

## FOREWORD

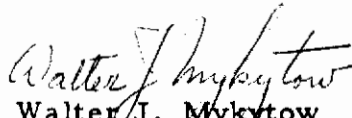
This report covers the research conducted from 1 April 1963 to 15 November 1964 by the Space and Information Systems Division of North American Aviation, Inc., Downey, California, for the Aerospace Dynamics Branch, Vehicle Dynamics Division, Air Force Flight Dynamics Laboratory, Wright-Patterson Air Force Base, Ohio, under Contract No. AF 33(657)-10399.

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Mr. L. V. Andrew was the Program Manager for North American Aviation. Mr. M. T. Moore formulated the approach based on an outline provided by Professor H. Ashley who acted as consultant during the development of this report and on previous work by L. V. Andrew. Mr. Moore also directed and contributed to the writing of the program for the IBM 7094 computer. He was assisted by Mr. J. B. O'Neill. Dr. E. R. Rodemich contributed many enlightening discussions.

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This report has been reviewed and approved.

  
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## ABSTRACT

In this volume, Ashley's approach to mutual interference theory by source superposition methods has been applied to the prediction of supersonic air loads on intersecting thin lifting surfaces in steady or oscillatory motion. Steady loading is regarded as the special case of zero frequency of oscillation. Each surface may be oscillating in a mode of rigid or elastic vibration or linear combinations thereof. Evvard's diaphragm concept has been extended to treat the out-of-plane interference problem. As a result, any leading or side edge on any of the intersecting surfaces may be subsonic. The study reported herein has led to the formulation of a method by which diaphragm regions can be selected that eliminate the need for calculating out-of-plane velocity potentials. Based on mutual interference theory, the method requires only the calculation of out-of-plane velocities so that tangential flow conditions may be met.

The Mach-box method has been used to obtain an approximate solution to the problem. In following the aerodynamic influence coefficient procedure of Zartarian and Hsu, each surface and diaphragm is overlaid with a grid of rectangular Mach boxes, the diagonals of which are parallel to the Mach lines. For purposes of calculating induced velocities and velocity potentials, the source strength over the area of each box is assumed to be constant. Out-of-plane velocity influence coefficients are calculated for the center of each box in the interference region, and the tangential flow condition is satisfied there. For purposes of calculating generalized forces, the resulting velocity potential over the area of each box is also assumed to be constant and equal to the value at the center of the box.

A computer program that calculates a matrix of generalized forces for the special case of a wing with symmetrically folded tips is also presented. The generalized forces are those due to motion of the surface which may consist of a linear combination of as many as ten modes of motion. By proper specification of the modes and frequencies, the user may obtain lift coefficients, longitudinal stability derivatives, loads due to symmetric sinusoidal gusts, drag due to surface warpage, or an array of generalized forces which may be used to calculate flutter speeds. If the surface is free from flutter, the same arrays may be used in a Fourier series approach to calculate the responses to random or discrete gusts.

Favorable comparisons of the results have been obtained for delta wings with folded tips and without folded tips. Results for rectangular and delta wings in both steady and oscillatory motion are presented and compared.

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favorably with theoretical forces calculated from both exact and other approximate numerical methods. A smooth decrease in the steady state lift and moment on a delta wing is calculated when the tips are folded down to the vertical position. These results agree with the trends that have been observed in recent experiments.

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## SYMBOLS

<u>Symbol</u>	<u>Definition</u>
$a_{\infty}$	Free-stream speed of sound
$b$	Box length
$b/\beta$	Box width
$C_p$	Surface pressure coefficient
$\frac{D}{Dt}$	Linearized substantial derivative
$f_i$	Nondimensional displacement mode of lifting surface
$G$	Strength of unit source
$H$	Source strength per unit area
$k$	Reduced frequency based on box length, $\omega b/U_{\infty}$
$\bar{k}$	Modified reduced frequency, $k M_{\infty}^2 / \beta^2$
$k_r$	Reduced frequency based on root chord, $\omega c/U_{\infty}$
$M_{\infty}$	Free-stream Mach number, $U_{\infty}/a_{\infty}$
$N$	Normal velocity influence coefficient
$n, m, l$	Location of receiving box center in $x_1, y_1, z_1$ coordinate system
$p$	Local pressure at the airfoil surface
$p_{\infty}$	Free-stream pressure
$Q_i$	Generalized force
$\vec{q}$	Unsteady perturbation velocity
$R$	$\sqrt{(x-\xi)^2 + (y-\eta)^2 + (z-\zeta)^2}$
$R_1$	$\sqrt{\xi_1^2 + \eta_1^2 + \zeta_1^2}$

List of Symbols continued on next page.

<u>Symbol</u>	<u>Definition</u>
$S$	Reference area
$U_{\infty}$	Free-stream velocity
$U, L$	Superscripts referring to upper or lower source sheet, respectively
$u, v, w$	Components of $\vec{q}$ in the $x, y, z$ directions, respectively
$\bar{u}, \bar{v}, \bar{w}$	Time-independent factors of $u, v, w$
$V$	Horizontal velocity influence coefficient
$W$	Vertical velocity influence coefficient
$\left. \begin{matrix} X, Y, Z \\ \xi, \eta, \zeta \end{matrix} \right\}$	Cartesian coordinate system (applicable in Section 2 only)
$\left. \begin{matrix} x, y, z \\ \xi, \eta, \zeta \end{matrix} \right\}$	Cartesian coordinate system transformed to $M = \sqrt{2}$
$\left. \begin{matrix} x_1, y_1, z_1 \\ \xi_1, \eta_1, \zeta_1 \end{matrix} \right\}$	Cartesian coordinate system transformed to $M = \sqrt{2}$ and nondimensionalized by box length
$\bar{z}$	Dimensional deflection mode of lifting surface
$\beta$	$\sqrt{M_{\infty}^2 - 1}$
$\delta C_w$	Incremental virtual work coefficient
$\delta r$	Nondimensional virtual displacement
$\delta q_i$	Virtual displacement in the generalized coordinate, $q_i$
$\rho_{\infty}$	Free-stream density
$\omega$	Oscillatory frequency - $2\pi f$
$\phi$	Unsteady perturbation velocity potential

List of Symbols continued on next page.



SymbolDefinition

$\bar{\phi}$	Time-independent factor of $\phi$
$\Phi$	Velocity potential influence coefficient
$v, u, l$	Location of sending box center in $x_1, y_1, z_1$ coordinate system
$\bar{v}, \bar{u}, l$	Fixed values of $\bar{\xi}, \bar{\eta}, \bar{\zeta}$ respectively
$\bar{\xi}_1, \bar{\eta}_1, \bar{\zeta}_1$	Relative position coordinates

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## 1. INTRODUCTION

Garrick and Rubinow (Reference 1) have shown that every solution of the boundary-value problem describing small disturbances in a parallel supersonic flow can be built up by superposing elementary solutions. One of these elementary solutions is that of a pulsating sound source fixed in the flow; however, exact numerical evaluations of the integral equations resulting from the source superposition method are laborious, if not impossible, to obtain.

Pines, Dugundji, and Neuringer (Reference 2) published the first source superposition method to successfully approximate the aerodynamic forces on an oscillating thin planar surface in supersonic flow. They employed Evvard's diaphragm concept (Reference 3) to handle subsonic leading edges and overlaid the surface and diaphragm with a grid of square boxes. For purposes of calculating pressures, they assumed that the source strength over the area of each box is a constant value which satisfies the condition of tangential flow at the center of the box. Thus, they established a method for calculating aerodynamic influence coefficients similar to structural stiffness influence coefficients.

Li (Reference 4) then published the Mach-box formulation, that greatly simplified calculation of the aerodynamic influence coefficients. Li's formulation was fully developed for planar surfaces by Zartarian and Hsu (Reference 5). The Mach-box procedure is basically the same as Pines' method, differing only in that the surface and diaphragm is overlaid with a grid of rectangular boxes the diagonals of which are parallel to the Mach lines.

Zartarian and Hsu were the first to establish the minimum number of boxes on a swept, low-aspect-ratio lifting surface that can be expected to yield reasonable results when the leading edge of the surface is represented by the jagged pattern of box leading edges. They were also the first to use velocity potential influence coefficients (VPIC's). Because the velocity potential is a better behaved function, near the leading edge, than the pressure, the VPIC's yield more accurate generalized forces than pressure influence coefficients for the same grid. The Mach-box approach, using VPIC's, evolved as the most efficient way to calculate generalized forces on planar lifting surfaces in supersonic flow.

The Mach-box approach to the source superposition method is currently accepted as a standard procedure for determining the generalized

forces on planar lifting surfaces in either steady or oscillatory motion at supersonic speeds. Ashley (Reference 6) has shown how the source superposition method may be applied to two or more planar surfaces flying in the disturbance field of each other. He also outlined application of the method to two intersecting planar surfaces when the line of intersection is parallel to the free-stream. The development contained in this report is the outgrowth of several iterations on the procedure and several consultations with Professor Ashley.

The extension of the source superposition method to intersecting lifting surfaces creates some difficulty in understanding the physics as well as the mathematics of the problem. This is primarily due to the artificiality of the method of using sheets of symmetric disturbances to obtain antisymmetric air loads.

An attempt was made to reduce the diaphragm area to a minimum in the expectation that the number of numerical operations would thereby be reduced to a minimum. Another desirable property of this arrangement of disturbance sheets is that the flow region is separated into a single upper and a single lower region. Further study revealed that this arrangement does not lead to a minimum number of operations because of the more complicated geometry.

A simpler geometric arrangement can be achieved by constructing the Mach envelopes of each of the separate components of the lifting surface and then removing redundant diaphragm regions. By this means the flow region is separated into a minimum of three regions. Removal of the redundant diaphragms results in the necessity to define separate sending regions for upper and lower parts of each sheet of disturbances in the interference region. The computer program contained in this report is based on this geometric arrangement.

The Mach-box method is used in the computer program to obtain aerodynamic forces on a wing with non-coplanar tips. In the computer program, discussed in Section 4, the velocity potential distribution over all surfaces is calculated and then integrated to obtain generalized forces. These quantities may be used to obtain (1) steady-state air loads on each surface for any small angle of attack and camber distribution, (2) stability derivatives for 20 different frequencies of rigid pitch or plunge motion, or (3) matrices of generalized forces for up to 10 modes of rigid or elastic vibration. For each mode, the distribution of velocity potentials, containing all the interference effects, may then be smoothed by fitting to the point values a polynomial that has the proper edge conditions. Operating on the resulting polynomial with the linearized substantial derivative provides an expression for the distribution of lifting pressure over the entire surface.

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After the computer program was written it was discovered that a more elegant approach is one in which the redundant diaphragms are retained. In this approach (1) the definition of sending surfaces is general and simple, (2) satisfaction of tangential flow conditions on each of the surfaces and the continuity conditions on diaphragms is a straightforward sequential process, (3) the necessity for calculating out-of-plane velocity potentials disappears (thus reducing the number of numerical operations), and (4) insight is provided for further extension. A discussion of this approach is contained in Section 6 of this report.

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## 2. THIN AIRFOIL THEORY IN SUPERSONIC FLOW

### LINEARIZED EQUATIONS

Thin airfoil or small perturbation theory is applied to describe the flow patterns that result when small disturbances are superposed on parallel uniform flow. The small perturbation methods used to linearize the second-order partial differential equation that characterizes compressible fluid flow are completely described for steady flow by Sears (Reference 7) and for unsteady flow by Ashley (Reference 8). These methods are summarized in Part I of this report and presented in a form consistent with the following application to a slightly perturbed, uniform, parallel, supersonic flow.

Consider such a perturbed supersonic flow in the direction of the positive X-axis of a Cartesian coordinate system (Figure 1). The linearized partial differential equation describing this flow may be expressed in the form

$$\phi_{XX} + \phi_{YY} + \phi_{ZZ} = \frac{1}{a_{\infty}^2} \left( U_{\infty}^2 \phi_{XX} + 2U_{\infty} \phi_{XT} + \phi_{TT} \right) \quad (1)$$

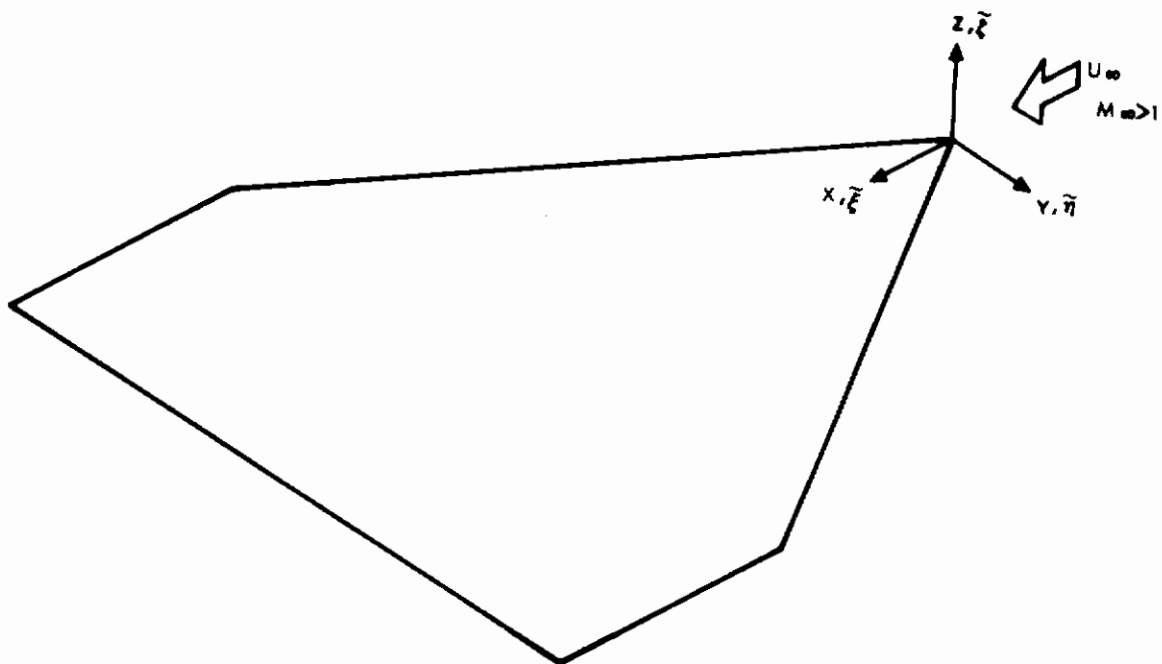
where  $\phi$  is the perturbation velocity potential,  $a_{\infty}$  is the fixed acoustic speed far upstream, and  $U_{\infty}$  is the uniform flow velocity. By collecting terms and introducing the free-stream Mach number

$$M_{\infty} = U_{\infty}/a_{\infty}$$

Equation 1 can be cast in the form of a hyperbolic differential equation

$$\left( M_{\infty}^2 - 1 \right) \phi_{XX} - \phi_{YY} - \phi_{ZZ} = -\frac{1}{a_{\infty}^2} \left( 2U_{\infty} \phi_{XT} + \phi_{TT} \right) \quad (2)$$

which, for the supersonic case, is satisfied only within a characteristic region. This region is the downstream Mach cone; its semivertex angle is equal to the Mach angle and its axis is parallel to the X-axis.



**Figure 1. Lifting Surface in  $Z = 0$  Plane in Parallel Supersonic Flow**

The small disturbance (Figure 2) has a velocity given by the expression

$$\vec{q} = u\vec{i} + v\vec{j} + w\vec{k} \quad (3)$$

where,

$$u = \frac{\partial \phi}{\partial X}, \quad v = \frac{\partial \phi}{\partial Y}, \quad w = \frac{\partial \phi}{\partial Z}$$

and is placed in a uniform parallel flow at the point  $(\tilde{\xi}, \tilde{\eta}, \tilde{\zeta})$ . The signals from the disturbance can only be received at the points  $(X, Y, Z)$  within the downstream Mach cone which has its vertex at the sending point  $(\tilde{\xi}, \tilde{\eta}, \tilde{\zeta})$ . Conversely, a given receiving point feels only those signals transmitted by sending points in the upstream Mach cone which has its vertex at the receiving point.

#### BOUNDARY CONDITIONS

Because the fluid is unbounded, the solution to Equation 2 must satisfy boundary conditions at infinity as well as at the surface of the immersed airfoil. The conditions at infinity specify: (1) the fluid must be in uniform motion with no disturbances ahead of the downstream Mach lines from the



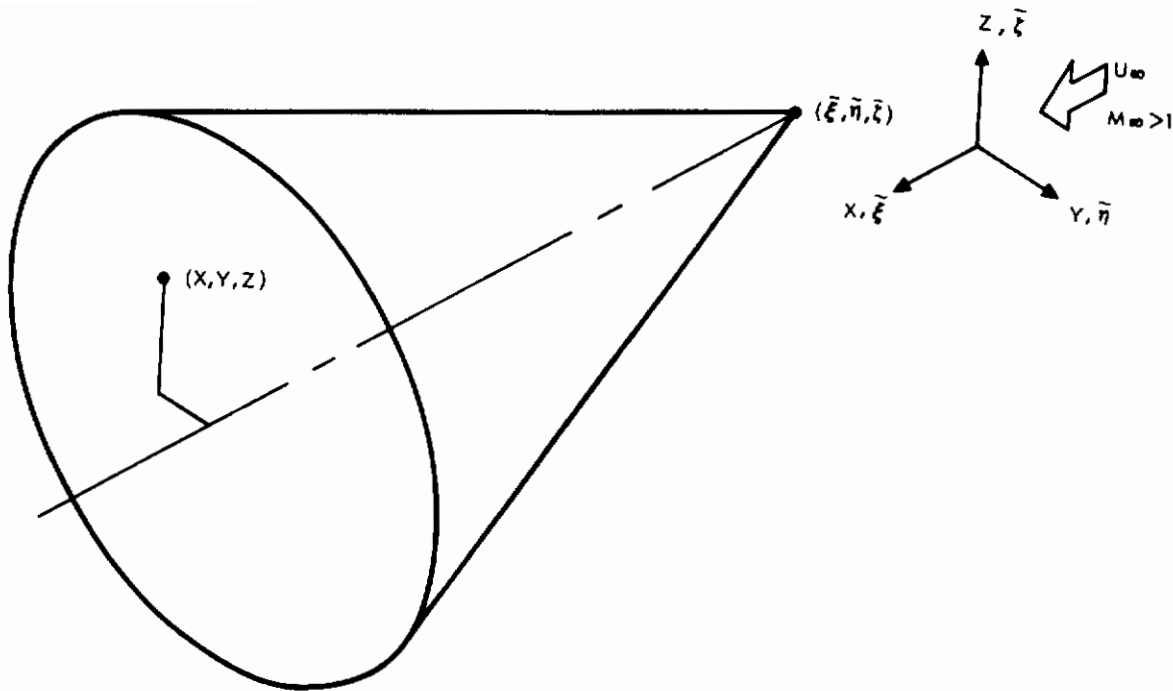


Figure 2. Small Disturbance and its Zone of Influence

leading edge of the surface, and (2) the disturbances behind these downstream Mach lines must be directed outwardly from their sources. The conditions at the surface of the thin airfoil require the flow to be tangent at all points on the surface. Setting the fluid velocity component normal to the surface equal to the normal velocity of the surface ensures this condition. The linearized expression of this condition for an airfoil lying close to the XY plane is

$$w(X, Y, Z, T) = \left[ U_\infty \frac{\partial}{\partial X} + \frac{\partial}{\partial T} \right] Z(X, Y, T) \quad (4)$$

where  $Z(X, Y, T)$  represents the instantaneous position of the mean surface moving normal to the XY plane. The condition in Equation 4 is applicable only when the slopes over the airfoil surface are small.

#### THE PRESSURE RELATION

The pressure relation, applicable to this problem and consistent with the assumptions of small perturbation theory, is given by Ashley (Reference 8) in the form

$$C_p = \frac{P - P_\infty}{\frac{1}{2} \rho_\infty U_\infty^2} = - \frac{2}{U_\infty} \phi_X - \frac{2}{U_\infty^2} \phi_T \quad (5)$$

where the pressure coefficient,  $C_p$ , is the ratio of the difference between the free-stream pressure and the local pressure,  $(p - p_\infty)$ , to the free-stream dynamic pressure,  $(\frac{1}{2} \rho_\infty U_\infty^2)$ . The difference between the pressure coefficient on the upper side and the lower side of the surface is obtained by replacing the velocity potential in Equation 5 with the change in the potential across the surface.

$$\Delta C_p = \frac{p^U - p^L}{\frac{1}{2} \rho_\infty U_\infty^2} = - \frac{2}{U_\infty} \Delta(\phi_X) - \frac{2}{U_\infty^2} \Delta(\phi_T) \quad (6)$$

This pressure difference coefficient is zero everywhere except at the airfoil surface.

