIC3045
   R = C/X                                          SNIC3050
   B = X*(A1+A2*X)                                  SNIC3055
   TO = FY(NC1)**2                                  SNIC3060
   TM1 = A3-R*B                                     SNIC3065
   TM2 = 2.*A4-R*A3                                 SNIC3070
   TM3 = R*A4                                       SNIC3075
   TM4 = 2.*A4+R*(R*B-2.*A3)                       SNIC3080
   TM5 = R*(R*A3-4.*A4)                             SNIC3085
```



# Contrails

|   |          |
|---|----------|
| TM6 = R*R*A4  | SNIC3090 |
| IMAX =1   | SNIC3095 |
| 12 ET =EXP (-R*TO)  | SNIC3100 |
| FT= ET*(B+A3*TO+A4*TO*TO)+FMO-1.                                      | SNIC3105 |
| FPT =ET*(TM1+TM2*TO-TM3*TO**2)  | SNIC3110 |
| FPPT =ET*(TM4+TM5*TO+TM6*TO**2)                                       | SNIC3115 |
| HO = -FT/FPT  | SNIC3120 |
| IF((FT*FPPT) .GE. 0.0) GO TO 14                                       | SNIC3125 |
| HO = .75*HO   | SNIC3130 |
| 14 TO =TO+HO  | SNIC3135 |
| IMAX =IMAX +1   | SNIC3140 |
| 1000 FORMAT(52HO COMPUTATION FOR SONIC LINE WILL NOT CONVERGE, HO = E | SNIC3145 |
| 116.8 )   | SNIC3150 |
| IF(IMAX .LT. 10 ) GO TO 18  | SNIC3155 |
| WRITE (6,1000) HO   | SNIC3160 |
| GO TO 1   | SNIC3165 |
| 18 IF(HO .GT. .0001) GO TO 12   | SNIC3170 |
| FY(NCR) = SQRT(TO)  | SNIC3175 |
| IF(NCR .GE. NM ) GO TO 20   | SNIC3180 |
| IF(FY(NCR) .GE. YM) GO TO 20  | SNIC3185 |
| IF(FX(NCR) .GE. 1.0) GO TO 20   | SNIC3190 |
| NCR = NCR +1  | SNIC3195 |
| GO TO 10  | SNIC3200 |
| 20 RETURN   | SNIC3205 |
| END   | SNIC3210 |

|  |          |
|--|----------|
| \$IBFTC POTE   | SNIC3215 |
| SUBROUTINE POT (NFR,FR,P)                                | SNIC3220 |
| COMMON   | SNIC3225 |
| */XYZ/ SX(101),SXP(101),SY(101),SYP(101),AL(41),TIM(101) | SNIC3230 |
| */CM/ CM(6)  | SNIC3235 |
| */ICNT/ IVAR,NCNT,ISORS,IBR,ITRAP,NMAX                   | SNIC3240 |
| */SOURCE/ XO(20),YO(20)                                  | SNIC3245 |
| */EPS/ E1,E2,FM,YMAX                                     | SNIC3250 |
| */NNN/ NSS,NLCS,NLLS                                     | SNIC3255 |
| C  | SNIC3260 |
| DIMENSION FR(10),P(101,2,10)                             | SNIC3265 |
| CON=-.25/3.14159   | SNIC3270 |
| XS=XO(NSS)   | SNIC3275 |
| YS=YO(NSS)   | SNIC3280 |
| DO 100 N=1,NCNT  | SNIC3285 |
| X=SX(N)  | SNIC3290 |
| Y=SY(N)  | SNIC3295 |
| T = TIM(N)   | SNIC3300 |
| RBAR = SYP(N)  | SNIC3305 |
| 10 DO 30 NF=1,NFR  | SNIC3310 |
| IF (RBAR) 12,14,16                                       | SNIC3315 |
| 12 P(N,1,NF)=0.  | SNIC3320 |
| P(N,2,NF)=0.   | SNIC3325 |
| GO TO 30   | SNIC3330 |
| 14 P(N,1,NF)=UNDEF                                       | SNIC3335 |
| P(N,2,NF)=UNDEF  | SNIC3340 |
| GO TO 30   | SNIC3345 |
| 16 IF (RBAR .LE. 1.E-9) GO TO 14                         | SNIC3350 |
| FACT = CON/RBAR  | SNIC3355 |
| ARG = FR(NF)*T   | SNIC3360 |
| CO = COS(ARG)  | SNIC3365 |