## GENERALIZED FORCES

If a point on the airfoil surface undergoes a small virtual displacement  $(\delta r)$  in the direction of the pressure difference across the surface, virtual work will be done on the system.

$$\delta C_W = \Delta C_p \delta r \quad (7)$$

The nondimensional displacement  $\delta r$ , can be expressed in terms of the  $N$  independent generalized coordinates,  $q$ , as

$$\delta r = \sum_{i=1}^n \frac{\partial r}{\partial q_i} \delta q_i \quad (8)$$

For the case of the airfoil surface

$$\frac{\partial r}{\partial q_i} \equiv f_i$$

where  $f_i$  is the displacement normal to the surface in the  $i^{\text{th}}$  vibration mode. (See Section 1.4 of Reference 9 for a complete discussion of generalized coordinates and generalized forces.) The total virtual work done by all the applied pressures on the airfoil surface is then

$$\delta C_W = \frac{1}{S} \iint_S \Delta C_p \delta r \, dX \, dY \quad (9)$$

which becomes, upon substitution of Equation 8 and 8a for  $\delta r$

$$\delta C_W = \sum_{i=1}^n Q_i \delta q_i \quad (10)$$

The symbol  $Q_i$  refers to the generalized force which results from integrating the pressure times the displacements in the  $i^{\text{th}}$  vibration mode.

$$Q_i = \frac{1}{S} \iint_S \Delta C_p f_i dX dY \quad (11)$$

where  $S$  is the area of the airfoil surface.

## SIMPLE HARMONIC MOTION

The time dependence of the motion, if small in amplitude, can be completely arbitrary, since the solutions to Equation 2 are valid for any small motion; however, simple harmonic motion is assumed because of the wide applicability of the results. These results may be used in flutter analyses and in determination of transient response to externally applied loads such as gusts where the response can be expressed as a Fourier series consisting of simple harmonic terms.

Simple harmonic motion of the airfoil is expressible as

$$Z(X, Y, T) = \bar{Z}(X, Y) e^{i\omega T} \quad (12)$$

or in terms of generalized coordinates,

$$Z(X, Y, T) = \sum_{j=1}^n f_j \bar{q}_j e^{i\omega T} \quad (12a)$$

where  $i = \sqrt{-1}$  and  $\omega$  is the frequency of motion. Since the boundary conditions relate the disturbance velocity to the airfoil motion, substitution of Equation 12 into Equation 4 gives the following relation for the perturbation velocity

$$\bar{w}(X, Y) e^{i\omega T} = \left[ U_\infty \frac{\partial}{\partial X} + i\omega \right] \bar{Z}(X, Y) e^{i\omega T} \quad (13)$$

where the operator

$$U_\infty \frac{\partial}{\partial X} + i\omega = \frac{D}{DT} \quad (14)$$

is the linearized substantial derivative for simple harmonic motion and, therefore,  $\bar{\omega}(X, Y)$  is complex.

The velocity potential at a receiving point is solely determined by the free-stream velocity and the disturbances within the fore Mach cone from that point. Since the only disturbances are those resulting from the harmonically oscillating airfoil, the variation of the linearized perturbation velocity potential is also simple harmonic and may be written in the form:

$$\phi(X, Y, T) = \bar{\phi}(X, Y) e^{i\omega T} \quad (15)$$

Substitution of this expression for the velocity potential into Equation 2 results in

$$\beta^2 \bar{\phi}_{XX} - \bar{\phi}_{YY} - \bar{\phi}_{ZZ} = -\frac{1}{a_\infty^2} (2i\omega U_\infty \bar{\phi}_X - \omega^2 \bar{\phi}) \quad (16)$$

where

$$\beta = \sqrt{M_\infty^2 - 1}$$

The equation is now expressed in terms of the complex velocity potential,  $\bar{\phi}(X, Y)$ .

Substituting Equation 15 into the pressure-difference velocity-potential relation, Equation 6, reduces the expression for the change in pressure coefficient across the surface to:

$$\Delta \bar{C}_p = -\frac{2}{U_\infty} \Delta(\bar{\phi}_X) - \frac{2i\omega}{U_\infty^2} \Delta(\bar{\phi}) \quad (17)$$

where  $\Delta \bar{C}_p$  is the complex pressure difference coefficient.

## THE EQUIVALENT PROBLEM AT MACH NO. $\sqrt{2}$

As the final step in the development, a transformation of the Cartesian coordinates is introduced

$$x = X, \quad y = \beta Y, \quad z = \beta Z, \quad t = T \quad (18)$$

This step transforms all supersonic flow problems to equivalent problems at  $M_\infty = \sqrt{2}$ . Upon substitution of Equation 18, Equation 16 becomes

$$\bar{\phi}_{xx} - \bar{\phi}_{yy} - \bar{\phi}_{zz} = -\frac{1}{\beta^2 a_\infty^2} (2i\omega U_\infty \bar{\phi}_x - \omega^2 \bar{\phi}) \quad (19)$$

Equation 19 is satisfied within a 45-degree Mach cone downstream from the point disturbance. Substitution of Equations 18 and 17 into the expression for the generalized force in Equation 11 gives an expression in terms of the frequency of motion ( $\omega$ )

$$Q_i = \frac{-2}{\beta U_\infty S} \iint_S \left[ \Delta(\bar{\phi}_x) + \frac{i\omega}{U_\infty} \Delta(\bar{\phi}) \right] f_i \, dx \, dy \quad (20)$$

Now that the linearized equations describing the fluid motion are complete, there remains the problem of obtaining solutions to Equation 19 which satisfy the boundary conditions at infinity and at the surface of planar airfoils in supersonic flow. All equations and expressions in the ensuing development will be functions of the transformed coordinates  $(x, y, z), (\xi, \eta, \zeta)$  unless otherwise noted.

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### 3. SOURCE SUPERPOSITION METHOD

Consider a pulsating sound source to be fixed in an otherwise undisturbed supersonic flow that has been transformed to  $M_\infty = \sqrt{2}$ . Each spherical pulse is emitted  $t = t_i$  and expands with radius,  $a_\infty (t - t_i)$ , while its center moves downstream at the free-stream velocity,  $U_\infty$ . The boundary of the positions of the sphere is the familiar downstream Mach cone.

The basic solution for spherical waves emitted from a stationary sound source is

$$\phi_s(x, y, z, t) = \frac{H(\xi, \eta, \zeta)}{R} g(t)$$

where

$$R = \sqrt{(x - \xi)^2 + (y - \eta)^2 + (z - \zeta)^2}$$

Superposition of a parallel supersonic flow on the sound source transforms the time coordinate, and the solution becomes

$$\phi_s(x, y, z, t) = \frac{H(\xi, \eta, \zeta)}{R} \left[ g\left(t - \frac{R}{a_\infty}\right) + g\left(t + \frac{R}{a_\infty}\right) \right] \quad (21)$$

The source is fixed at the point  $(\xi, \eta, \zeta)$  with strength  $H(\xi, \eta, \zeta) g(t)$ . The two time functions,  $g(t - R/a_\infty)$  and  $g(t + R/a_\infty)$ , represent the receding and advancing portions of the waves, respectively. The explicit form of the fundamental solution at  $(x, y, z)$  corresponding to the amplitude of a harmonically pulsating source at  $(\xi, \eta, \zeta)$  in a parallel supersonic stream is given by Stewart (Reference 10) as,

$$\bar{\phi}(x, y, z, \xi, \eta, \zeta) = H(\xi, \eta, \zeta) G(x - \xi, y - \eta, z - \zeta) \quad (22)$$

where

$$G = -\frac{1}{\pi R} \exp \left[ \frac{-i\omega U_\infty}{2\beta^2} (x - \xi) \right] \cos \left( \frac{\omega}{a_\infty \beta^2} R \right)$$

which is parametrically dependent upon the free-stream conditions and the frequency of oscillation.



$H$  is the space variation of the source strength;  $\bar{\phi}$  is assigned the value zero everywhere outside the downstream Mach cone from  $(\xi, \eta, \zeta)$  and is singular all over the conical surface because the signals from a concentrated source reinforce one another there to an infinite degree. The elementary source, Equation 22, may be integrated with respect to any of the variables  $\xi, \eta, \zeta$  to give extended solutions to the transformed linearized equation of fluid motion, Equation 19.

## SINGLE PLANAR SURFACES

For the purpose of solving the thin planar lifting surface problem in supersonic flow, the mean surface of the airfoil is considered to lie nearly in the  $\zeta = 0$  plane. It can then be replaced by two sheets of the discontinuities given by Equation 22. One of the two sheets is placed on the upper surface of the airfoil and the other on the lower surface. The strength distributions of the source sheets are determined by the boundary conditions at the upper and lower sides of the surface, respectively.

The condition of tangential flow to the harmonically oscillating surface suggests that the strength of each source sheet is proportional to the normal velocity of the surface. Stewart (Reference 10) has shown that for a single planar surface lying in the  $\zeta = 0$  plane, the value of the strength of the upper source sheet is

$$H^U(\xi, \eta) = \bar{w}(\xi, \eta) \quad (23)$$

where  $\bar{w}(\xi, \eta)$  is the local normal component of the perturbation velocity. The antisymmetry of the normal velocity at any instant in time, with respect to the outward pointing normal on the upper and lower sides of the surface, implies that the lower source sheet is of opposite strength to the upper sheet, i.e.,

$$H^L(\xi, \eta) = -\bar{w}(\xi, \eta) \quad (24)$$

Keep in mind that a single source sheet causes disturbances in the flow which are symmetric relative to the plane of the source sheet. Therefore, it is essential to the source solution of the problem that we consider only those disturbances in the upper half space due to the upper source sheet and only those disturbances in the lower half space due to the lower source sheet. By placing two antisymmetric sheets of sources on the  $\zeta = 0$  plane, it is possible to obtain the proper lifting antisymmetry of the disturbance due to a thin airfoil oscillating in a supersonic main stream. It follows that the velocity potentials will also be antisymmetric.

If any portion of the edges of the airfoil is subsonic (i.e., the component of the free-stream velocity normal to some point on the edge is less than the acoustic velocity) the source superposition method seems to be inadequate. The underlying assumption that the upper and lower source sheets be isolated or not feel each other is violated in the region near a subsonic edge.

To apply the source sheet simulation to airfoils at all supersonic speeds Evvard (Reference 3) suggested extending the source sheets to



completely cover that portion of the  $\zeta = 0$  plane within the Mach envelope from the leading and trailing edges of the surface. The edges of the source sheets then become either sonic or supersonic, and disturbances from the upper and lower sheets are confined to the upper and lower half spaces, respectively, of the Mach envelope. The regions of the source sheet between a subsonic edge of the lifting surface and the Mach envelope are referred to as diaphragms and are placed thereto satisfy conditions of continuity of pressure and velocity and to prevent signals from being propagated from one half space to the other.

Since there is no surface motion in the diaphragm region, the perturbation velocity must be determined from other boundary conditions. These conditions are that the diaphragm causes no discontinuities in either the pressure,

$$\Delta \bar{p} = \bar{p}(\xi, \eta, 0+) - \bar{p}(\xi, \eta, 0-) = 0 \quad (25)$$

or the normal component of perturbation velocity

$$\Delta \bar{w} = \bar{w}(\xi, \eta, 0+) - \bar{w}(\xi, \eta, 0-) = 0 \quad (26)$$

The diaphragm normal velocity distribution must preserve the slopes which the streamlines would have if flow passed through the diaphragm regions. This condition is satisfied by making the source strength distribution proportional to the normal velocity distribution.

On the part of the diaphragm that is neither in the wake of its associated lifting surface nor in the wake of another body or surface upstream, an even stronger condition than that in Equation 25 has been formulated by Evvard (Reference 3). This stronger condition, which then replaces Equation 25, specifies that there can be no jump in velocity potential

$$\Delta \bar{\phi} = \bar{\phi}(\xi, \eta, 0+) - \bar{\phi}(\xi, \eta, 0-) = 0 \quad (27)$$

at any point in a diaphragm not in a wake region.

With the  $\zeta = 0$  plane of discontinuities dividing the Mach envelope into two half spaces, it is now possible to write the complete solution to Equation 19 for the velocity potential at a point  $(x, y, z)$  due to the presence of an oscillating lifting surface in a supersonic flow. The solution at any receiving point may be obtained by integrating both sides of Equation 22 over those portions of the airfoil and associated diaphragms that are within the zone of influence for that point. The velocity potential is given by

$$\bar{\phi}^{U, L}(x, y, z) = \frac{1}{\beta} \iint_A H^{U, L}(\xi, \eta) G(\bar{\xi}, \bar{\eta}, \bar{\zeta}) d\eta d\xi \quad (28)$$

where  $G$  is given in Equation 22 with  $\bar{\xi} = x - \xi$ ,  $\bar{\eta} = y - \eta$ ,  $\bar{\zeta} = z - \zeta$ . The sending area,  $A$ , is the region of the  $\zeta = 0$  plane bounded by the Mach envelope and the hyperbola  $\bar{\xi}^2 - \bar{\eta}^2 - \bar{\zeta}^2 = 0$ . Figure 3 shows a typical area on the upper side of the  $\zeta = 0$  plane that influences the point  $(x, y, z)$  in the upper half space. It is emphasized that the upper source sheet (with strength  $= \bar{w}(\xi, \eta)$ ) defines the velocity potential in the upper half space ( $\bar{\phi}^U$ ) and the lower source sheet (with strength  $= -\bar{w}(\xi, \eta)$ ) defines the potential in the lower half space ( $\bar{\phi}^L$ ).

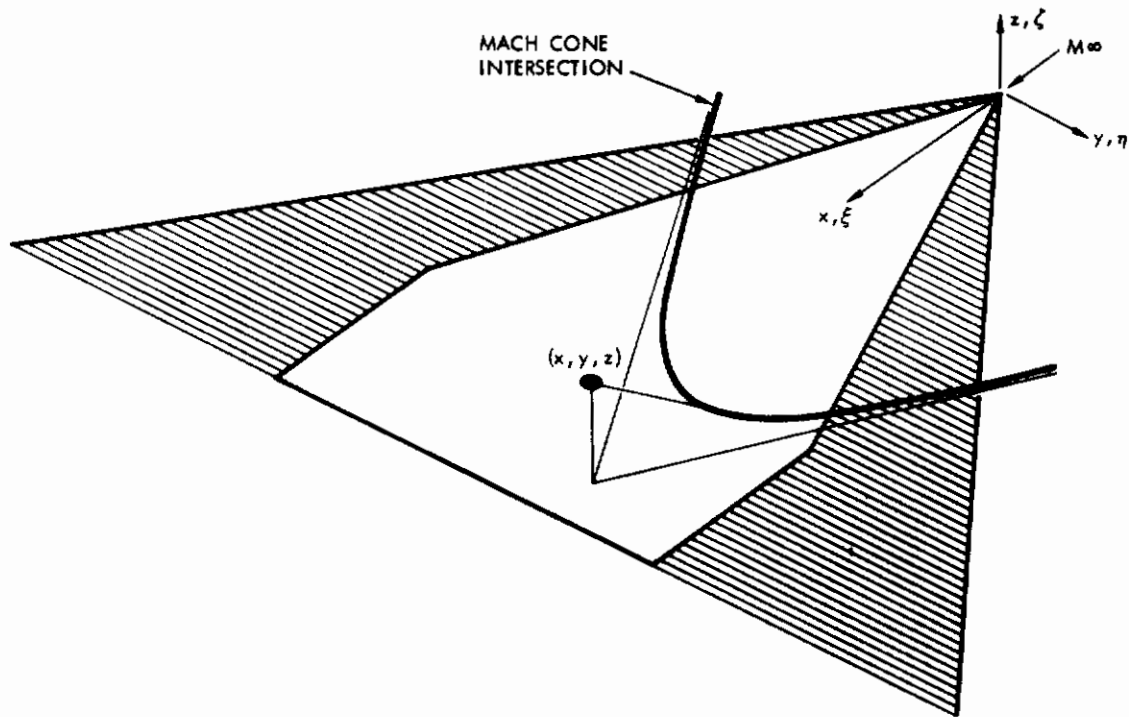


Figure 3. Source Sheet and Area of Influence for  $(x, y, z)$

Before the velocity potential can be calculated, using Equation 28, a complete description of the strength distribution over the source sheet will be necessary. The source strengths in the region that replaces the lifting surface result from substituting the tangential flow boundary condition (Equation 13) into Equation 23 for the upper surface

$$H^U(\xi, \eta) = \frac{D}{Dt} \bar{Z}(\xi, \eta) \quad (29)$$

and into Equation 24 for the lower surface

$$H^L(\xi, \eta) = - \frac{D}{Dt} \bar{Z}(\xi, \eta) \quad (30)$$

The source-sheet strength distribution in the diaphragm regions can be obtained by applying to the pressure difference coefficient the general condition of no pressure jump across the region. However, if we allow only subsonic leading and side edges, Equation 28 can be substituted directly into Equation 27. For a receiving point  $(x, y, 0)$  ahead of or beside the airfoil surface, but within the Mach envelope, Equation 27 becomes

$$0 = \iint_A \Delta H(\xi, \eta) G(\bar{\xi}, \bar{\eta}, \bar{\zeta}) d\eta d\xi \quad (31)$$

where the area of integration,  $A$ , includes portions of both the airfoil surface and its associated diaphragms within the fore-Mach triangle from the point  $(x, y, 0)$ . Separation of the integral into a sum of integrals over the airfoil and diaphragms, respectively, and substitution of the strength-per-unit area in the region representing the airfoil, produces an integral equation,

$$\begin{aligned} & \iint_{A_D} \Delta H(\xi, \eta) G(\bar{\xi}, \bar{\eta}, \bar{\zeta}) d\eta d\xi \\ & = -2 \iint_{A_S} \frac{D}{Dt} \bar{Z}(\xi, \eta) G(\bar{\xi}, \bar{\eta}, \bar{\zeta}) d\eta d\xi \end{aligned} \quad (32)$$

for the source-sheet strength per unit area in the diaphragm region. Areas of integration,  $A_D$  and  $A_S$ , for a typical lifting surface are shown in Figure 4. For the single-planar lifting surface, separate integral equations for the upper and lower diaphragm strength distributions can be obtained. Applying the condition of antisymmetry to the normal component of perturbation velocity, Equation 32 becomes

$$\begin{aligned} & \iint_{A_D} H^{U,L}(\xi, \eta) G(\bar{\xi}, \bar{\eta}, \bar{\zeta}) d\eta d\xi \\ & = \mp \iint_{A_S} \frac{D}{Dt} \bar{Z}(\xi, \eta) G(\bar{\xi}, \bar{\eta}, \bar{\zeta}) d\eta d\xi \end{aligned} \quad (33)$$

where the plus sign (+) goes with the lower source sheet and the minus sign (-) goes with the upper source sheet.

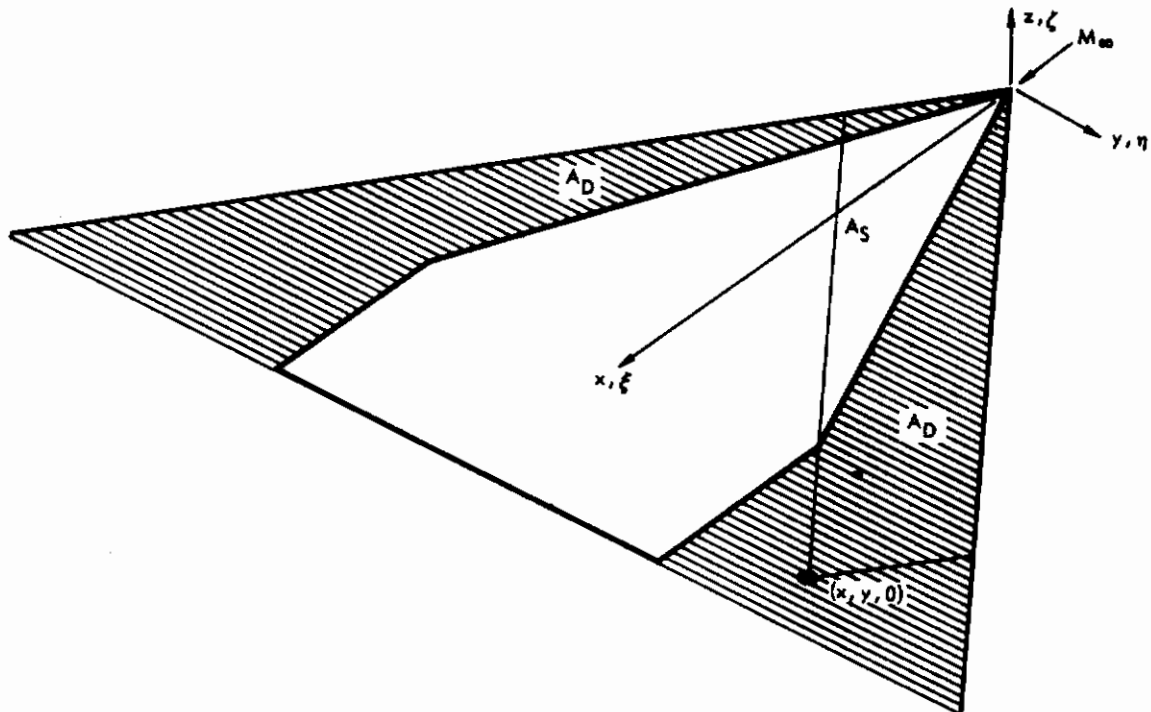


Figure 4. Source Sheet and Area of Influence for  $(x, y, 0)$

It is emphasized that the foregoing development of the source-sheet solution to Equation 19 for a single-planar lifting surface at any supersonic speed provides the proper lifting antisymmetry only through proper use of the diaphragm in the  $\zeta = 0$  plane between a subsonic edge and the Mach envelope. The diaphragm must isolate the upper and lower airfoil surface source sheets while ensuring that no pressure loading acts on regions that are incapable of sustaining it.