| SI = SIN(ARG)  | SNIC3370 |
| P(N,1,NF) = CO*FACT                                      | SNIC3375 |
| P(N,2,NF) = -SI*FACT                                     | SNIC3380 |
| 30 CONTINUE  | SNIC3385 |
| 100 CONTINUE   | SNIC3390 |
| RETURN   | SNIC3395 |
| END  | SNIC3400 |

# Contrails

```
$IBFTC RKS3*   RUNGE-KUTTA, FORTRAN IV, VERSION 13, SHARE D2*ATFRKS3   SNIC3405
  SUBROUTINE RKS3 (DERIV,CNTRL,Y,DY,ATABL,RTABL,WORK,X,DX,N,IFVD   SNIC3410
  1           ,IBKP,NTRY,IERR)   SNIC3415
    EXTERNAL DERIV,CNTRL   SNIC3420
    INTEGER N,NTRY,IERR   SNIC3425
    LOGICAL IFVD,IBKP   SNIC3430
    REAL Y,DY,ATABL,RTABL,X,DX   SNIC3435
    DIMENSION Y(N),DY(N),ATABL(N),RTABL(N)   SNIC3440
    DIMENSION WORK(1)   SNIC3445
C   DIMENSION WORK(9*N+8)   SNIC3450
    CALL RKINT (DERIV,CNTRL,Y,DY,ATABL,RTABL,WORK(1),WORK(3),WORK(5)   SNIC3455
  1           ,WORK(7),WORK(9),WORK(2*N+9),WORK(4*N+9),WORK(6*N+9)   SNIC3460
  2           ,WORK(7*N+9),WORK(8*N+9),X,DX,N,IFVD,IBKP,NTRY,IERR)   SNIC3465
    RETURN   SNIC3470
    END   SNIC3475
$IBFTC RKINT*   CALLED BY RKS3, RUNGE-KUTTA, F 4, V13, SHARE D2*ATFRKS3   SNIC3480
  SUBROUTINE RKINT (DERIV,CNTRL,REALY,DY,ATABL,RTABL,DELTAX,X,XHALF   SNIC3485
  1           ,XZERO,Y,YHALF,YZERO,DYHALF,DYZERO,DELTAY,REALX   SNIC3490
  2           ,DX,N,IFVD,IBKP,NTRY,IERR)   SNIC3495
    EXTERNAL DERIV,CNTRL   SNIC3500
    INTEGER N,NTRY,IERR   SNIC3505
    LOGICAL IFVD,IBKP   SNIC3510
    REAL REALY,DY,ATABL,RTABL,DELTAX,DYHALF,DYZERO,DELTAY,REALX,DX   SNIC3515
    DOUBLE PRECISION X,XHALF,XZERO,Y,YHALF,YZERO   SNIC3520
    DIMENSION REALY(N),DY(N),ATABL(N),RTABL(N),Y(N),YHALF(N),YZERO(N)   SNIC3525
  1           ,DYHALF(N),DYZERO(N),DELTAY(N)   SNIC3530
    IERR = 0   SNIC3535
  10 DELTAX = DX   SNIC3540
    X = REALX   SNIC3545
    DO 20 I=1,N   SNIC3550
  20 Y(I) = REALY(I)   SNIC3555
    CALL DERIV   SNIC3560
    GO TO 200   SNIC3565
  30 IF (DX .EQ. 0.) GO TO 230   SNIC3570
    DELTAX = DX   SNIC3575
    DX2 = DX/2.   SNIC3580
    DX4 = DX/4.   SNIC3585
    ZXERO = X   SNIC3590
    DO 40 I=1,N   SNIC3595
  40 DYZERO(I) = Y(I)   SNIC3600
    DO 110 J=1,2   SNIC3605
  110 XHALF = X   SNIC3610
    X = X+DX4   SNIC3615
    REALX = X   SNIC3620
    DO 50 I=1,N   SNIC3625
  50 DELTAY(I) = DY(I)*DX4   SNIC3630
    YHALF(I) = Y(I)   SNIC3635
    Y(I) = Y(I)+DELTAY(I)   SNIC3640
  50 REALY(I) = Y(I)   SNIC3645
```

# Contrails

```
CALL DERIV
DO 60 I=1,N
DELTAY(I) = DELTAY(I)+DY(I)*DX2
Y(I) = YHALF(I)+DY(I)*DX4
60 REALY(I) = Y(I)
CALL DERIV
X = XHALF+DX2
REALX = X
DO 70 I=1,N
DELTAY(I) = DELTAY(I)+DY(I)*DX2
Y(I) = YHALF(I)+DY(I)*DX2
70 REALY(I) = Y(I)
CALL DERIV
DO 80 I=1,N
DELTAY(I) = (DELTAY(I)+DY(I)*DX4)/3.
Y(I) = YHALF(I)+DELTAY(I)
80 REALY(I) = Y(I)
CALL DERIV
GO TO (90,110),J
90 DO 100 I=1,N
100 DYHALF(I) = DY(I)
110 CONTINUE
IF (IFVD) GO TO 200
ERRMAX = 0
DO 120 I=1,N
ERR = ATABL(I)+ABS(RTABL(I)*REALY(I))
IF (ERR .EQ. 0.) GO TO 220
SR = (YZERO(I)+4.*DYHALF(I)+DY(I))/3.*DX2
120 ERRMAX = AMAX1(ERRMAX,ABS (SR-(REALY(I)-SNGL(YZERO(I))))/ERR)
IF (ERRMAX-1.) 130,170,160
130 IF (ERRMAX-.75) 140,200,170
140 IF (ERRMAX-.075) 150,200,200
150 DX = DX*1.5845932
GO TO 200
160 DX = DX/1.5848932
IF (.NOT. IEKP) GO TO 180
ERRMAX = ERRMAX/10.
IF (ERRMAX .GT. 1.) GO TO 160
GO TO 180
170 DX = DX/1.5848932
GO TO 200
180 X = XZERO
DO 190 I=1,N
Y(I) = YZERO(I)
190 DY(I) = DYZERO(I)
GO TO 30
200 NTRY = 1
CALL CNTRL (NTRY)
GO TO (30,210,180,10),NTRY
210 RETURN
```

SNIC3655  
SNIC3660  
SNIC3665  
SNIC3670  
SNIC3675  
SNIC3680  
SNIC3685  
SNIC3690  
SNIC3695  
SNIC3700  
SNIC3705  
SNIC3710  
SNIC3715  
SNIC3720  
SNIC3725  
SNIC3730  
SNIC3735  
SNIC3740  
SNIC3745  
SNIC3750  
SNIC3755  
SNIC3760  
SNIC3765  
SNIC3770  
SNIC3775  
SNIC3780  
SNIC3785  
SNIC3790  
SNIC3795  
SNIC3800  
SNIC3805  
SNIC3810  
SNIC3815  
SNIC3820  
SNIC3825  
SNIC3830  
SNIC3835  
SNIC3840  
SNIC3845  
SNIC3850  
SNIC3855  
SNIC3860  
SNIC3865  
SNIC3870  
SNIC3875  
SNIC3880  
SNIC3885  
SNIC3890  
SNIC3895  
SNIC3900

# Contrails

220 IERR = 1  
RETURN  
230 IERR = -1  
RETURN  
END

SNIC3905  
SNIC3910  
SNIC3915  
SNIC3920  
SNIC3925

# FORTRAN FIXED 10 DIGIT DECIMAL DATA

APPENDIX II. *Control* Sample Input and Output

| DECK NO. | PROGRAMMER | DATE | PAGE | of | JOB NO. | DESCRIPTION                                      | DO NOT KEY PUNCH |
|----------|------------|------|------|----|---------|--|------------------|
| 1        |            |      |      |    |         | NSOURCE (NUMBER OF SOURCE POINTS, 20 MAXIMUM)    |                  |
| 13       |            |      |      |    |         | NLA (NUMBER OF A PER SOURCE, 40 MAXIMUM)         |                  |
| 25       |            |      |      |    |         | NPL (NUMBER OF PLATFORM COORDINATES, 8 MAXIMUM)  |                  |
| 37       |            |      |      |    |         | NMAX (LIMIT NUMBER OF POINTS PER PLOT, 100 MAX)  |                  |
| 49       |            |      |      |    |         | NF (NUMBER OF ASSUMED FREQUENCIES, 10 MAXIMUM)   |                  |
| 61       |            |      |      |    |         |  |                  |
| 1        |            |      |      |    |         | IFVD LOGICAL WORDS - VARIABLE INTERVAL MODE      |                  |
| 13       |            |      |      |    |         | IBKP IF IFVD = FALSE AND IBKP = TRUE THIS CHOICE |                  |
| 25       |            |      |      |    |         | IS RECOMMENDED. FIXED INTERVAL IF                |                  |
| 37       |            |      |      |    |         | IFVD = TRUE.                                     |                  |
| 49       |            |      |      |    |         |  |                  |
| 61       |            |      |      |    |         |  |                  |
| 1        |            |      |      |    |         |  |                  |
| 13       |            |      |      |    |         |  |                  |