The following review of the Mach-box scheme outlines the method for determining planar aerodynamic influence coefficients for velocity potentials and applying them to obtain lifting pressures and generalized forces.

If we let the wing and diaphragm in the  $\zeta = 0$  plane (Figure 4) be covered with a grid of boxes of length  $b$  in the flow direction and width  $b/\beta$ , their diagonals are then parallel to the Mach lines (Figure 5). The coordinate system has been nondimensionalized by the box length  $b$  so that the integral equation for the perturbation potential, Equation 28, at the point  $(x_1, y_1, z_1)$  becomes

$$\bar{\phi}^{U, L}(x_1, y_1, z_1) = \frac{b}{\beta} \iint_A H^{U, L}(\xi_1, \eta_1) G(\xi_1, \bar{\eta}_1, \bar{z}_1) d\eta_1, d\xi_1 \quad (34)$$

where

$$G(\bar{\xi}_1, \bar{\eta}_1, \bar{\zeta}_1) = \frac{-1}{\pi R_1} \exp \left[ -ik\bar{\xi}_1 \right] \cos \left( \frac{\bar{k}}{M_\infty} R_1 \right)$$

$$R_1 = \sqrt{\bar{\xi}_1^2 - \bar{\eta}_1^2 - \bar{\zeta}_1^2}$$

$$\bar{k} = \frac{\omega b}{U_\infty} \frac{M_\infty^2}{\beta^2}$$

and the new coordinates are

$$x_1, \xi_1 = \frac{x, \xi}{b} ; \quad y_1, \eta_1 = \frac{y, \eta}{b} ; \quad z_1, \zeta_1 = \frac{z, \zeta}{b}$$

$$\bar{\xi}_1 = x_1 - \xi_1 ; \quad \bar{\eta}_1 = y_1 - \eta_1 ; \quad \bar{\zeta}_1 = z_1 - \zeta_1$$

Now, assume that the strength distribution over each rectangular portion of the superposed source sheets is constant and equal to the value at the box center, and consider a sending area or Mach box to be centered at

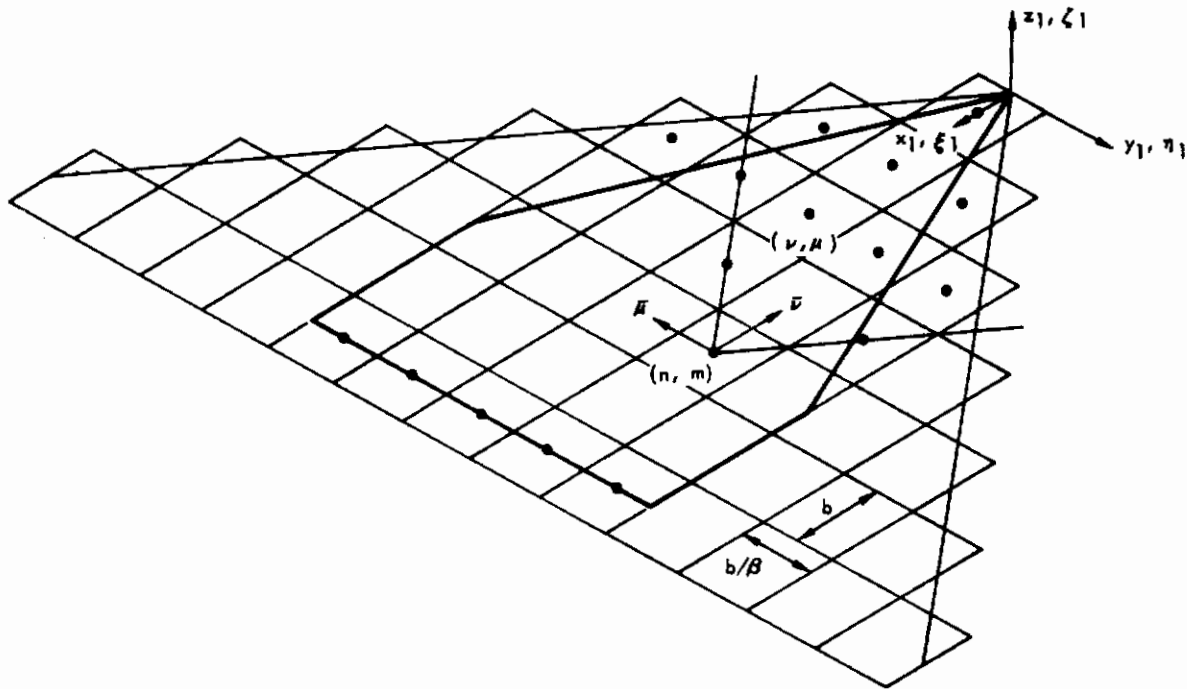


Figure 5. Boxes on Figure 4



the point  $\xi_1 = \nu$ ,  $\eta_1 = \mu$ ,  $\zeta_1 = 0$ , while a receiving point is located at the position  $x_1 = n$ ,  $y_1 = m$ ,  $z_1 = 0$ . When the source strength,  $H^{U,L}(\nu, \mu)$ , is constant, the velocity potential at the receiving point is

$$\phi^{U,L}(n, m, 0) = \frac{b}{\beta} \sum H^{U,L}(\nu, \mu) \Phi(\bar{\nu}, \bar{\mu}, 0) \quad (35)$$

where  $\Phi(\bar{\nu}, \bar{\mu}, 0)$  is the planar velocity potential influence coefficient (VPIC) at the receiving point, due to a source box of unit strength with the center located  $\bar{\nu} = n - \nu$  box lengths ahead and  $\bar{\mu} = m - \mu$  box widths to the left. The VPIC is dependent only on the relative distances  $\bar{\nu}$  and  $\bar{\mu}$ , and the Mach number and reduced frequency parameters. The summation in Equation 35 is extended over all boxes wholly or partially within the fore Mach cone (Figure 5). This then divides the area of integration in Equation 34 into several subareas over which the planar VPIC may be evaluated using the equation

$$\Phi(\bar{\nu}, \bar{\mu}, \ell) = \iint_{A_{\text{box}}} G(\bar{\xi}_1, \bar{\eta}_1, \ell) d\bar{\eta}_1 d\bar{\xi}_1 \quad (36)$$

where  $\ell = 0$  and where  $\bar{\xi}_1 = n - \xi_1$  and  $\bar{\eta}_1 = m - \eta_1$ . When a point lies just above

When a point lies just above a source sheet, the only perturbation velocity there is due to the outwardly directed strength at the adjacent point on the sheet.

$$H^U(n, m) = \frac{D}{Dt} \bar{Z}(n, m) \quad (37)$$

and

$$H^L(n, m) = - \frac{D}{Dt} Z(n, m) \quad (38)$$

On the other hand, the strength distribution over a diaphragm box not in a wake must be determined by application of the condition that the velocity potential is continuous across the diaphragm. This condition is satisfied by summing Equation 35 over all upper and lower boxes with known source strengths in the fore cone of the receiving point and over the upper and lower sides of the box of unknown source strength at the receiving point and setting the total equal to zero. Solution for the unknown source strength gives:

$$\Delta H(n, m) = - \sum \Delta H(\nu, \mu) \Phi(\bar{\nu}, \bar{\mu}, 0) / \Phi(0, 0, 0) \quad (39)$$

because continuity of velocity through the diaphragm means that the upper and lower source strengths are equal in magnitude and opposite in sign, Equation 39 can be written

$$H^{U,L}(n, m) = \mp \sum H^{U,L}(\nu, \mu) \Phi(\bar{\nu}, \bar{\mu}, 0) / \Phi(0, 0, 0) \quad (40)$$

where the minus or plus signs of Equation 40 are to be used with the superscripts U or L, respectively.

Calculation of the velocity potential discontinuity across the airfoil surface may now be performed by applying to the appropriate boxes the relationships given in Equations 35, 37, 38, and 40. To eliminate the need for solving sets of simultaneous equations for the unknown strength distribution over the diaphragm, the following rules for calculation are followed:

1. Starting at the foremost centerline box, calculate the change in velocity potential across the surface at all box centers in the first row
2. Return to the centerline box in the next row; calculate the velocity potential change at all surface box centers in that row and then the source strengths at box centers on the diaphragm
3. Repeat procedure 2 for each subsequent row of boxes

The resulting velocity potential distribution over the surface can then be substituted into Equation 17 to obtain the pressure difference coefficient, or into Equation 20 to obtain the generalized force.

## INTERSECTING PLANAR SURFACES

The following development parallels that of Ashley's (References 6, 11, and 12) and presents the results in a form amenable to high-speed computational techniques.

The success of the extension of the source superposition method to intersecting surfaces depends heavily on careful application of the diaphragm concept to regions between subsonic edges and Mach envelopes. Before developing the mathematical expressions for the necessary source strength distributions and resulting velocity potentials over these surfaces, the following general statements are made concerning the construction of diaphragm regions:

1. A diaphragm is a device used to ensure continuity of pressure and velocity between every pair of adjacent points in the flow field except those on opposite sides of lifting surfaces. When adjacent points are not on opposite sides of a surface, or the wake of a surface, the stronger condition of continuity of velocity potential replaces the condition of continuity of pressure. In particular, the diaphragm is placed between all pairs of adjacent points that receive signals from some pair of opposite half-envelopes.

2. There is no unique set of diaphragms for a particular configuration and Mach number. Any set which provides the required continuity of pressure and stream lines is acceptable. Diaphragms do not necessarily have to be coplanar with any lifting surface to meet this requirement; furthermore, diaphragms do not even have to be planar surfaces.
3. The concept is easiest to envision and much easier to apply if diaphragms are parallel to the free-stream direction.

Also, the following general statements and definitions are made concerning interference regions:

1. A portion of a surface is in an interference region when it is within the Mach envelope of another surface.
2. Mutual interference exists between two surfaces when a portion of each surface lies within the Mach envelope of the other surface.

The basic difference between application of the source superposition method to isolated planar lifting surfaces and to a collection of intersecting lifting surfaces lies in the way in which the strength distribution of the various superposed source sheets is determined. In both cases, the strength is determined from either the tangential flow or continuous flow boundary conditions.

When a point that lies just above a source sheet is not in a region of interference, the only perturbation velocity there is due to the outwardly directed strength at the adjacent point in the sheet; however, if the point is in an interference region it will not only be within the half-envelope associated with its source sheet, but also within the upper or lower half-envelopes associated with one or more of the other surfaces. In addition to the velocity due to the point source, velocities are induced there by outwardly directed disturbances on portions of the other source sheets. The portions of these source sheets referred to will be those areas within the Mach hyperbola formed on each sheet by its intersection with the fore Mach cone from the receiving point.

The strength of the local source sheet is still unknown and is designated  $H^U(x, y)$ . On the other hand, the strengths of the other source sheets are presumed to be known and, therefore, the induced normal velocity at the receiving point,  $\partial\bar{\phi}/\partial z$ , may be determined. The induced velocity potential is  $\bar{\phi}$ . The sum of all the perturbation velocities in the direction of the outwardly directed normal to the upper surface may be written as

$$\bar{w}^U(x, y) = H^U(x, y) + \left(\frac{\partial\bar{\phi}}{\partial z}\right)^U \quad (41)$$



The outwardly directed normal on the lower side of the surface is opposite in sense to that on the upper side. Since the two sides move together the velocity of the lower side, in the direction of the outwardly directed normal on that side, is the negative of its upper counterpart. Thus,  $\bar{w}^L(x, y) = -\bar{w}^U(x, y)$ . The source strength at the lower point is still unknown and not necessarily proportional to its upper counterpart. To complete the description of the normal velocity at the lower point, the induced velocities in the direction of the lower normal are  $\partial \bar{\phi}^L / \partial n^L = -\partial \bar{\phi}^L / \partial z$ . Thus, the expression for the sum of all the perturbation velocities normal to the lower surface becomes, after multiplying both sides by -1.0,

$$\bar{w}^U(x, y) = -H^L(x, y) + \left(\frac{\partial \bar{\phi}}{\partial z}\right)^L \quad (42)$$

The superscripts on w in Equations 41 and 42 may now be dropped and the unknown local source strengths on the upper and lower parts of the surface may be determined separately.

All portions of each surface, whether or not they are in the interference region, will have flow tangent to both upper and lower sides if the perturbation velocity at the surface in the positive  $Z_i$  direction is set equal to the velocity of the mean surface in that direction,

$$\bar{w}(x_i, y_i) = \frac{D}{Dt} \bar{Z}_i(x_i, y_i) \quad (43)$$

where the coordinates  $x_i, y_i, z_i^*$  are placed so that the  $i$ th surface lies near the  $z_i = 0$  plane and  $Z_i$  represents its variation from that plane. For the region covering the  $i$ th intersecting lifting surface, the  $i$ th upper source-sheet strength is given by substitution of Equation 41 into Equation 43,

$$\frac{D}{Dt} \bar{Z}_i(x_i, y_i) = H^U(x_i, y_i) + \left(\frac{\partial \bar{\phi}}{\partial z_i}\right)^U \quad (44)$$

and the  $i$ th lower source-sheet strength is given by substitution of Equation 44 into Equation 43

$$\frac{D}{Dt} \bar{Z}_i(x_i, y_i) = -H^L(x_i, y_i) + \left(\frac{\partial \bar{\phi}}{\partial z_i}\right)^L \quad (45)$$

The velocity potential,  $\bar{\phi}$ , at any point may be determined by adding contributions induced there by the various source sheets associated with the Mach half-envelopes within which the point lies. For a point on the  $i$ th upper

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\* The subscripts  $i, j$  are a convenient way to distinguish the  $i$ th receiving point or plane from the  $j$ th sending point or plane and are not to be confused with the subscripts used for normalization of coordinates.

surface, the velocity potential is due to all the disturbances on the  $i$ th upper source sheet within the fore Mach triangle as well as those disturbances from upstream out-of-plane source sheets, i.e.,

$$\begin{aligned} \bar{\phi}^U(x_i, y_i) = & \frac{1}{\beta} \iint_{A_i} H^U(\xi_i, \eta_i) G(\bar{\xi}_i, \bar{\eta}_i, \bar{\zeta}_i) d\eta_i d\xi_i \\ & + \left[ \frac{1}{\beta} \sum_j \iint_{A_j} H^{U,L}(\xi_j, \eta_j) G(\bar{\xi}_j, \bar{\eta}_j, \bar{\zeta}_j) d\eta_j d\xi_j \right]^U \end{aligned} \quad (46)$$

where  $H^{U,L}(\xi_j, \eta_j)$  is the  $j$ th upper or lower source sheet strength distribution within the area,  $A_j$ , that affects the point  $(x_i, y_i)$ .  $G$  is the velocity potential at the point  $(x_i, y_i)$  due to a source of unit strength located at a sending point  $(\xi_i, \eta_i, \zeta_i)$  upstream. In addition to its parametric dependence upon the free-stream Mach number and frequency of oscillation,  $G$  is dependent only on the relative positions of the receiving and sending points. The velocity potential at the point  $(x_i, y_i)$  on the  $i$ th lower source sheet is given by

$$\begin{aligned} \bar{\phi}^L(x_i, y_i) = & \frac{1}{\beta} \iint_{A_i} H^L(\xi_i, \eta_i) G(\bar{\xi}_i, \bar{\eta}_i, \bar{\zeta}_i) d\eta_i d\xi_i \\ & + \left[ \frac{1}{\beta} \sum_j \iint_{A_j} H^{U,L}(\xi_j, \eta_j) G(\bar{\xi}_j, \bar{\eta}_j, \bar{\zeta}_j) d\eta_j d\xi_j \right]^L \end{aligned} \quad (47)$$

where the area,  $A_i$ , is the Mach triangle on the  $i$ th lower source sheet. The second term on the right side of Equations 46 and 47 represents the velocity potential induced at the point by the upstream disturbances.

For points on diaphragms that are not in the wake of any upstream disturbances, the condition for continuous velocity across the diaphragm is given by Equation 26 which, upon substitution of Equations 41 and 42, becomes

$$0 = H^U(x_i, y_i) + H^L(x_i, y_i) + \left[ \left( \frac{\partial \bar{\phi}}{\partial z} \right)^U - \left( \frac{\partial \bar{\phi}}{\partial z} \right)^L \right] \quad (48)$$

An additional equation is necessary to solve the unknown upper and lower source strengths at a point on the diaphragm. The condition of continuity of velocity potential across the diaphragm is obtained by substitution of Equations 47 and 46 into Equation 27 to get

$$\begin{aligned} \iint_{A_i} \Delta H(\xi_i, \eta_i) G(\bar{\xi}_i, \bar{\eta}_i) d\eta_i d\xi_i &= 0 \\ &= - \sum_j \iint_{A_j} \Delta H^{U,L}(\xi_j, \eta_j) G(\bar{\xi}_j, \bar{\eta}_j, \bar{\zeta}_j) d\eta_j d\xi_j \quad (49) \end{aligned}$$

Once the velocity potential distribution over the surface is obtained, the pressure loading and generalized forces may be obtained by substituting the discontinuity in velocity potential  $\Delta \bar{\phi} = \bar{\phi}^U - \bar{\phi}^L$ , across the surface into Equations 17 and 20, respectively.