| 25       |            |      |      |    |         |  |                  |
| 37       |            |      |      |    |         |  |                  |
| 49       |            |      |      |    |         |  |                  |
| 61       |            |      |      |    |         |  |                  |

FORTRAN FIXED 10 DIGIT DECIMAL DATA

| DECK NO. _____ |                | PROGRAMMER _____ | DATE _____       | PAGE _____ | of _____ | JOB NO. _____ |
|----------------|----------------|------------------|------------------|------------|----------|---------------|
| NUMBER         | IDENTIFICATION | DESCRIPTION      | DO NOT KEY PUNCH |            |          |               |
| 1              |                |                  |                  |            |          |               |
| 13             |                |                  |                  |            |          |               |
| 25             |                |                  |                  |            |          |               |
| 37             |                |                  |                  |            |          |               |
| 49             |                |                  |                  |            |          |               |
| 61             |                |                  |                  |            |          |               |
| 1              |                |                  |                  |            |          |               |
| 13             |                |                  |                  |            |          |               |
| 25             |                |                  |                  |            |          |               |
| 37             |                |                  |                  |            |          |               |
| 49             |                |                  |                  |            |          |               |
| 61             |                |                  |                  |            |          |               |
| 1              |                |                  |                  |            |          |               |
| 13             |                |                  |                  |            |          |               |
| 25             |                |                  |                  |            |          |               |
| 37             |                |                  |                  |            |          |               |
| 49             |                |                  |                  |            |          |               |
| 61             |                |                  |                  |            |          |               |

XO(1) COORDINATES OF SOURCE POINTS.  
 YO(1) (ALL GEOMETRY IS NORMALIZED ON b<sub>0</sub>, THE  
 XO(2) DISTANCE FROM MOST FORWARD TO MOST AFT  
 YO(2) PORTION OF THE WING.)