#### A Discussion of Possible Approaches

The previous discussion implies that a single diaphragm may be associated with two or more lifting surfaces. As an example, consider a thin wing with foldable tips and subsonic leading edges (Figure 6). As the tips fold from 0 to 90 degrees, the diaphragm might also be folded and then extended to a hyperbolic intersection with the Mach envelope of the wing (Figure 7). The diaphragms that are coplanar with the folded tips actually

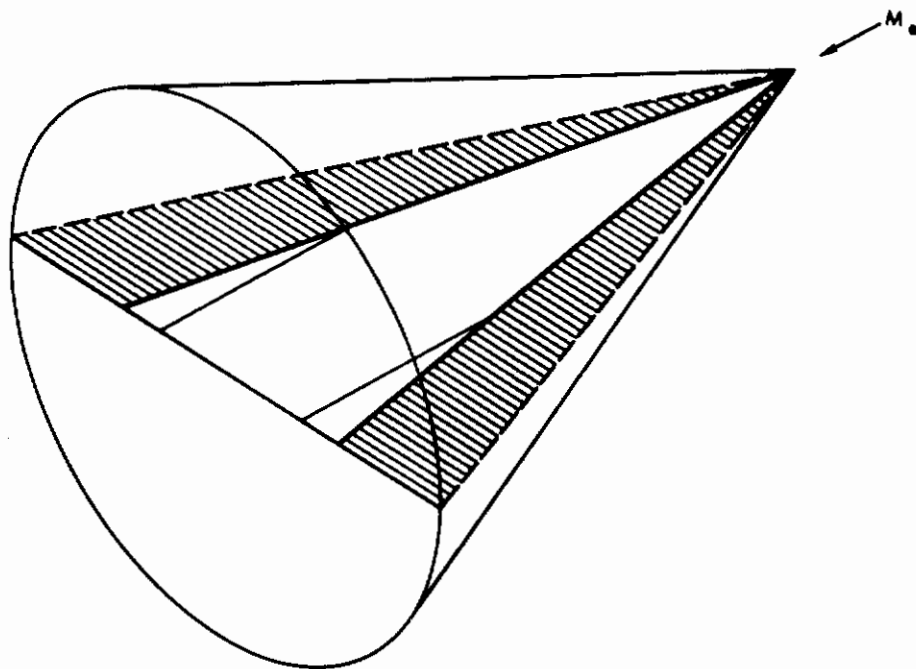


Figure 6. Planar Wing With Tips at 0-Degree Fold

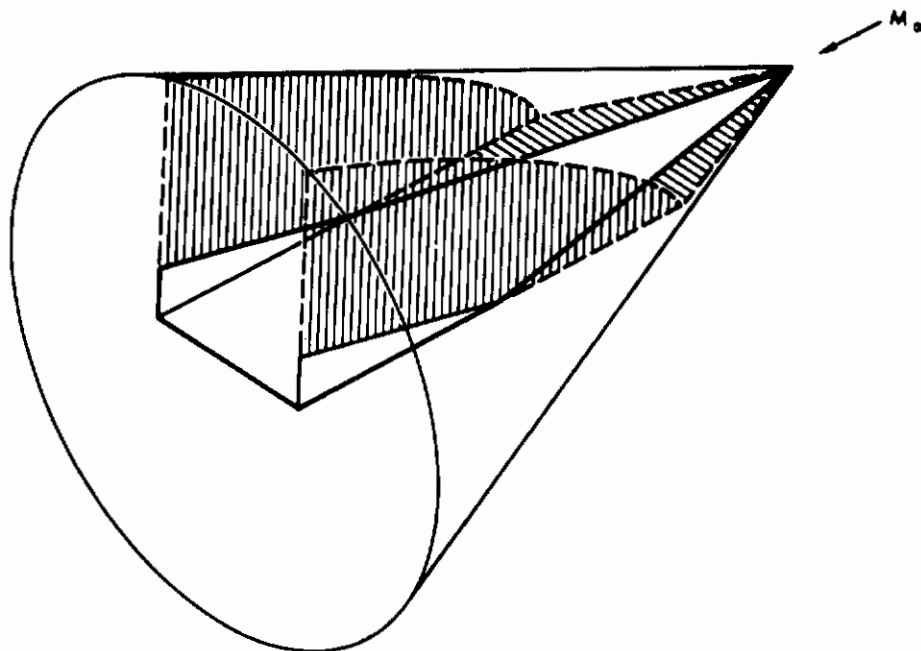


Figure 7. Planar Wing With Tips at 90-Degree Fold

serve to complete isolation of the upper and lower sides of the wing as well as both sides of the tip. The flow is divided into two regions, and points in one region are not influenced by disturbances in the other region except through the diaphragm regions which are now nonplanar and will require three-dimensional calculations. Extending diaphragms until they intersect with the largest Mach envelope not only causes an unnecessarily large interference region between source sheets but also makes the area of the diaphragm dependent upon the angle of intersection between the surfaces.

Another approach is one in which the diaphragm remains fixed in the plane of the wing as the tip folds from zero degrees. To complete the isolation of the upper surface from the lower surface with a small amount of out-of-plane diaphragm, a diaphragm is placed between the leading edge of the tip and the wing diaphragm (Figure 8). This diaphragm is not parallel with the free-stream direction which is permissible as long as the replacement source sheets do not alter the steady-state direction of the free stream. The nonparallel-to-flow diaphragm does complicate the three-dimensional calculations required for out-of-plane source sheet interference. While this diaphragm construction represents an apparent minimum in out-of-plane calculations, its area and angle with the flow are both dependent upon the angle of intersection of the two surfaces. These problems together with the numerical difficulties associated with the three-dimensional construction offset any advantages obtained by this approach.

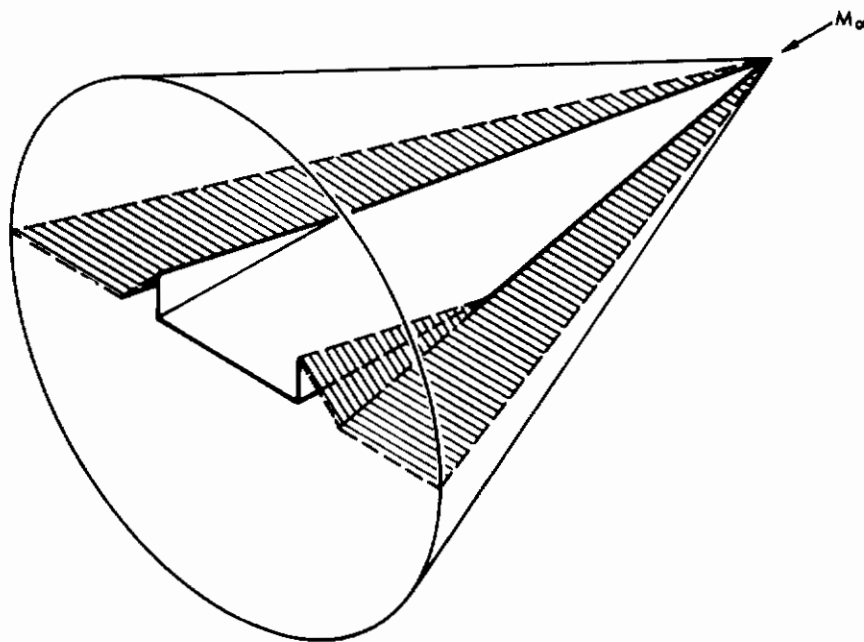


Figure 8. Alternate Diaphragm Construction



Another approach which is also the result of trying to reduce the diaphragm area to a practical minimum is one step removed from the approach suggested for future development. In this approach the several intersecting surfaces are separated into isolated planar regions and a Mach envelope is constructed for each surface. After superposing source sheets on the upper and lower sides of the surface, each sheet is extended to intersect the associated Mach envelope. If the surfaces are actually isolated, the disturbances from each upper or lower source sheet will be confined to the respective half-envelope (Figure 9). The configuration is then reassembled and the relationship between each surface and its Mach envelope is retained. (Figure 10)

At this point a step is performed which is shown later to increase rather than decrease the number of required numerical operations. Since one region may serve to isolate the upper and lower sides of more than one surface, it may be that some of the diaphragms are unnecessary when the surfaces are rejoined. They will be unnecessary and may be removed if continuity of pressure and velocity can be ensured without them.

To demonstrate this approach consider the wing with folded tips to be separated into three isolated surfaces. Figure 9 shows the three surfaces, their Mach envelopes, and their diaphragm regions for a particular Mach number. When the two tips are reattached to the wing, the tip diaphragms in the lower half-envelope of the wing may be folded into the wing diaphragm. This particular set of diaphragms is shown in Figure 11, where the tips are folded to 90 degrees.

With this set of boundaries, a point on the upper surface of the wing receives from the upper parts of the tip source sheets as well as from the upper part of the wing source sheets. The adjacent point on the lower surface of the wing receives only from the lower parts of the wing source sheet. The boundary condition applied on each side of the wing surface is that of tangential flow.

A point on either side of the tip surface receives from the upper part of the wing source sheets as well as from the tip source sheets on its side. Finally, a point on that portion of the upper surface of the wing diaphragm which is in the interference region receives from the lower parts of the tip source sheets as well as from the upper parts of the wing source sheets. Table 1 shows the sending and receiving relationships described above and shown in Figure 12.

This approach, described previously, is used in the computer program. Subsequent development has shown that it is not the most efficient approach but it does serve to demonstrate one of the options.

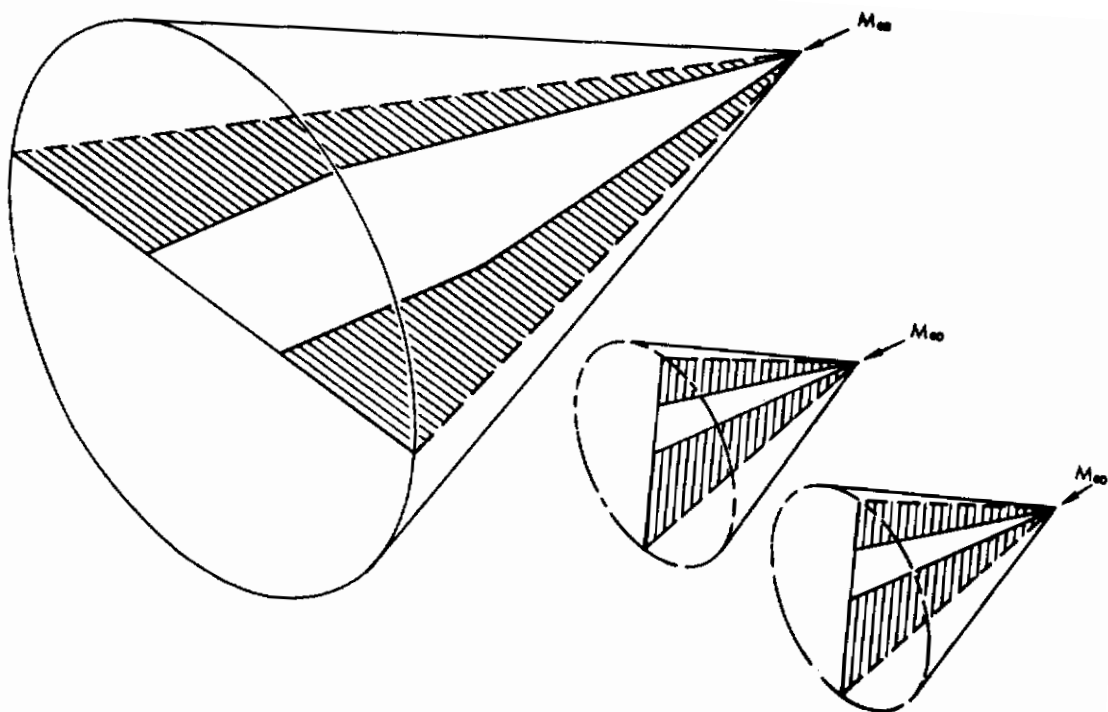


Figure 9. Wing and Tips Isolated

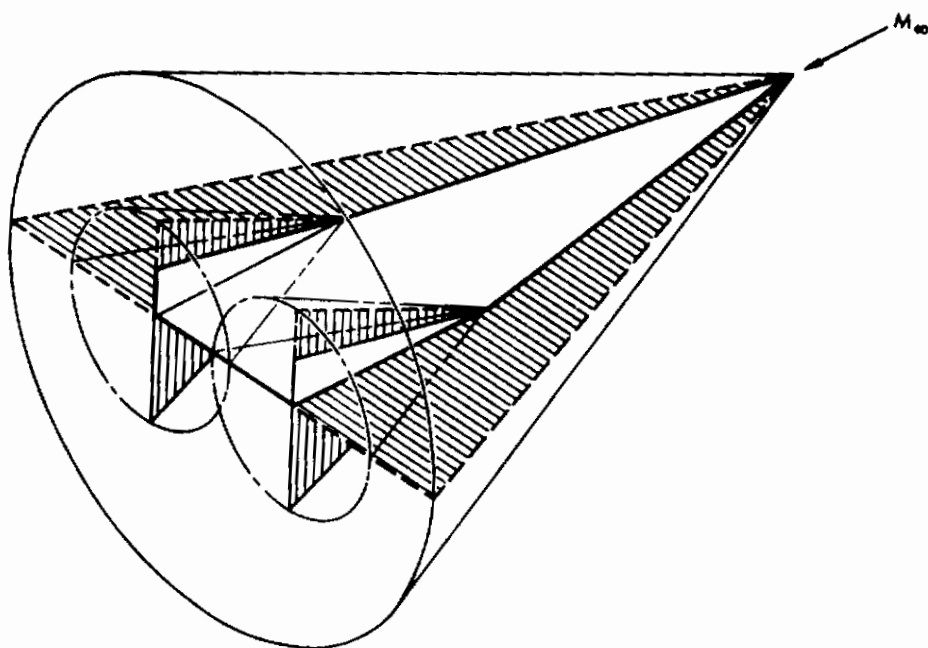


Figure 10. Wing and Tips Reattached

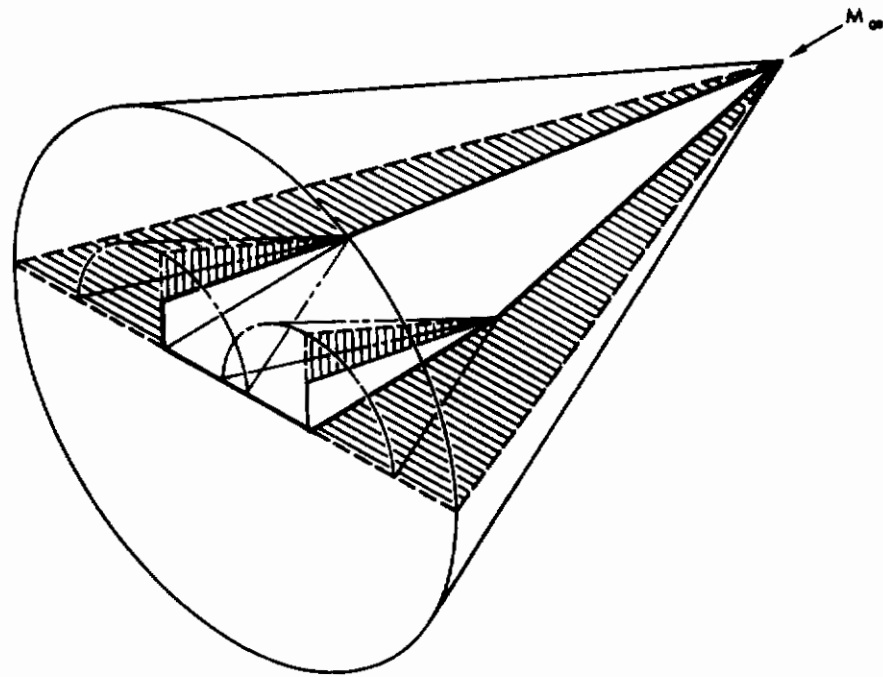


Figure 11. Wing and Tips Reattached - Lower Diaphragm Removed

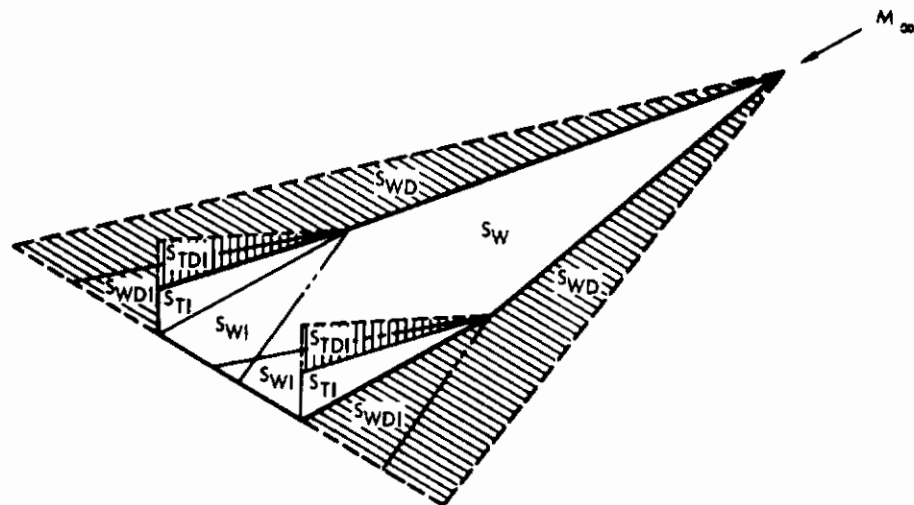


Figure 12. Zones on Figure 11



Table 1. Legend to Determine the Sending and Receiving Relationship Between Regions of the Configuration Shown in Figure 11

$\begin{array}{c} \nearrow j \\ i \end{array}$		UPPER				LOWER			
		SW & SWD	SWI	SWDI	STI & STDI	SW & SWD	SWI	SWDI	STI & STDI
U P P E R	SW & SWD	0	0	0	0				0
	SWI & SWDI		0	0	0				0
	STI & STDI		0		0				
L O W E R	SW & SWD					0	0	0	
	SWI & SWDI						0	0	
	STI & STDI			0					0

### The Aerodynamic Influence Coefficient Method

The wing with symmetrically folded tips is analyzed and presented as an example of one of the more efficient of the possible approaches to the problem of intersecting lifting surfaces. The discussion concerns only one-half of the wing because the configuration has a plane of symmetry through the centerline.

We now overlay a grid of Mach boxes on the wing, tips, and associated diaphragms of the surfaces shown in Figure 12. The result is shown in Figure 13. Note that the grid is arranged so that box centers lie along the centerline and box edges coincide with the line of intersection of the wing and tip.

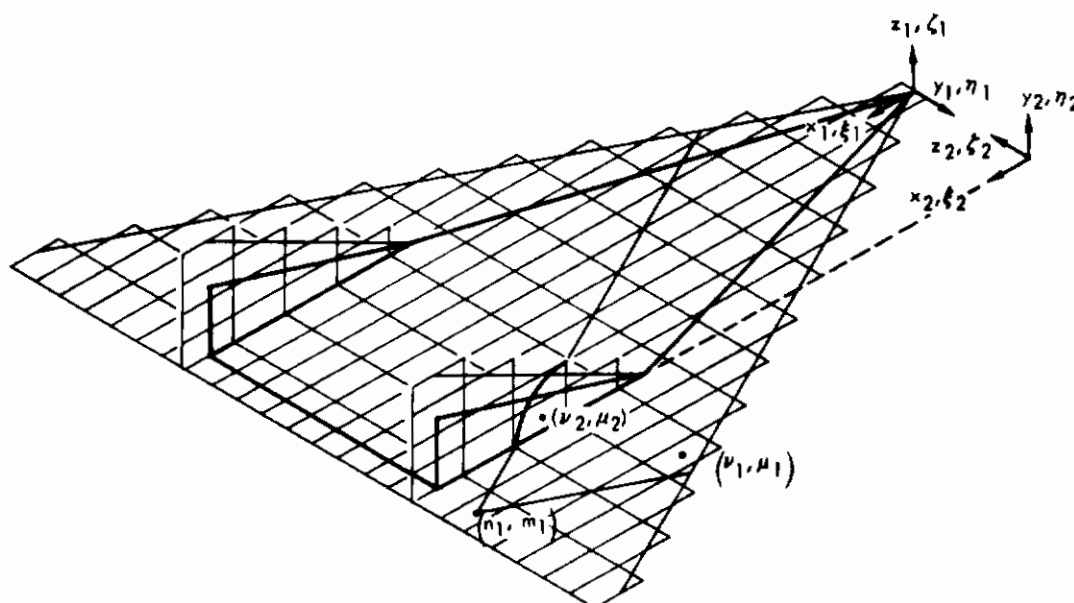


Figure 13. Boxes on Figure 12 and Area of Influence for  $(v, \mu)$

The wing coordinate system  $(x_1, y_1, z_1)$  and the tip coordinate system  $(x_2, y_2, z_2)$  enable relative distances between sending and receiving boxes to be set up as follows:

Wing boxes  $(n_1, m_1)$  receiving from wing boxes  $(v_1, \mu_1)$

$$\bar{v}_{11} = n_1 - v_1 \quad (50)$$

$$\bar{\mu}_{11} = m_1 - \mu_1$$

$$l_{11} = 0$$

Wing boxes  $(n_1, m_1)$  receiving from tip boxes  $(v_2, \mu_2)$

$$\bar{v}_{12} = n_1 - v_2 \quad (51)$$

$$\bar{\mu}_{12} = -(y_{ff} - m_1) \cos \gamma - \mu_2$$

$$l_{12} = -(y_{ff} - m_1) \sin \gamma$$

Tip boxes  $(n_2, m_2)$  receiving from wing boxes  $(v_1, \mu_1)$

$$\bar{v}_{21} = n_2 - v_1 \quad (52)$$

$$\bar{\mu}_{21} = m_2 \cos \gamma + (y_{ff} - \mu_1)$$

$$l_{21} = m_2 \sin \gamma$$

Tip boxes  $(n_2, m_2)$  receiving from tip boxes  $(v_2, \mu_2)$

$$\bar{v}_{22} = n_2 - v_2 \quad (53)$$

$$\bar{\mu}_{22} = m_2 - \mu_2$$

$$l_{22} = 0$$

The fold angle,  $\gamma$ , is measured positive when the tips fold up from the wing plane and  $y_{fl}$  is the  $y_1$  distance to the fold line or  $x_2$  axis.

The continuous strength distribution over the upper and lower source sheets on both tips and the wing is replaced by the discontinuous distribution that results when the source sheet on each box is assumed to have a constant strength equal to the value computed for its center. With this approximation and the information contained in Table 1, Equations (28), (46), and (47) are written for the upper and lower sides of the wing and tip surfaces as follows:

$$S_W: \bar{\phi}^U(n_1, m_1) = \frac{b}{\beta} \sum H^U(\nu_1, \mu_1) \Phi(\bar{\nu}_{11}, \bar{\mu}_{11}, 0) \quad (54)$$

$$S_W: \bar{\phi}^L(n_1, m_1) = \frac{b}{\beta} \sum H^L(\nu_1, \mu_1) \Phi(\bar{\nu}_{11}, \bar{\mu}_{11}, 0) \quad (55)$$

$$\begin{aligned} S_{WI}: \bar{\phi}^U(n_1, m_1) &= \frac{b}{\beta} \sum H^U(\nu_1, \mu_1) \Phi(\bar{\nu}_{11}, \bar{\mu}_{11}, 0) \\ &+ \frac{b}{\beta} \sum H^U(\nu_2, \mu_2) \Phi(\bar{\nu}_{12}, \bar{\mu}_{12}, \ell_{12}) \end{aligned} \quad (56)$$

$$\bar{\phi}^L(n_1, m_1) = \frac{b}{\beta} \sum H^L(\nu_1, \mu_1) \Phi(\bar{\nu}_{11}, \bar{\mu}_{11}, 0) \quad (57)$$

$$\begin{aligned} S_{TI}: \bar{\phi}^U(n_2, m_2) &= \frac{b}{\beta} \sum H^U(\nu_2, \mu_2) \Phi(\bar{\nu}_{22}, \bar{\mu}_{22}, 0) \\ &+ \frac{b}{\beta} \sum H^U(\nu_1, \mu_1) \Phi(\bar{\nu}_{21}, \bar{\mu}_{21}, \ell_{21}) \end{aligned} \quad (58)$$

$$\begin{aligned} \bar{\phi}^L(n_2, m_2) &= \frac{b}{\beta} \sum H^L(\nu_2, \mu_2) \Phi(\bar{\nu}_{22}, \bar{\mu}_{22}, 0) \\ &+ \frac{b}{\beta} \sum H^L(\nu_1, \mu_1) \Phi(\bar{\nu}_{21}, \bar{\mu}_{21}, \ell_{21}) \end{aligned} \quad (59)$$

The sums are over all boxes that influence the receiving point. An example of the sending areas for a typical receiving point on the tip is shown in Figure 13.

The source strengths on those boxes outside the interference region are determined by the relationships given in the preceding section. However, when a box is within the interference region, the strength distribution is modified to account for the interaction with other boxes. In this way, the mutual interaction effects will be properly represented in the evaluation of the velocity potentials.

In anticipation of the need to calculate velocities induced at points above or below the sending plane, we derive here the formulas for velocity components normal and parallel to the plane of disturbances. When the source strengths over the areas of sending boxes are constant, these induced velocities are

$$\bar{w}(n, m, t) = \sum_{\text{induced}} H^{U, L}(\nu, \mu) W(\bar{\nu}, \bar{\mu}, t) \quad (60)$$

and

$$\bar{v}(n, m, t) = \sum_{\text{induced}} H^{U, L}(\nu, \mu) V(\bar{\nu}, \bar{\mu}, t) \quad (61)$$

The  $W$  and  $V$  are derived in Appendix I. The sum is over all the boxes partially or wholly within the forward Mach hyperbola associated with the receiving point  $(n, m, t)$  and the strength is on the upper or lower sheet depending on the sign of  $t$ .

The source strengths at the box centers on the tip and on the portion of the wing within the interference region are corrected using Equations 44 and 45 in the following forms:

$$S_{wI}: \frac{D}{Dt} \bar{z}_1(n_1, m_1) = H^U(n_1, m_1) + \sum H^U(\nu_2, \mu_2) N(\bar{\nu}_{12}, \bar{\mu}_{12}, t_{12}) \quad (62)$$

$$S_{rI}: \frac{D}{Dt} \bar{z}_2(n_2, m_2) = H^U(n_2, m_2) + \sum H^U(\nu_1, \mu_1) N(\bar{\nu}_{21}, \bar{\mu}_{21}, t_{21}) \quad (63)$$

$$\frac{D}{Dt} \bar{z}_2(n_2, m_2) = -H^L(n_2, m_2) + \sum H^U(\nu_1, \mu_1) N(\bar{\nu}_{21}, \bar{\mu}_{21}, t_{21}) \quad (64)$$

where the velocity influence coefficient (VIC) is written

$$N(\bar{\nu}, \bar{\mu}, t) = W(\bar{\nu}, \bar{\mu}, t) \cos \gamma + V(\bar{\nu}, \bar{\mu}, t) \sin \gamma \quad (65)$$

The VIC gives the component of velocity normal to a source sheet at a Mach box center. The magnitude of the velocity is that which is induced there by an out-of-plane Mach box  $\bar{v}$  box lengths ahead,  $\bar{\mu}$  box widths to the port side, and  $l$  box widths down. The source strength at the sending box is unity and the dihedral angle between the two planes is  $\gamma$ .

The source strengths on boxes on the wing and tip diaphragms that are within the interference region must also be adjusted for interaction effects. This correction has to be done in such a way as to preserve the continuity of the normal velocity at the diaphragm centers. The perturbation velocity at the center of the diaphragm boxes is given by

$$S_{WDI}: \bar{w}(n_1, m_1) = H^U(n_1, m_1) + \sum H^L(v_2, \mu_2) N(\bar{v}_{12}, \bar{\mu}_{12}, l_{12}) \quad (66)$$

$$S_{TDI}: \bar{w}(n_2, m_2) = H^U(n_2, m_2) + \sum H^U(v_1, \mu_1) N(\bar{v}_{21}, \bar{\mu}_{21}, l_{21}) \quad (67)$$

$$\bar{w}(n_2, m_2) = -H^L(n_2, m_2) + \sum H^U(v_1, \mu_1) N(\bar{v}_{21}, \bar{\mu}_{21}, l_{21}) \quad (68)$$

To ensure continuity of the normal velocity across those diaphragm boxes that are in the interference region, substitute Equations 66 and 24 into Equation 26 to get

$$0 = H^U(n_1, m_1) + H^L(n_1, m_1) + \sum H^L(v_2, \mu_2) N(v_{12}, \mu_{12}, l_{12}) \quad (69)$$

for the wing diaphragm and substitute Equations 67 and 68 into Equation 26 to get

$$0 = H^U(n_2, m_2) + H^L(n_2, m_2) \quad (70)$$

for the tip diaphragm. There is no net interference effect across the tip diaphragm because it is completely within the wing Mach envelope, and is, therefore, exposed on both sides to the same portions of the wing and wing diaphragm.

The condition of zero jump in velocity potential across the wing and tip diaphragms can be applied to the portions in the interference region because neither is in the wake of the other. This condition can be easily written by applying the approximation of constant strength at box centers on the wing diaphragm to Equation 49. After separating out the  $n_1, m_1$  terms we get

$$\begin{aligned} \Delta H(n_1, m_1) = & - \sum \Delta H(v_1, \mu_1) \Phi(\bar{v}_{11}, \mu_{11}) / \Phi(0, 0) \\ & - \sum H^L(v_2, \mu_2) \Phi(\bar{v}_{12}, \mu_{12}, l_{12}) / \Phi(0, 0) \end{aligned} \quad (71)$$

A similar relationship can be obtained for the tip diaphragm by approximating Equation 4<sup>o</sup> and isolating the  $n_2, m_2$  terms,

$$\Delta H(n_2, m_2) = - \sum \Delta H(v_2, \mu_2) \Phi(\bar{v}_{22}, \bar{\mu}_{22}) / \Phi(0, 0) \quad (72)$$

Here again, there is no net interference effect across the tip diaphragm. This result is consistent with the result observed in Equation 70.



#### 4. COMPUTER PROGRAM

##### USE AND LIMITATIONS

An IBM 7094 Fortran IV language computer program is presented which calculates supersonic unsteady aerodynamic forces on a wide variety of wings with symmetrical folded tips. The calculation procedure is based on the source superposition method that has been extended to account for the interference effects between intersecting lifting surfaces. The Mach box approximation is employed to reduce the integral equations to sums of constant values of source strength at box centers times certain integrals dependent upon relative position, Mach number, and reduced frequency. These integrals are the aerodynamic influence coefficients which express the velocity potential influence coefficient (VPIC) or velocity influence coefficient (VIC) induced at a receiving box center by a unit strength source sheet covering the area of influence of a sending box. The VPIC's and VIC's are developed and presented in the form used for programming in Appendix I.

The configuration to be analyzed must have a plane of symmetry, wing and tip leading edges that are not swept forward, and supersonic trailing edges. The surfaces may have any small angle of attack or camber distribution or may be oscillating in an arbitrary mode of rigid or elastic vibration. Each tip may have a side edge, subsonic leading edge, and fore or aft swept trailing edge. The wing may also have a subsonic leading edge and fore or aft swept trailing edge.

The program has been specifically designed to calculate generalized forces to be used in determination of steady-state lift and moment coefficients, oscillatory stability derivatives, gust loads and aeroelastic stability. Also computed are source strength distributions corrected for interference effects and velocity potential distributions over the entire surface. The velocity potential values are used to determine the generalized forces and may also be employed to calculate surface pressure or pressure difference distributions over the configuration. For one run, these quantities are calculated for one supersonic Mach number, up to 20 reduced frequencies (steady-state loadings result when the oscillatory frequency is zero for a nonzero mode shape) and up to 10 modes of deflection.

The Mach box approximation is achieved by overlaying the surfaces and diaphragms with a grid of rectangular boxes with chordwise length  $b$  and spanwise width,  $b/\beta$ , which makes the diagonals parallel to Mach lines.

The program will subdivide the source sheets into as many as 20 chordwise and 30 spanwise boxes with the condition that box edges lie along the fold line. Past experience indicates that 12 chordwise boxes adequately defines the motion at low Mach numbers and as few as 8 boxes along the chord will in some cases describe an arbitrary vibration mode.

The method of isolating the upper and lower sides with diaphragms used in the computer program is shown in Figure 11. The program was written for a set of diaphragms that are a practical minimum in area and it was not discovered until after completion that the more general diaphragm construction shown in Figure 10 greatly simplifies the algebra and logic at only a slight increase in diaphragm area. Consequently, the equations used in the program are more complicated and will be listed here for completeness along with the pictures of the actual diaphragm regions used.

The folded tips wing and diaphragms actually used in the program are shown in Figure 14 with the zones of influence identified. The portions of the wing diaphragms outboard of the tips in the zero-degree fold position do not affect any portion of the surfaces and are, therefore, eliminated. The Mach box overlay for this configuration is shown in Figure 15. The following relative position coordinates will be used in the influence coefficients

$$\bar{v}_{11} = n_1 - v_1; \bar{\mu}_{11} = m_1 - \mu_1; \ell_{11} = 0$$

$$\bar{v}_{12} = n_1 - v_2; \bar{\mu}_{12} = (m_1 - y_{fl}) \cos \gamma - \mu_2; \ell_{12} = (m_1 - y_{fl}) \sin \gamma$$

$$\bar{v}_{21} = n_2 - v_1; \bar{\mu}_{21} = m_2 \cos \gamma - (\mu_1 - y_{fl}); \ell_{21} = -m_2 \sin \gamma$$

$$\bar{v}_{22} = n_2 - v_2; \bar{\mu}_{22} = m_2 - \mu_2; \ell_{22} = 0$$

where  $\gamma$  is the fold angle positive tips up and  $y_{fl}$  is the distance from the wing centerline to the tip foldline.

The equations for velocity potential difference are

$$S_W: \Delta\phi(n_1, m_1) = \frac{b}{\beta} \sum \Delta H(v_1, \mu_1) \Phi(\bar{v}_{11}, \bar{\mu}_{11}) \quad (73)$$

$$S_{WI}: \Delta\phi(n_1, m_1) = \frac{b}{\beta} \sum \Delta H(v_1, \mu_1) \Phi(\bar{v}_{11}, \bar{\mu}_{11}) \quad (74)$$

$$+ \frac{b}{\beta} + \sum H^U(v_2, \mu_2) \Phi(\bar{v}_{12}, \bar{\mu}_{12}, \ell_{12})$$



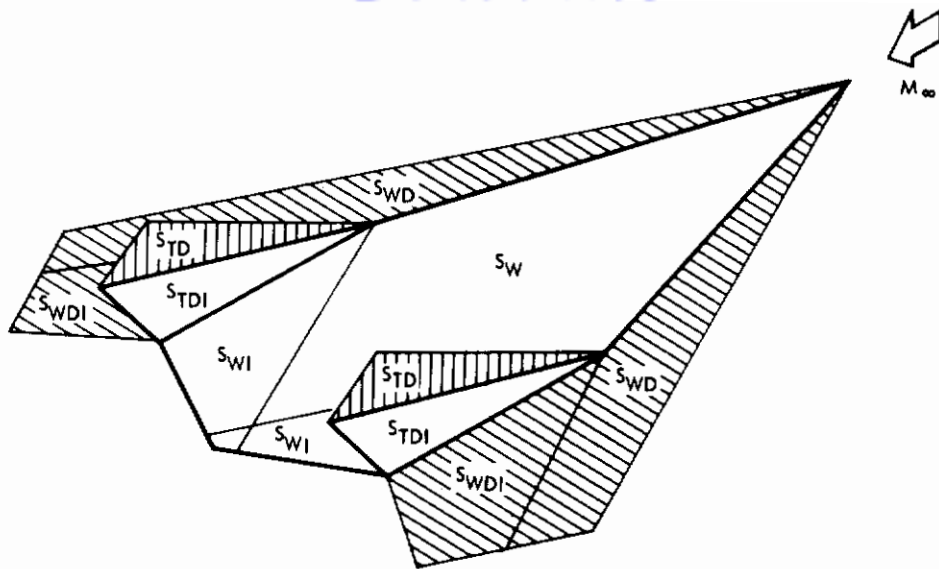


Figure 14. Zones on General Case In Program

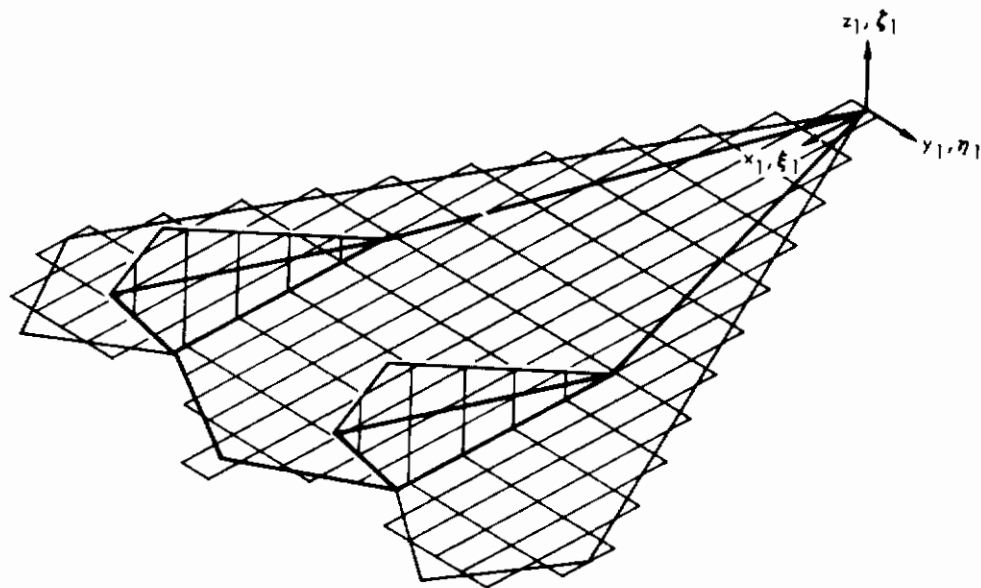


Figure 15. Boxes on General Case In Program

$$S_{TI}: \Delta \phi(n_2, m_2) = \frac{b}{\beta} \sum \Delta H(\nu_2, \mu_2) \Phi(\bar{\nu}_{22}, \bar{\mu}_{22}) \quad (75)$$

The source strengths at the surface box centers are given by

$$S_w: \frac{D}{Dt} Z(n_1, m_1) = H^U(n_1, m_1) \quad (76)$$

$$\frac{D}{Dt} Z(n_1, m_1) = -H^L(n_1, m_1) \quad (77)$$

$$S_{wI}: \frac{D}{Dt} Z(n_1, m_1) = H^U(n_1, m_1) + \sum H^U(\nu_2, \mu_2) N(\bar{\nu}_{12}, \bar{\mu}_{12}, \ell_{12}) \quad (78)$$

$$\frac{D}{Dt} Z(n_1, m_1) = -H^L(n_1, m_1) \quad (79)$$

$$S_{TI}: \frac{D}{Dt} Z(n_2, m_2) = H^U(n_2, m_2) + \sum H^U(\nu_1, \mu_1) N(\bar{\nu}_{21}, \bar{\mu}_{21}, \ell_{21}) \quad (80)$$

$$\frac{D}{Dt} Z(n_2, m_2) = -H^L(n_2, m_2) + \sum H^U(\nu_1, \mu_1) N(\bar{\nu}_{21}, \bar{\mu}_{21}, \ell_{21}) \quad (81)$$

where  $N(\bar{\nu}, \bar{\mu}, \ell) = W(\bar{\nu}, \bar{\mu}, \ell) \cos \gamma + V(\bar{\nu}, \bar{\mu}, \ell) \sin \gamma$  and the VIC's  $W$  and  $V$  are given in Appendix I.

The diaphragm source strength differences are calculated from

$$S_{wD}: \Delta H(n_1, m_1) = - \sum \Delta H(\nu_1, \mu_1) \Phi(\bar{\nu}_{11}, \bar{\mu}_{11}) / \Phi(0, 0) \quad (82)$$

$$S_{wDI}: \Delta H(n_1, m_1) = - \sum \Delta H(\nu_1, \mu_1) \Phi(\bar{\nu}_{11}, \bar{\mu}_{11}) / \Phi(0, 0) \\ - \sum H^L(\nu_2, \mu_2) \Phi(\bar{\nu}_{12}, \bar{\mu}_{12}, \ell_{12}) / \Phi(0, 0) \quad (83)$$

$$S_{TDI}: \Delta H(n_2, m_2) = - \sum \Delta H(\nu_2, \mu_2) \Phi(\bar{\nu}_{22}, \bar{\mu}_{22}) / \Phi(0, 0) \quad (84)$$

These equations reflect the slight differences between the program and the general method. The outline presented below details the logical and computational operations of the program. Also provided is a flow chart and list of

*Continued*

sequential operations that parallel the isolated surface Mach box technique. The rules for computation given in Section 2 are for the general case, whereas the outline given below is for the special case of a wing with symmetrical folded tips.

## THE LOGICAL FLOW

The computer program, MBX, consists of an executive or main program with several subroutine subprograms. The purpose of some of the subprograms is purely logical or decision making while other subprograms are developed for the many repetitive calculations necessary in this type of program. A flow chart with descriptive statements is shown in Figure 16 and will be used as reference in this outline. A complete set of program listings in Fortran IV language is presented in Appendix IV.

The main program contains all the input and output statements with their associated format statements. Tape 5 is used for input and Tape 6 is used for output in the version of the program included in the listings.

After the data arrays are initialized, the main program, MBX, reads the data describing the flight conditions, the configuration geometry, the number of frequencies and modes, and the number of boxes to be fitted in the chordwise direction. Also determined, at this point, is whether the wing and tip mode shapes are represented by polynomial coefficients or deflection patterns and the various printing options.

The geometry is converted by BOUNDS to the transformed coordinate system,  $x = X/b$ ,  $y = \beta Y/b$ ,  $z = \beta Z/b$ , and the surfaces and diaphragms are approximated with a grid of boxes based on the input maximum number in the chordwise direction. The number of boxes is adjusted so that box edges will coincide with the fold line. The subroutine then defines the inner and outer wing, tip and outer diaphragm boxes as those boxes with centers just inside the actual planform and diaphragm limits (Figure 15). If either trailing edge is not subsonic, BOUNDS then calculates the exact streamwise location of the trailing edge for each row of box centers.

MBX then writes the heading, flight conditions, and geometry as well as the outer box limits of the wing, tip, and their associated diaphragms. After the  $x_1$ ,  $y_1$  locations of the wing and tip trailing edges are printed, the list of frequencies and number of fold angle changes for each frequency are read. MBX begins the frequency loop by converting the input frequency to reduced frequencies based on the box length ( $k = b\omega/U_\infty$  and  $\bar{k} = k M_\infty^2/\beta^2$ ) and based on the root chord ( $k_r = c\omega/U_\infty$ ).