XO(NSØRCE)  
 YO(NSØRCE)

# FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. \_\_\_\_\_ PROGRAMMER \_\_\_\_\_ DATE \_\_\_\_\_ PAGE \_\_\_\_\_ of \_\_\_\_\_ JOB NO. \_\_\_\_\_

| NUMBER | IDENTIFICATION | DESCRIPTION DO NOT KEY PUNCH                           |
|--------|----------------|--|
| 1      |                | CM(1) COEFFICIENTS IN MACH NUMBER EQUATION             |
| 13     |                | CM(2) $M = CM(2) + EXP \{-CM(1)(Y^2/X)\}$ .            |
| 25     |                | $\{CM(3)X + CM(4)X^2 + CM(5)Y^2$                       |
| 37     |                | $+ CM(6)Y^4\}$   |
| 49     |                | THE COEFFICIENTS WERE DETERMINED BY A                  |
| 61     |                | CM(6) LEAST-SQUARE PROCEDURE.                          |
| 1      |                | DZ INITIAL VALUE OF INCREMENT                          |
| 13     |                | E1 TEST WORD FOR TRAPPING SIGNAL ON LOCAL M.L.         |
| 25     |                | E2 TEST WORD FOR STOPPING TRAPPED SIGNAL ON SONIC LINE |
| 37     |                | YMAX NORMALIZED SEMI-SPAN                              |
| 49     |                |  |
| 61     |                |  |
| 1      |                |  |
| 13     |                |  |
| 25     |                |  |
| 37     |                |  |
| 49     |                |  |
| 61     |                |  |



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DECK NO. \_\_\_\_\_ PROGRAMMER \_\_\_\_\_ DATE \_\_\_\_\_ PAGE \_\_\_\_\_ of \_\_\_\_\_ JOB NO. \_\_\_\_\_

| NUMBER | IDENTIFICATION    | DESCRIPTION DO NOT KEY PUNCH                       |
|--------|-------------------|--|
| 1      | - 0 4             | ATABLE(1) ATABLE AND RTABLE DETERMINE THE ACCURACY |
| 13     | - 0 4             | ATABLE(2) REQUIREMENTS FOR DECREASING OR IN-       |
| 25     | - 0 4             | ATABLE(3) CREATING THE INTERVAL IN THE VARIABLE    |
| 37     | - 0 4             | ATABLE(4) INTERVAL MODE.                           |
| 49     | - 0 3<br>73       | RTABLE(1) 90                                       |
| 61     | - 0 3<br>7        | RTABLE(2)  |
| 1      | - 0 3             | RTABLE(3)  |
| 13     | - 0 3             | RTABLE(4)  |
| 25     |                   |  |
| 37     |                   |  |
| 49     | 73                |  |
| 61     | 8                 |  |
| 1      | + 0 1             | PLX(1) PLANFORM COORDINATES. LIST ALL CORNERS      |
| 13     | - 4 6 6 3 1 + 0 0 | PLY(1) STARTING FROM LEFT ALONG LEADING EDGE       |
| 25     | + 0 0             | PLX(2) TOWARD THE RIGHT, AND AGAIN STARTING        |
| 37     | + 0 0             | PLY(2) AT LEFT, ALONG TRAILING EDGE TOWARD         |
| 49     | + 0 1<br>73       | PLX(3) THE RIGHT. ALL GEOMETRY IS NORMALIZED       |
| 61     | 4 6 6 3 1 + 0 0   | PLY(3) on b <sub>0</sub> .                         |
| 1      |                   |  |

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| NUMBER | IDENTIFICATION | DESCRIPTION DO NOT KEY PUNCH              |
|--------|----------------|---|
| 1      |                | PLX(4)                                    |
| 13     |                | PLY(4)                                    |
| 25     |                | PLX(5)                                    |
| 37     |                | PLY(5)                                    |
| 49     | 73             |   |
| 61     | 1 0            |   |
| 1      |                | CINF REMOTE SPD. OF SND. IN 10 UNITS/SEC. |
| 13     |                | FMINF REMOTE MACH NUMBER                  |
| 25     |                | TAU MAXIMUM THICKNESS RATIO (T/C)         |
| 37     |                | TSAH TANGENT SEMI-APEX ANGLE              |
| 49     | 73             |   |
| 61     | 1 1            |   |
| 1      |                | FREQ(1) FREQUENCIES, FOR COMPUTATION      |
| 13     |                | FREQ(2) OF POTENTIALS, (RADIAN/SEC)       |
| 25     |                | FREQ(NF)                                  |
| 37     |                |   |
| 49     | 73             |   |
| 61     | 1 2            |   |

86

IVAR NCNT ISORS IBR ITRAP NLCS =

2 12 1 1 0 9

XO = 0.18000000E 00  
 YO = 0.00000000E-38

MACH NO. = 0.97311907E 00

ACOUSTIC RAY PATH FOR LAMBDA = 0.57119817E 00

| X              | Y              | X-PRIME        | R-BAR          | TIME           |
|----------------|----------------|----------------|----------------|----------------|
| 0.1800000E 00  | 0.0000000E-38  | 0.21141810E 00 | 0.0000000E-38  | 0.0000000E-38  |
| 0.18424259E 00 | 0.2000000E-01  | 0.21257251E 00 | 0.72580551E-02 | 0.38412277E-02 |
| 0.19096695E 00 | 0.51691864E-01 | 0.21112035E 00 | 0.18791161E-01 | 0.99172437E-02 |
| 0.20135082E 00 | 0.10191058E 00 | 0.20083363E 00 | 0.37255190E-01 | 0.19553514E-01 |
| 0.20530157E 00 | 0.12189922E 00 | 0.19445186E 00 | 0.44702658E-01 | 0.23404072E-01 |
| 0.20915295E 00 | 0.14188787E 00 | 0.19213422E 00 | 0.52104825E-01 | 0.27296191E-01 |
| 0.21302884E 00 | 0.16187273E 00 | 0.19647354E 00 | 0.59213898E-01 | 0.31281433E-01 |
| 0.21700866E 00 | 0.18185759E 00 | 0.20163166E 00 | 0.66174245E-01 | 0.35289611E-01 |
| 0.22350139E 00 | 0.21352546E 00 | 0.20816299E 00 | 0.77123095E-01 | 0.41616165E-01 |
| 0.23415963E 00 | 0.26370616E 00 | 0.21631019E 00 | 0.94319344E-01 | 0.51590675E-01 |
| 0.25178554E 00 | 0.34322216E 00 | 0.22669061E 00 | 0.12128799E 00 | 0.67289645E-01 |
| 0.28126384E 00 | 0.46922266E 00 | 0.24104052E 00 | 0.16353054E 00 | 0.91917480E-01 |

VELOCITY POTENTIALS ALONG A RAY PATH FOR A SOURCE AT

XO = 0.18000000E 00  
 YO = 0.00000000E-38  
 LAMBDA = 0.57119817E 00

ALTERNATING REAL AND IMAGINARY

OMEGA = 0.10000000E 02

|               |              |               |              |               |              |               |              |
|---------------|--------------|---------------|--------------|---------------|--------------|---------------|--------------|
| .....         | .....        | .....         | .....        | .....         | .....        | .....         | .....        |
| -0.209531E 01 | 0.415009E 00 | -0.109559E 02 | 0.421050E 00 | -0.421403E 01 | 0.419291E 00 | -0.147071E 01 | 0.411726E 00 |
| -0.127868E 01 | 0.413568E 00 | -0.173162E 01 | 0.412835E 00 | -0.943756E 00 | 0.417118E 00 | -0.295125E 00 | 0.386914E 00 |
| -0.733892E 00 | 0.416219E 00 | -0.112844E 01 | 0.415620E 00 | -0.408919E 00 | 0.408919E 00 | 0.128649E 00  | 0.469308E 00 |

OMEGA = 0.20000000E 02

|               |              |               |              |               |              |               |              |
|---------------|--------------|---------------|--------------|---------------|--------------|---------------|--------------|
| .....         | .....        | .....         | .....        | .....         | .....        | .....         | .....        |
| -0.197475E 01 | 0.814201E 00 | -0.109317E 02 | 0.841478E 00 | -0.415181E 01 | 0.834462E 00 | -0.130527E 01 | 0.792965E 00 |