After MBX sets the value of reduced frequency the in-plane, VPIC's are computed by CAPHI and stored in a table. These coefficients depend

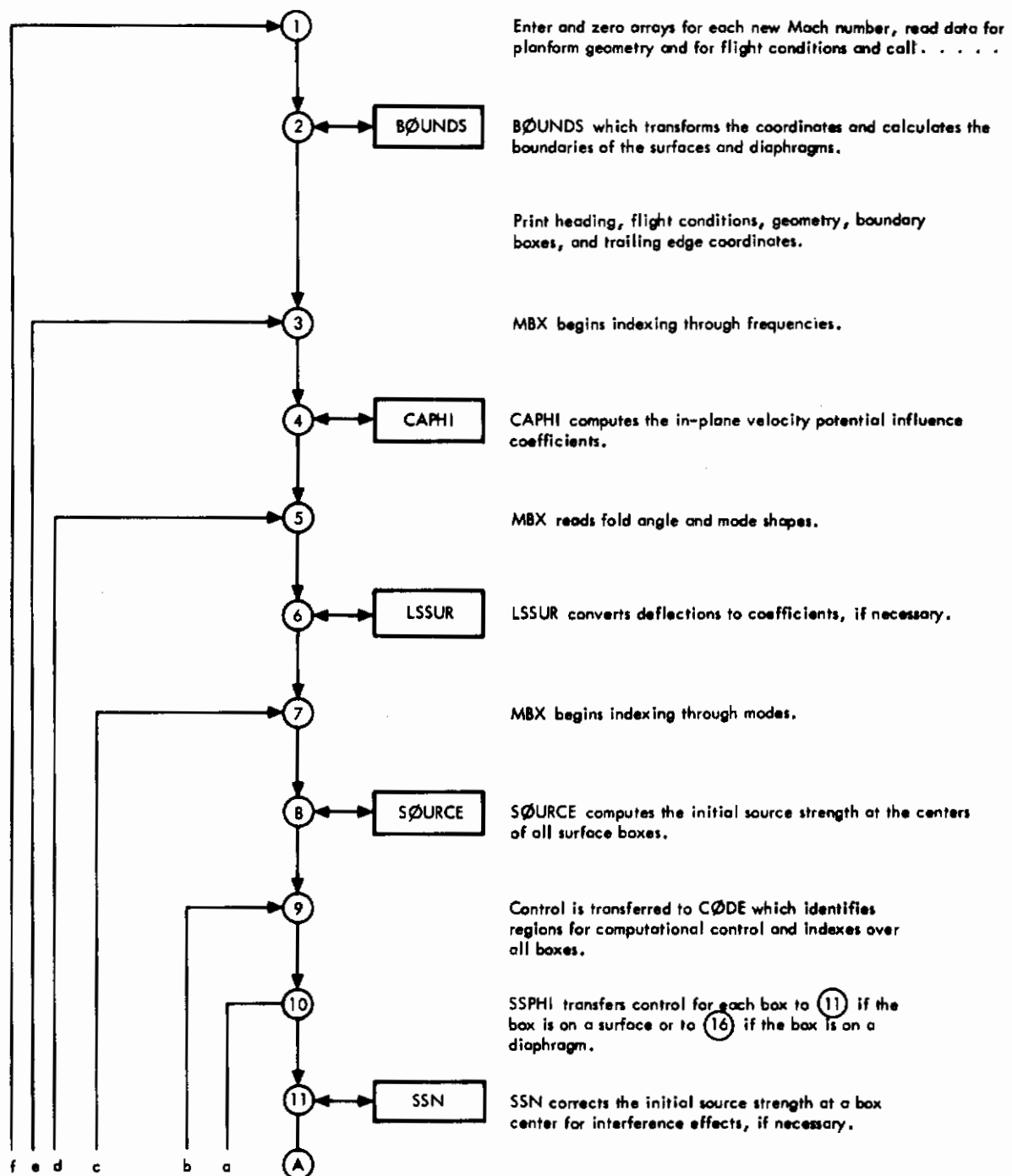


Figure 16. Flow Chart

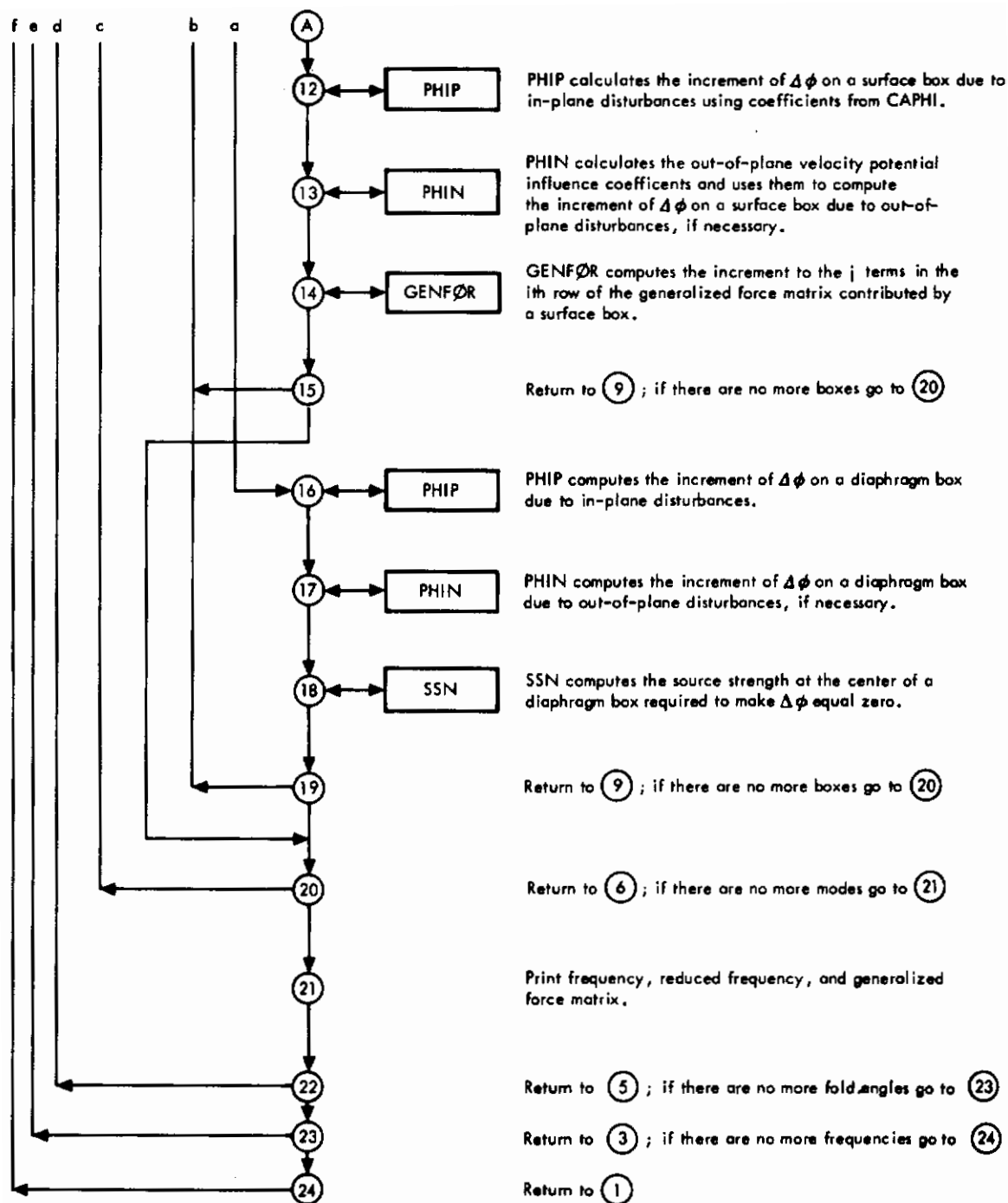


Figure 16. Flow Chart (Cont)

only upon relative position and may be computed for each Mach number-reduced frequency combination and used with all modes. These coefficients are optionally printed out by MBX.

The fold angle ( $0^\circ$  to  $90^\circ$ ) and up to ten sets of mode shape coefficients or deflections are read into the program by MBX. Previously read indicators for both the wing and tip are used to read the mode shapes before the deflections at the box centers can be determined. If the mode shapes are to be determined from the input deflection coordinates, a smoothing process is used. The input deflection data for all modes is supplied to a least squares polynomial fit subroutine (LSSUR) where the coefficients,  $A_t$ ,  $t = (r^2 + r + 2s + 2)/2$ , for a prespecified degree polynomial are determined. The coefficients for the deflection polynomials are

$$\bar{Z}_i(\xi_1, \eta_1) = \sum_{r=0}^R \sum_{s=0}^r A_t \xi_1^{r-s} \eta_1^s \quad i=1, \dots, \text{MODES} \quad (85)$$

where  $0 \leq R \leq s$  is the input value, for the degree of the wing polynomials or of the tip polynomials. The polynomial coefficients, either input or calculated are then optionally printed by MBX.

MBX sets the mode number and if the velocity potential is to be determined for this mode transfers control to SOURCE. This subroutine calculates the initial source strength values on the upper and lower sides of the surface boxes for the particular mode. This is accomplished by evaluating the polynomials and their derivatives at box centers from the coefficients.

The task of passing over all the boxes from centerline outboard and front to aft is carried out by CODE. This subroutine assigns a KODE number to each box which indicates the equations that will be used to determine the quantities at its center. With the KODE value and location of the box, SSPHI selects the proper computation subroutines depending on whether the box is on a surface or a diaphragm. The surface boxes have the source strength corrected by SSN, if necessary. The in-plane velocity potential is calculated by PHIP and, if necessary the out-of-plane velocity potential is calculated by PHIN. The increment to the mode row in the generalized force matrix due to the change in velocity potential across the box is then calculated by GENFOR.

GENFOR calculates the contribution to the generalized force integral by each box. This is accomplished by forming a separate integral for each of the deflection polynomial coefficients that consists of the potential times the appropriate powers of  $x$  and  $y$ . (Part II of this report details this method for obtaining generalized forces.) To calculate the generalized forces using the velocity potential and the displacements the trailing edge values must be determined. Since the trailing edge is generally not located along box



centers a second order interpolation procedure is employed to calculate the velocity potential at the trailing edge of each column of boxes. These values along with the previously determined x and y coordinates of the trailing edge are added to the integrals. GENFOR also weights the centerline and trailing edge box contributions to the generalized forces by the partial area of the box on the planform.

If the box is a diaphragm box, SSPHI calls PHIP and, if necessary, PHIN to calculate the velocity potential increment due to all boxes upstream that influence the box center. SSN then divides this value by  $\Phi(0, 0)$  to calculate the source strength at the diaphragm box.

When the pass over the wing is completed, there will be a set of velocity potential differences and corrected source strengths available for optional printout. Also in common is the completed jth row in the generalized force matrix.

If there are more modes the next pass over the wing will be with the same frequency but with a new mode to obtain the next row in the matrix of generalized forces. When the mode loop is complete the matrix is printed with the reduced frequency based on the root chord.

A new fold angle and modes may be read in at this point, if not, the next frequency is set by MBX and the mode loop restarted so that a new matrix may be computed from the same coefficients still in storage.

The operator must be warned that the maximum number of boxes, modes, and frequencies will take a considerable length of time to run. The surest way to save time is to cut down the number of boxes since the computational time varies with the fourth power of the number of boxes. The large number of boxes is provided so that a single critical mode and frequency may be evaluated with a very dense grid.

Input data sheets with a short explanation of each item are provided in the following text.

## INPUT AND OUTPUT

A standard input format of six 12-column fields per card is used in this program. The floating point numbers are to be to the left of the field starting in the second column and the fixed point numbers are to be to the right of the field ending in the twelfth column. Sample standard input sheets are provided in Appendix II.

The first two cards will provide the flight conditions, and planform geometry that is to remain fixed throughout the complete run. The Mach number (EMACH) and speed of sound (AS) are the free stream conditions



while the root chord (CRØØT), leading edge sweep angle (SLEW), trailing edge sweep angle (STEW), and distance from centerline to fold line (YFL) determine the wing planform. The tip geometry constants are leading edge sweep angle (SLET), trailing edge sweep angle (STET) and YFL plus the distance from the fold line to the tip line (YTIP). The length units on the speed of sound should agree with those on the chord and span distances.

The next card contains integer constants that specify the number of boxes to be fitted in the chordwise direction (NBØX), the number of frequencies (NFREQ), the number of deflection mode shapes (MØDES), and the print option indicators. The velocity potential influence coefficients (LVPIC), the initial upper and lower source strengths (LSSUL), and the final velocity potential differences with the corrected source strength (LDPHI) can be printed out by simply setting each indicator to a non zero number.

The fourth and fifth cards contain the integers that determine the mode shape input. The indicator for the method of determining the mode shapes for the wing (MDEW) and tip (MDET) is a negative integer for supplying a set of deflection points that will be smoothly fit with the specified degree polynomial or a positive integer for directly supplying the polynomial coefficients in accordance with Equation 85. The degree of the polynomial for the wing (NPØLW) and for the tip (NPØLT) need not be equal but each applies to all modes for the respective surface. MBX either reads in the correct number of coefficients computed from the degree of the polynomial or calculates the coefficients for a polynomial of that degree for each mode. If coefficients are input all modes must have enough coefficients or zeros to satisfy this requirement. If the mode shape polynomial coefficients are to be computed by the program from deflection points the number of points for the wing (NPØW) and for the tip (NPØT) is specified. Finally the coefficients, either input or calculated, can be printed out by setting the indicator for the wing (LCØW) or for the tip (LCØT) to a non zero number.

The program can be used to calculate forces for a single surface with only one leading edge and one trailing edge sweep angle by simply setting the tip indicator (MDET) equal to zero and leaving out all subsequent tip data. The tip line is then coincident with the fold line and the surface has the geometric characteristics described by the wing input data.

The next card or cards contain the lists of frequencies (ØMEGA(I)) and number of fold angles or sets of modes per frequency (ANGS(I)). There are to be two times NFREQ numbers on these cards for MBX to read a pair of numbers for each of NFREQ values. A generalized force matrix will be calculated for each ANGS(I) at the ith frequency (ØMEGA(I)) and it will contain MØDES by MØDES terms. This provision enables the user to repeat the modes or put in new modes at different frequencies without calculating all the meaningless off diagonal terms obtained when all the modes are included in one set. Also the fold angle can be varied at different frequencies and a

*Contrails*

complete set of reduced frequency, fold angle, and mode shape variations can be performed at the same Mach number during the same run .

The ensuing data cards must be carefully provided to assure successful operation of the program. There is to be one set of fold angle and mode shape data for each of the ANGS(I) and  $\Phi$ MEGA (I) combinations. Each of these sets must have the current value of the fold angle on the first card even if there are no tips to be folded.

Special care is necessary to be sure that the thickness indicator is correctly specified for both wing and tip portions of each mode shape. These numbers tell whether the coefficients are to be used to calculate the velocity potential due to a thickness distribution or due to a deflection distribution. If any of the thickness indicators is a negative number the program will use the input as a symmetric thickness distribution and calculate the potential and generalized forces accordingly.

If any of the wing thickness indicators has a zero value there will be no velocity potential calculate for either the wing or tip but the mode will be used to obtain generalized forces due to velocity potentials in other modes.

The thickness indicators for the wing and tip must be either zero or positive for all vibratory mode shapes. However, if the input frequency is zero the wing or tip may have thickness to determine the various steady state thickness effects.

If the mode shapes are to be given as a set of deflections the card or cards after the fold angle card will contain the x-y coordinates of the deflections on the wing, 3 points or 6 numbers per card until NP $\Phi$ W points have been entered, which will apply to all modes. The dimensions of the coordinates should be compatible with those of the chord and span distances input earlier.

The next cards will contain the values of the deflections at the several x-y points on the wing for all the modes, one mode at a time. Each mode should start on a new card with the thickness indicator, THW(I), as the first number and the deflections in all the following consecutive locations until NP $\Phi$ W numbers have been listed.

If there is a tip (MDET $\neq$ 0) the x-y coordinates, measured from the wing axis system, of the tip deflections will be on the cards immediately following the last of the wing deflection data. There should be 2 NP $\Phi$ T numbers entered with 6 on each card and 6 or less on the last card.

The tip deflections are to be listed on the next cards with each mode starting on a new card. The first card of the mode will have the thickness indicator, THT(I), as its first number and the first 5 deflection points in the

remaining locations. The rest of the deflections will follow, 6 points per card until NPØT points are entered.

If the mode shape coefficients are to be input rather than computed the first cards after the fold angle card will contain the wing coefficients for all the modes, one mode at a time. Each mode should begin on a new card with the first number the wing thickness indicator for that mode. The wing coefficients for each mode are entered into the next  $1/2(NPØLW+2)$   $(NPØLW+1)$  consecutive locations.

The tip coefficients, if applicable are placed on the cards immediately following the wing coefficients for the same number of modes. As before, each mode should begin on a new card with the tip thickness indicator as the first number. There are  $1/2(NPØLT+2)$   $(NPØLT+1)$  coefficients for each mode that are to be entered.

The cards must then be placed in the order presented in this discussion at the back of the program cards. Appendix II provides further explanation of the input data as well as demonstrates the use of the coefficient input option.

The output from the sample data sheets is shown in Appendix III. The first page of the output will always contain the program title, the input flight conditions and geometry, calculated fold-line and tip-line chords and reference area, number of boxes in the chordwise and spanwise directions, and the box length and width.

The next page always depicts the placement of the boxes over the wing, tip, and associated diaphragms. The box number for each row that its center inboard of the wing, tip, tip diaphragm and wing diaphragm outer edge appears as MØBW, MØBT, MØBTD, and MØBWD, respectively. The box number for each row that has its center outboard of the wing and tip inner edge appears as MIBW and MIBT, respectively.

The nondimensionalized x and y coordinates of the wing and tip trailing edge will always be on the page following the box boundaries.

If there is none of the optional printout the ensuing pages will contain the generalized forces as well as the reduced frequency, fold angle, Mach number, and chordwise boxes information.

There will be one set of generalized forces for each frequency-fold angle combination that contains the real and imaginary components as well as the magnitude and phase angle of each generalized force. Each generalized force is the pressure due to displacement in the DPHI mode weighted by the displacement in the DEFL mode.

The optional output, which will precede each table of generalized force, includes any of the following information. The input or calculated polynomial coefficients will be printed out for all modes for the wing if  $LC\emptyset W \neq 0$  and then for the tip if  $LC\emptyset T \neq 0$ . The initial upper and lower source strengths at all box centers will then be printed for all modes except those that have  $THW(I) = 0$  if  $LSSUL \neq 0$ . Finally the velocity potential difference distribution and the source strengths with interference included will be printed for each mode with  $THW(I) \neq 0$  if  $LDPHI \neq 0$ . The source strengths and velocity potential differences are printed for each box with row (N) and column (M) locations starting with (1, 1) as the foremost centerline box and proceeding outward and aft. The boxes in the plane of the wing and its diaphragm are numbered (N, M) while the boxes in the intersecting plane of the tip and its diaphragm are numbered (N, M+MMAX) where MMAX equals the most outboard box on the wing or wing diaphragm. When the tip is folded into the plane of the wing the tip surface then replaces a portion of the wing diaphragm and the tip boxes then assume the numbers of the replaced wing diaphragm boxes.

Complete listings for the program and all non-systems subroutines are presented in Appendix IV.

*Contrails*



## 5. RESULTS

### SINGLE PLANAR SURFACES

The extension of the source superposition method with Mach box approximations to intersecting surfaces is verified on the basis of results for a single planar surface. The computer program calculates potential distributions and generalized forces that compare well with exact theoretical results, other analytical methods, and experimental results. Figure 17 compares the theoretical lift and moment slopes using the exact expression with the lift curve slope calculated from the MBX program. These data are for a 65-degree delta at Mach numbers from 1.05 to 2.37. Figure 18 shows the effect of reduced frequency on the stability derivatives for a 65-degree delta at  $M = 2.0$ .

Data for a rectangular wing calculated by the characteristic box method (Reference 13) are shown in Figure 19 compared with results from the MBX program for pitch and plunge stability derivative variations with Mach number. These data are at a reduced frequency of 0.3 based on a unit root chord.