| -0.108936E 01 | 0.786997E 00 | -0.158867E 01 | 0.803160E 00 | -0.694583E 00 | 0.763032E 00 | 0.128649E 00  | 0.469308E 00 |
| -0.433042E 00 | 0.724092E 00 | -0.915255E 00 | 0.780015E 00 | 0.639566E 00  | 0.639566E 00 | 0.128649E 00  | 0.469308E 00 |

OMEGA = 0.00000000E-38

|               |              |               |              |               |              |               |              |
|---------------|--------------|---------------|--------------|---------------|--------------|---------------|--------------|
| .....         | .....        | .....         | .....        | .....         | .....        | .....         | .....        |
| -0.213601E 01 | 0.000000E-38 | -0.109640E 02 | 0.000000E-38 | -0.423484E 01 | 0.000000E-38 | -0.152726E 01 | 0.000000E-38 |
| -0.134390E 01 | 0.000000E-38 | -0.178015E 01 | 0.000000E-38 | -0.103182E 01 | 0.000000E-38 | -0.486622E 00 | 0.000000E-38 |
| -0.843703E 00 | 0.000000E-38 | -0.120255E 01 | 0.000000E-38 | 0.000000E-38  | 0.000000E-38 | 0.000000E-38  | 0.000000E-38 |

## APPENDIX III. Application to the Boundary Value Problem

A procedure that may be used to match the tangential flow condition on a wing surface is, in principle, the same as that employed by Rodemich in the box method for uniform sonic flow (Reference 3). The velocity potential at a field point  $(x, y, z)$  due to a doublet sheet in its zone of influence, is

$$\bar{\phi}(x, y, z) = \frac{\partial}{\partial z} \int \int_{S+W} \Delta\phi(\xi, \eta) \phi_0(x-\xi, y-\eta, z) d\xi d\eta \quad (39)$$

where  $\Delta\phi(\xi, \eta)$  is the velocity potential discontinuity through the doublet sheet over the region  $S + W$  (the surface and its wake), and

$$\phi_0(x-\xi, y-\eta, z) = \frac{-1}{2\pi\bar{R}} \sum_{n=1}^N e^{-iag_n} \quad (40)$$

where  $\bar{R} = \sqrt{(x-\xi)^2 + [1-M_L^2(x, y, z)] [(y-\eta)^2 + z^2]}$

and where  $N$  represents the number of times the wave front passes the field point. In uniform subsonic flow  $N$  equals one, in uniform supersonic flow it equals two, and in the limiting case of uniform sonic flow it equals one. As discussed previously, in uniform sonic flow the stationary portion of the perturbation wave front is not augmented by high frequency signals that follow it; instead, the pressure discontinuity is dissipated by them.

When the local flow in a non-uniform flow field is sonic the wave front gradually becomes stationary and is dissipated. Rays of this type are shown in Figures 9, 12, and 13. In certain regions of non-uniform flow a wave front may pass field points more than twice as shown in Figures 6, 7, 9, 10, and 12. These regions may be in the region of subsonic flow or in supersonic flow. Multiple crossings normally occur on receding portions of the wave front. Ray lines on advancing portions normally pass over the trailing edge before they cross. In these regions of multiple crossings of the wave front, care must be taken to establish an accurate value of  $N$ , and of each of the corresponding  $g_n$ 's,  $n = 1, 2, \dots, N$ . A computer program that may be used to do this is contained herein. Figures 11 and 13 show that in some regions of both subsonic and supersonic flow even the receding ray lines do not cross. All of Figures 6 through 13 show that once a ray crosses the transition region at the edge of the planform it does not return to the wing region. This characteristic is important because when a doublet solution is employed a ray trace can be ignored once it reaches an edge that is not adjacent to the wake.

# Contrails

The next step in the procedure is to define a grid of square boxes over the region  $S + W$ , and assume that  $\Delta\phi(\xi, \eta)$  is constant over the area of each box. For this to be a valid assumption as many as 50 boxes along the root chord may be required. The upwash adjacent to the upper surface may be written

$$\bar{w}(x, y, 0+) = \lim_{z \rightarrow 0+} \frac{\phi(x, y, z)}{z}$$

or,

$$\bar{w}(x_1, y_1, 0+) = \sum_{i, j} \Delta\phi_{i, j} \iint_{B_{i, j}} \psi(x_1 - \xi, y_1 - \eta) d\xi d\eta \quad (41)$$

i.e., the upwash at  $(x_1, y_1)$  equals the summation (over all boxes  $B_{i, j}$  that influence it), of products of the constant velocity potential discontinuities and their downwash influence coefficients. The latter are represented by the double integral of the kernel  $\psi$  over the areas of the boxes. The limits of integration and  $\Delta\phi$  of Equation (39) are not functions of  $z$ , so from Equation (40) we get

$$\psi(x_1 - \xi, y_1 - \eta) = \frac{-1}{2\pi} \lim_{z \rightarrow 0+} \frac{1}{z} \frac{\partial}{\partial z} \frac{\sum e^{-i\omega g_n}}{\bar{R}} \quad (42)$$

At this point it is theorized that for non-uniform flow around a nearly planar surface the variation in signal transmission time with distance normal to the surface is approximately equal to the variation in uniform flow, i.e.,

$$\frac{\partial g_n}{\partial z} = \frac{\partial}{\partial z} \frac{M(x - \xi) \mp \bar{R}}{C(M^2 - 1)}$$

or, performing the differentiation

$$\frac{\partial g_n}{\partial z} = \frac{\pm z}{\bar{C}\bar{R}} \quad (43)$$

where the upper sign refers to the advancing portion of the wave front and the lower sign to the receding portion.  $C$  is the speed of sound. Making use of equation (43) when taking the derivative in equation (42),

$$\psi(x_1 - \xi, y_1 - \eta) = \frac{-1}{2\pi} \frac{\beta^2 C \pm i\omega \bar{R}}{\bar{C}\bar{R}^3} \sum_n e^{-i\omega g_n} \quad (44)$$

The  $g_n$ 's are those obtained by tracing ray paths through the non-uniform flow field.

# Contrails

One way in which Equation (44) may be evaluated and integrated is as follows: Say for nine values of  $(\xi, \eta)$  on each sending box, the values of the kernel at the center of the receiving box  $(x_i, y_j)$  are evaluated. Since the ray paths are not known in advance, each of these values must be interpolated from values in its neighborhood. It is then necessary to evaluate the integral in Equation (41) given the values of the integrand at nine points in the region of integration.

The unknowns in Equation (41) are the  $\overline{\Delta\phi}_{i,j}$ 's. When the center of a receiving box  $(x_i, y_j)$  lies in the subsonic flow region it lies in the zone of influence of every other point in the subsonic region and may lie in the zone of influence of a small portion of the supersonic region (Figure 9). All velocity potentials in zones of mutual influence must be determined simultaneously. Once velocity potentials have been established that meet the tangential flow conditions on the surface and the zero pressure difference condition on the wake they may be fitted with analytical expressions that have the proper edge behavior. Using these expressions, local oscillatory pressures and generalized forces may be obtained in the way outlined in Reference 3.

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|-----|--|--------|----|--------|----|--------|----|
|     |  | ROLE   | WT | ROLE   | WT | ROLE   | WT |
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