It is well known that the pressure distribution along the centerline of a delta wing at angle of attack is constant thereby implying that the velocity potential is linear. Figure 20 shows the centerline potential distribution on a 65-degree delta at  $M = 2.0$  as calculated by exact theory compared with the box center values calculated by the MBX program. The computed values oscillate around the exact line due to the jagged leading effects but the lift curve slope, obtained by integrating the velocity potential is seen to agree very well with exact theory. To obtain pressures using the MBX program one would have to fit a curve through the potential distribution by the methods described in Part II of this report. This representative function could then be differentiated to obtain the pressure distribution. Generalized forces are obtained by the MBX program by direct integration of the potential (Equation 20) which cancels the leading-edge induced oscillations.

### INTERSECTING PLANAR SURFACES

To demonstrate the effect of wing droop or fold on the stability derivatives, the tips of a 65-degree delta were folded at the 60 percent semispan line. These results are plotted in Figure 21 and 22 for  $C_{L\alpha}$  and  $C_{M\alpha}$  at various supersonic Mach numbers. The coefficients shown are based on a unit root chord and the reference area remains that of the unfolded



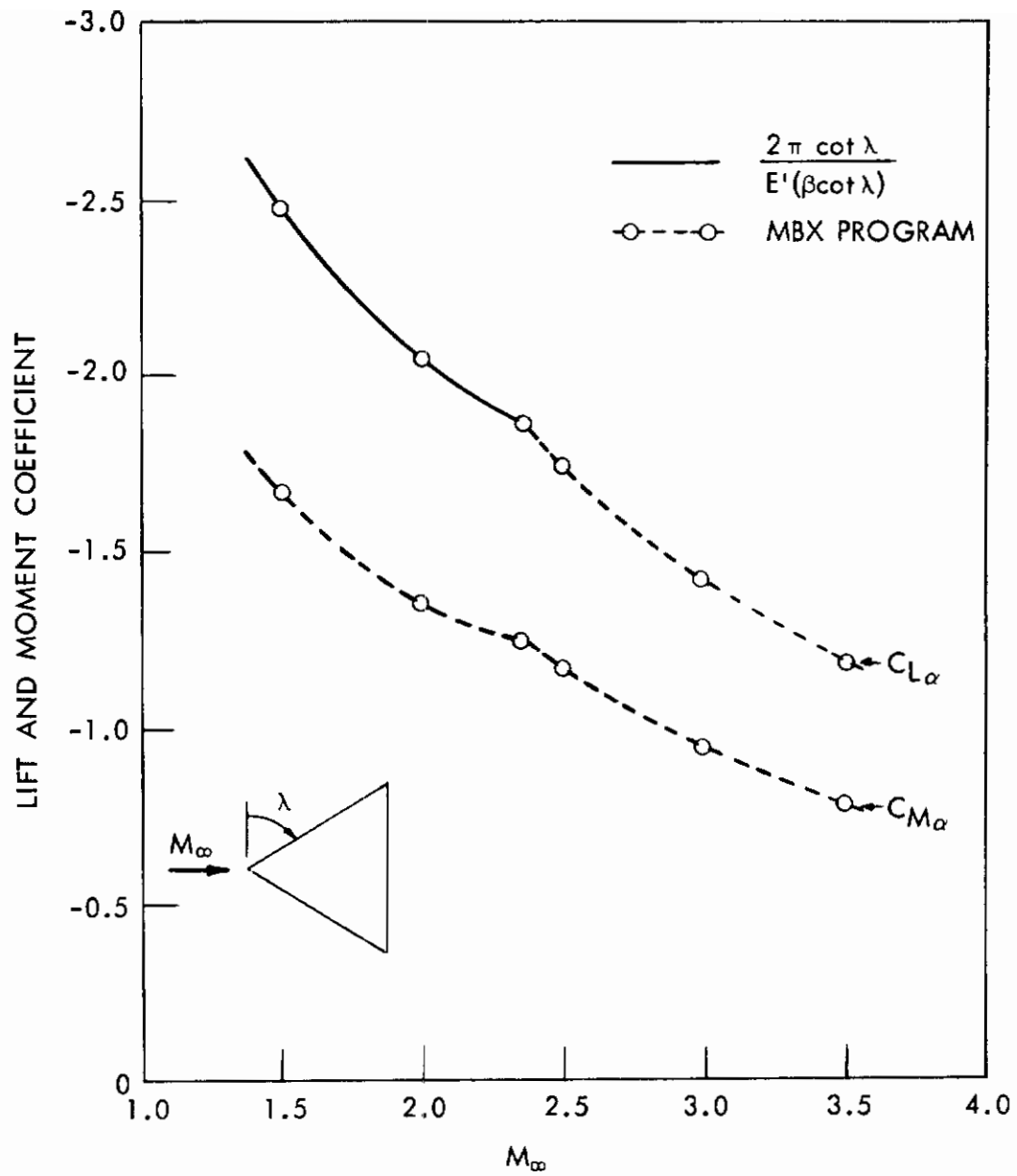


Figure 17.  $C_{M\alpha}$  and  $C_{L\alpha}$  Vs.  $M$  for  $65^\circ$  Delta

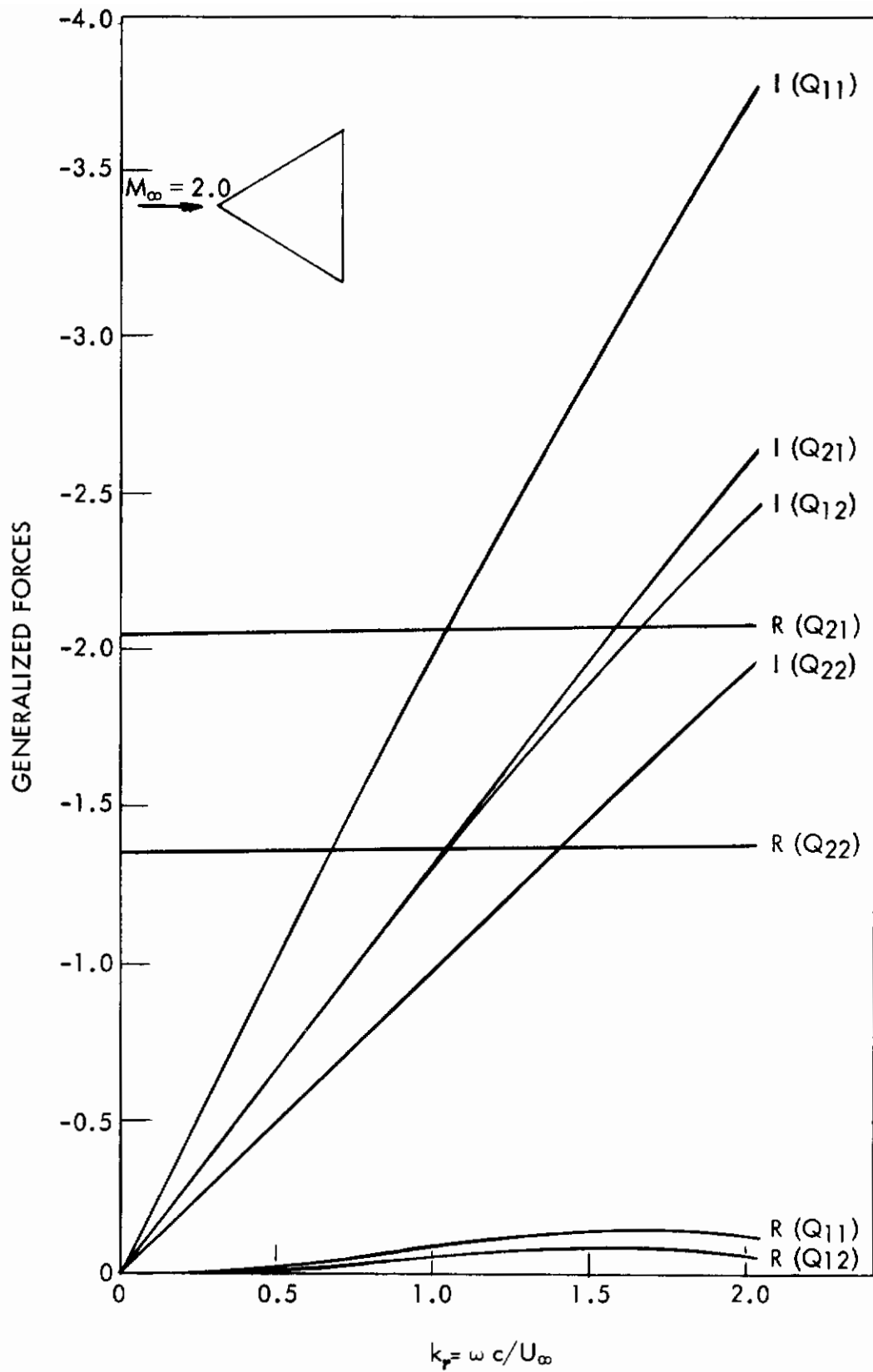


Figure 18. Forces on 65° Delta for  $Z_1 = 1.0$ ,  $Z_2 = X$

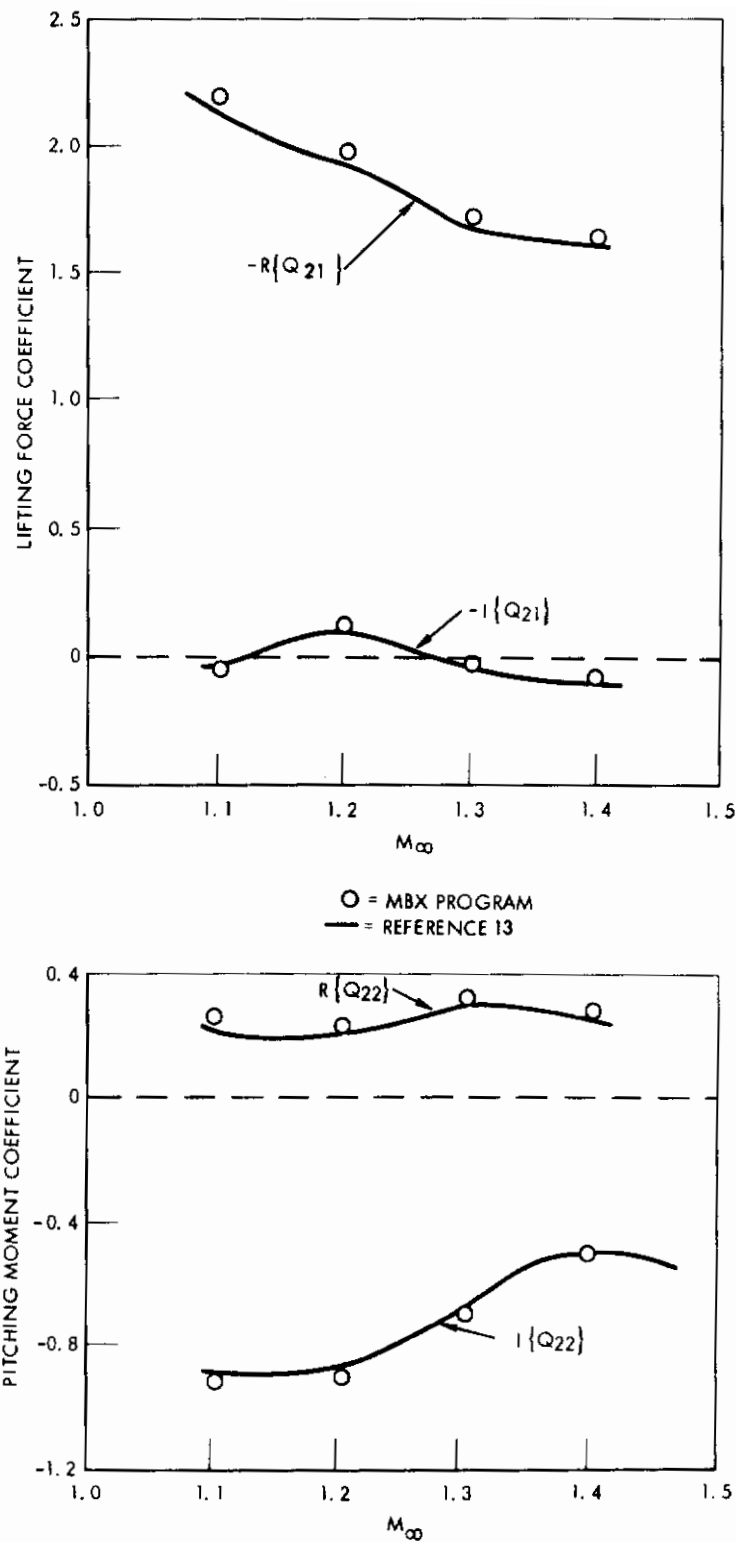


Figure 19. Aerodynamic Coefficients for a Rectangular  $AR = 2.0$  Wing Oscillating in Pitch About the Midchord With  $k = 0.3$

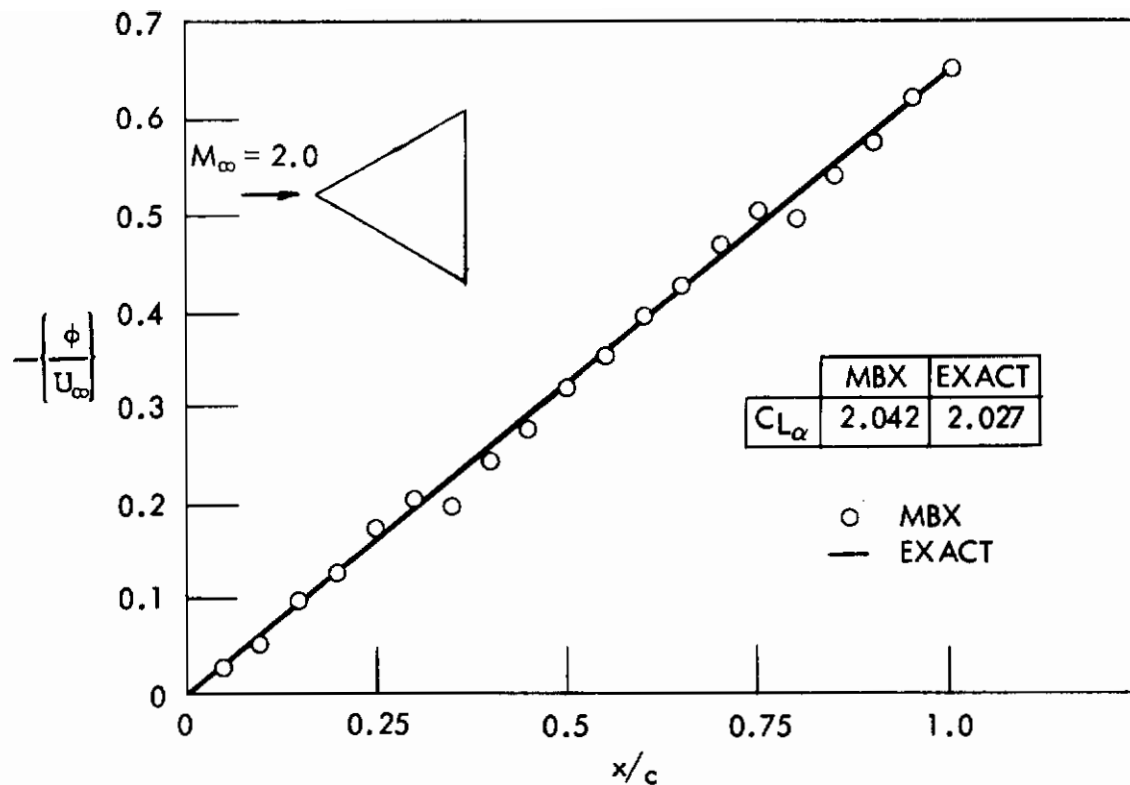


Figure 20. Potential Distribution on  $C_L$  of a  $65^\circ$  Delta at  $\alpha$

65-degree delta when the tips are taken off. Note that even without thickness the configuration has greater lift and pitching moment when the tips are folded to 90 degrees than when they are removed. These trends in lift and moment variation with fold angle are similar to the trends observed experimentally in tests on configurations that use the drooped tips in increase lift to drag ratios.

Figure 23 depicts the variation of lift and moment curve slopes,  $CL_{\alpha}$  and  $CM_{\alpha}$ , respectively, with tip fold angle,  $\gamma$ . The curve was produced by connecting the computed values at 5 degrees and above (up to 90 degrees), with the 0 degree point. The value for zero fold angle agrees with supersonic theory for planar surfaces (c.f. Figure 17). The coefficient values, which are based on the unfolded planar area, increase smoothly as the fold angle decreases from 90 to 0 degrees.

In order to obtain the smooth variation of aerodynamic force coefficients with fold angle using only a moderately dense grid of Mach boxes (up to 20 boxes chordwise), the wing diaphragm source strengths were computed using an equivalent but more accurate form of Equation 83. The expression more accurately accounts for the increasing lag in signals between the wing diaphragm and the lower side of the tip as the fold angle approaches 0 degrees.

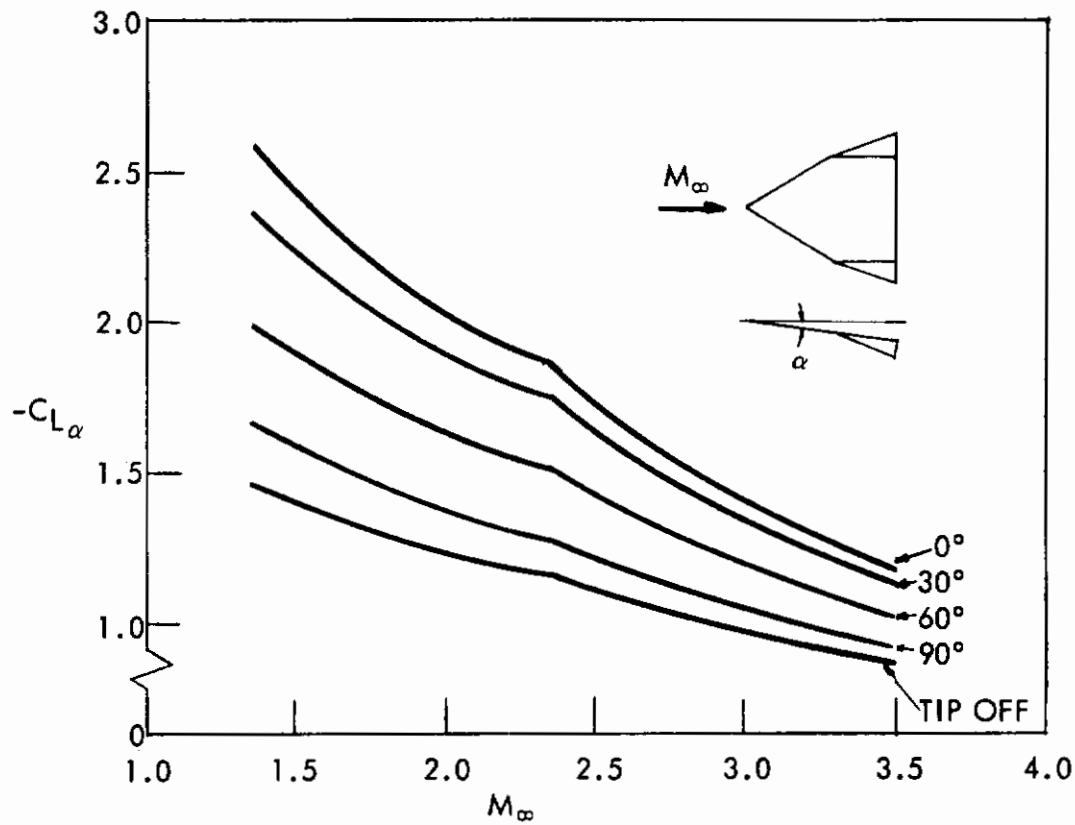


Figure 21. Variation of  $C_{L\alpha}$  on  $65^\circ$  Delta Due to Folding the Tips at the 60% Semispan

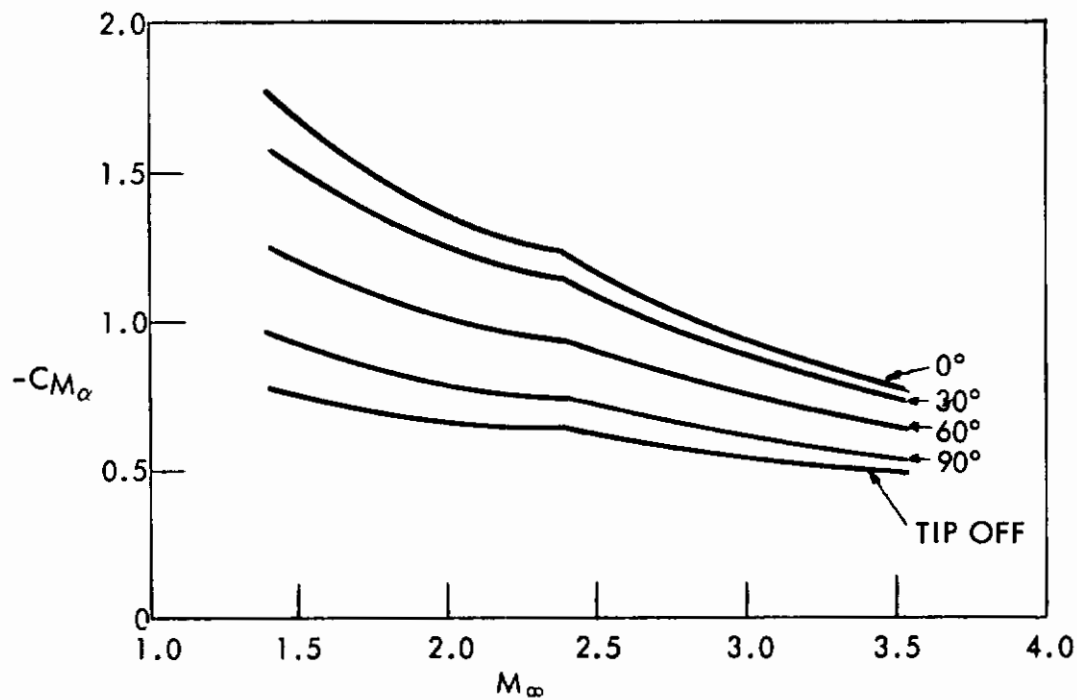


Figure 22. Variation of  $C_{M\alpha}$  on  $65^\circ$  Delta Due to Folding the Tips at the 60% Semispan

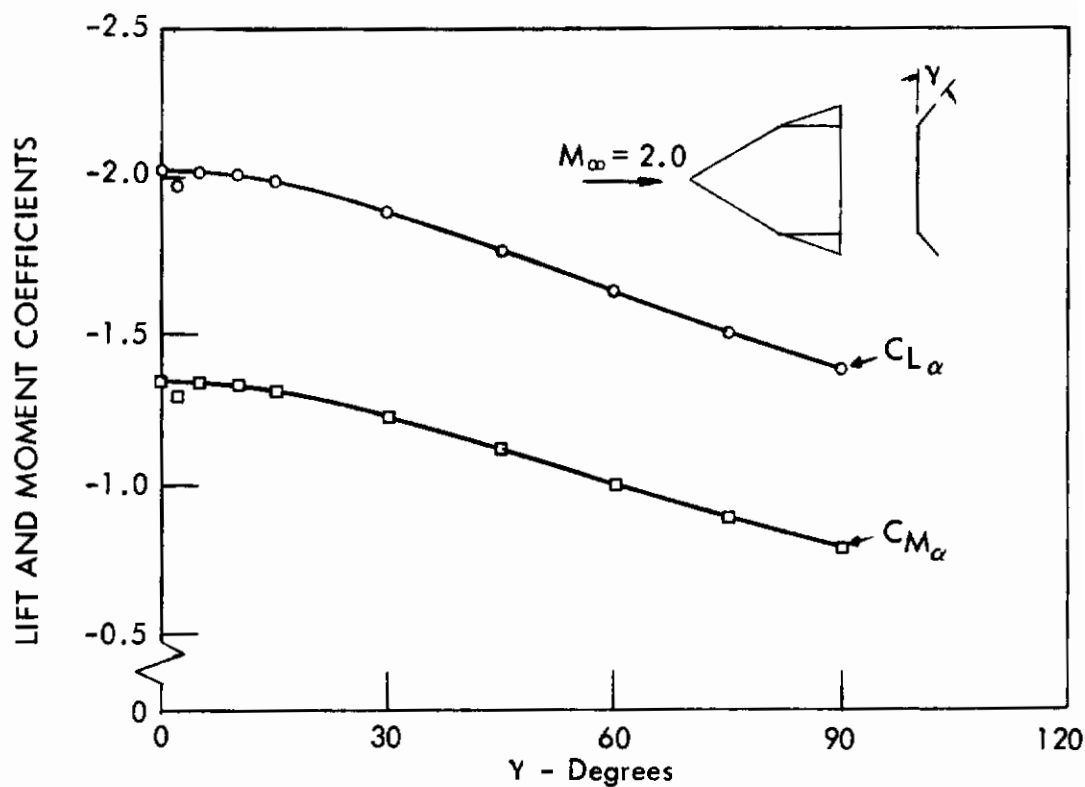


Figure 23. Variation of  $C_{L\alpha}$  and  $C_{M\alpha}$  on 65° Delta Due to Folding the Tips at the 60% Semispan

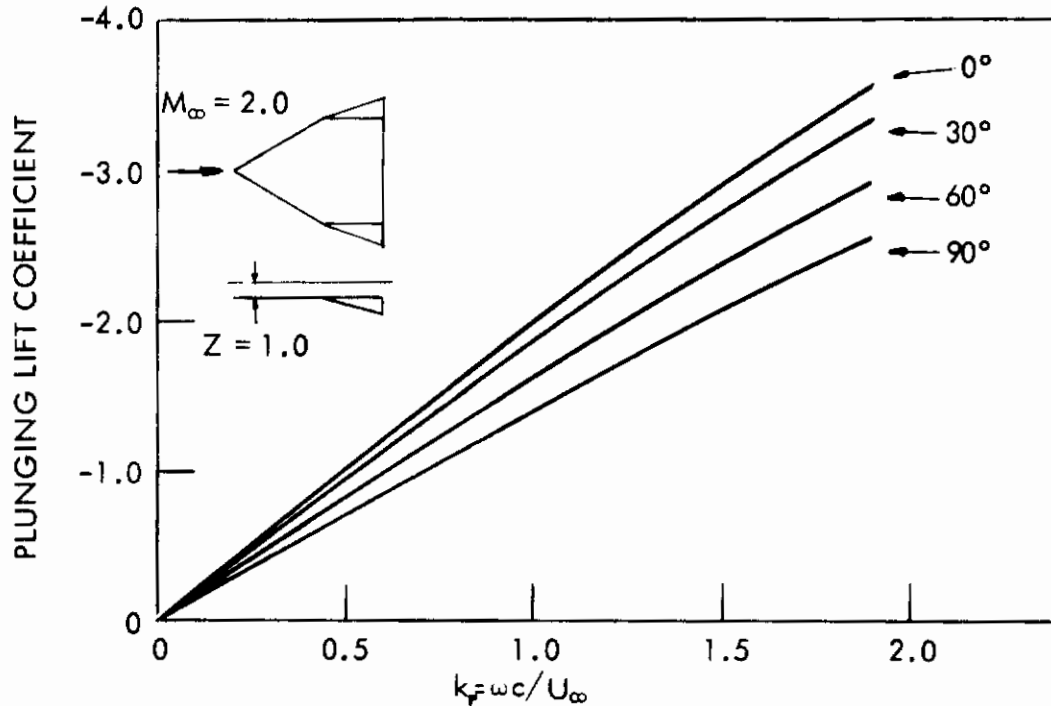


Figure 24. Variation of Plunging Lift on 65° Delta Due to Folding the Tips at the 60% Semispan



The effect of the grid size is apparent if coefficients are computed for fold angles below 5 degrees. Fortunately, one can take advantage of the fact that in the steady state case the values below 5 degrees are within 1 per cent of those for planar surfaces. Because this accuracy is not expected to be significantly different in the unsteady case, a slight modification was made to the computer program to ensure the accuracy of calculations at all values of fold angle. When the fold angle is less than 5 degrees in absolute value the program treats the configuration as if it were planar. The more complicated non-planar logic is used only if the fold angle is 5 degrees or greater in absolute value.

The program listings, Appendix IV, note the above changes at the appropriate places.

## 6. CONCLUSIONS AND RECOMMENDATIONS

As a result of this study, the practical technique for solving supersonic aerodynamic interference problems of intersecting planar surfaces using the source superposition method is now clear. Mach envelopes should be constructed for each of the separate components of the system as though the other components were nonexistent. Diaphragms should be constructed in the planes of the components to separate the upper and lower half-envelopes (Figure 10). Overlapping Mach envelopes should then be defined as interference regions. A grid of Mach boxes all of the same dimensions may then be overlaid on each of the surfaces and its diaphragm. Starting with the foremost row of boxes (rows of boxes if two or more leading edges intersect at the foremost point), the source sheets should be computed that satisfy the pertinent boundary conditions at the center of the boxes in that row. This process should then be repeated for each successive row (or rows) of boxes until the source strengths on the aftmost boxes on the surfaces have been determined.

This technique has several advantages:

1. It entirely eliminates the need for calculating source strengths on both sides of the surfaces that are in the interference region. Source strengths on one side are equal in magnitude and opposite in sense to those on the other.
2. When only the difference between upper and lower pressures on the surfaces is of concern, it eliminates the need for calculating out-of-plane velocity potentials. In most aeroelastic problems, the upper surface of an airfoil does not move relative to the lower surface, therefore, only the differences between pressures are needed.
3. The concept is simple and, therefore, provides insight which makes the technique more readily extendable. For instance, the extension that would be required to handle T-tails, V-tails, and top-mounted vertical tails is immediately apparent. Also, the extension to handle the wing-body interference problem should be a fairly simple one.

The following recommendations are made:

1. Perform an experimental study to verify the theoretical results

2. Modify the computer program to incorporate the pressure smoothing technique used in the transonic box method described in Volume II of this report
3. Modify the computer program to handle T-tails, V-tails, and top-mounted vertical tails as well as symmetrically folded tips; a single efficient computer program could probably handle all these configurations.
4. Modify the computer program to handle trailing edge control surfaces. This would involve only subdividing the boxes at the edges of the control surfaces and providing separate modal displacement functions to be used for the tangential flow conditions of those boxes on the control surfaces.
5. Apply the technique to other aerodynamic interference problems such as wing-body interference and wing-empennage interference problems
6. Apply the technique to very thick lifting surfaces

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## APPENDIX I. AERODYNAMIC INFLUENCE COEFFICIENTS

The velocity potential influence coefficient (VPIC) defines the velocity potential at a point in space due to that portion of a rectangular unit strength source sheet that lies within the upstream zone of influence from the point. The velocity induced at the point, by the rectangular sheet, is determined by differentiating the VPIC, with respect to the direction of the velocity, to obtain the velocity influence coefficient (VIC). The rectangular source sheet or box can always be positioned in the  $\zeta = 0$  plane with its length,  $b$ , parallel to the flow direction. The width of the box,  $b/\beta$ , is set so that its diagonals are parallel to Mach lines in the flow. The VPIC's and VIC's are dependent only upon the relative position of the sending area and the receiving point and the Mach number and reduced frequency parameters.

Consider a point  $(n, m, l)$  placed above a unit strength source sheet located in the  $\zeta = 0$  plane (Figure 25). The source sheet has been divided into Mach boxes whose centers lie at the points  $(v, \mu)$ . For a typical Mach number ( $M > 1$ ), the portion of the source sheet within the Mach hyperbola will influence the point  $(n, m, l)$ . This results in VPIC's and VIC's for both full and partial box areas sending to the point.

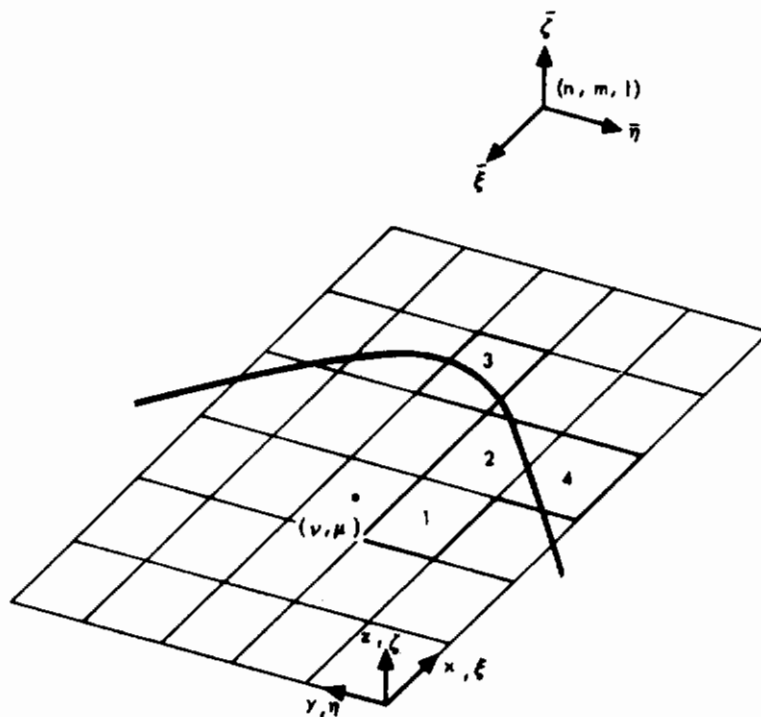


Figure 25. Boxes That Influence an Out-of-Plane Point



The expression for the VPIC is

$$\Phi(\bar{v}, \bar{\mu}, l) = - \frac{1}{\pi} \int_{\bar{\xi}_L}^{\bar{\xi}_U} \int_{\bar{\eta}_L}^{\bar{\eta}_U} \frac{e^{-i\bar{k}\bar{\xi}} \cos\left(\frac{\bar{k}}{M} \sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}\right)}{\sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}} d\bar{\eta} d\bar{\xi} \quad (86)$$

where

$$\bar{k} = \frac{\omega b}{U_\infty} \frac{M_\infty^2}{\beta^2}; \beta^2 = M_\infty^2 - 1$$

The limits of integration correspond to the area of the box that is within the Mach hyperbola.

Since the integrand is singular along the Mach hyperbola,  $\bar{\xi} = \sqrt{\bar{\eta}^2 + l^2}$  the integration of Equation 86 for the three types of intersected boxes (Figure 25) must be performed after the singularity has been removed. This may be accomplished by using the following identity

$$\frac{\cos\left(\frac{\bar{k}}{M} \sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}\right)}{\sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}} = -\frac{\partial}{\partial \bar{\eta}} \psi\left(\sqrt{\bar{\xi}^2 - l^2}, \bar{\eta}\right) \quad (87)$$

where

$$\psi(\Omega, \bar{\eta}) = J_0\left(\frac{\bar{k}}{M} \Omega\right) \sin^{-1}\left(\frac{\bar{\eta}}{\Omega}\right) + \sum_{r=1}^{\infty} \frac{(-1)^r}{r} J_{2r}\left(\frac{\bar{k}}{M} \Omega\right) \sin\left(2r \sin^{-1}\left(\frac{\bar{\eta}}{\Omega}\right)\right)$$

and  $\Omega = \sqrt{\bar{\xi}^2 - l^2}$

Upon substitution of Equation 87 into the VPIC expression the inner integration may be performed to obtain

$$\Phi = - \frac{1}{\pi} \int_{\bar{\xi}_L}^{\bar{\xi}_U} e^{-i\bar{k}\bar{\xi}} \left[ \psi(\Omega, \bar{\eta}_U) - \psi(\Omega, \bar{\eta}_L) \right] d\bar{\xi} \quad (88)$$

The integral is now in a form suitable for numerical evaluation and the values of the function,  $\psi$ , will correspond to the area within the Mach hyperbola on cut boxes if the following interpretation is adhered to,

$$-\frac{\pi}{2} < \sin^{-1} \frac{\bar{\eta}}{\Omega} < \frac{\pi}{2} \quad -\Omega < \bar{\eta} < \Omega$$

$$\sin \frac{\bar{\eta}}{\Omega} = \frac{\pi}{2} \quad \bar{\eta} \geq \Omega \quad (89)$$

$$\sin \frac{\bar{\eta}}{\Omega} = -\frac{\pi}{2} \quad \bar{\eta} \leq -\Omega$$

The VIC in the  $l$  direction is given by

$$W = -\frac{1}{\pi} \frac{\partial}{\partial l} \int_{\bar{\xi}_L}^{\bar{\xi}_U} \int_{\bar{\eta}_L}^{\bar{\eta}_U} \frac{e^{-ik\bar{\xi}} \cos \left( \frac{\bar{k}}{M} \sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2} \right)}{\sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}} d\bar{\eta} d\bar{\xi} \quad (90)$$

Substitution of the identity (Equation 87), performing the inner integration, and interchanging the differentiation with the outer integration results in

$$W = -\frac{1}{\pi} \int_{\bar{\xi}_L}^{\bar{\xi}_U} e^{-ik\bar{\xi}} \frac{\partial}{\partial l} \left[ \psi(\Omega, \bar{\eta}_U) - \psi(\Omega, \bar{\eta}_L) \right] d\bar{\xi} \quad (91)$$

Since  $l$  only appears in the parameter  $\Omega = \sqrt{\bar{\xi}^2 - l^2}$ , the differentiation variable may be changed, i.e.,

$$\frac{\partial}{\partial l} \Omega = -\frac{l}{\bar{\xi}} \frac{\partial}{\partial \bar{\xi}} \Omega$$

The expression for the  $l$  direction VIC may be integrated by parts after the variable change with the following relationships

$$\begin{aligned} u &= \frac{1}{\bar{\xi}} e^{-ik\bar{\xi}} & du &= \frac{1 + ik\bar{\xi}}{\bar{\xi}^2} e^{-ik\bar{\xi}} d\bar{\xi} \\ dv &= \frac{\partial}{\partial \bar{\xi}} \psi d\bar{\xi} & v &= \psi \end{aligned}$$

to obtain:

$$W = \frac{l}{\pi} \left\{ \left( \frac{1}{\bar{\xi}} e^{-i\bar{k}\bar{\xi}} \left[ \psi(\Omega, \bar{\eta}_U) - \psi(\Omega, \bar{\eta}_L) \right] \right) \right. \\ \left. + \int_{\bar{\xi}_L}^{\bar{\xi}_U} \frac{1 + i\bar{\xi}\bar{k}}{\bar{\xi}^2} e^{-i\bar{k}\bar{\xi}} \left[ \psi(\Omega, \bar{\eta}_U) - \psi(\Omega, \bar{\eta}_L) \right] d\bar{\xi} \right\} \quad (92)$$

The relationships in Equation 89 also apply to the function,  $\psi$ , in the  $l$ -direction VIC.

The unsteady VIC in the  $\bar{\eta}$ -direction is given by

$$V = -\frac{1}{\pi} \frac{\partial}{\partial \bar{\eta}} \int_{\bar{\xi}_L}^{\bar{\xi}_U} \int_{\bar{\eta}_L}^{\bar{\eta}_U} \frac{e^{-i\bar{k}\bar{\xi}} \cos\left(\frac{\bar{k}}{M} \sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}\right)}{\sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}} d\bar{\eta} d\bar{\xi} \quad (93)$$

where the integral becomes singular when  $\bar{\xi} = \sqrt{\bar{\eta}^2 + l^2}$  and,

$$F(\bar{\xi}, \bar{\eta}, l) = \frac{\cos\left(\frac{\bar{k}}{M} \sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}\right)}{\sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}}$$

$$V = -\frac{1}{\pi} \frac{\partial}{\partial \bar{\eta}} \int_{\bar{\xi}_L}^{\bar{\xi}_U} \int_{\bar{\eta}_L}^{\bar{\eta}_U} e^{-i\bar{k}\bar{\xi}} F(\bar{\xi}, \bar{\eta}, l) d\bar{\eta} d\bar{\xi} \quad (94)$$

The limits are not functions of  $\bar{\eta}$  therefore the differentiation can be performed inside the integral and then the inner integration performed to obtain

$$V = -\frac{1}{\pi} \int_{\bar{\xi}_L}^{\bar{\xi}_U} e^{-i\bar{k}\bar{\xi}} \left[ F(\bar{\xi}, \bar{\eta}_U, l) - F(\bar{\xi}, \bar{\eta}_L, l) \right] d\bar{\xi} \quad (95)$$

Integration by parts may be performed with the following relationships

$$\begin{aligned} u &= \frac{e^{-i\bar{k}\bar{\xi}}}{\bar{\xi}} & du &= \frac{1 + i\bar{\xi}\bar{k}}{\bar{\xi}^2} e^{-i\bar{k}\bar{\xi}} d\bar{\xi} \\ dv &= \frac{\bar{\xi} \cos\left(\frac{\bar{k}}{M} \sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}\right)}{\sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}} d\bar{\xi} & v &= \frac{M}{\bar{k}} \sin\left(\frac{\bar{k}}{M} \sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}\right) \\ & & v &= \frac{M}{\bar{k}} E(\bar{\eta}) \end{aligned}$$

to obtain for the VIC,

$$\begin{aligned} V &= -\frac{M}{\pi\bar{k}} \left\{ \left( \frac{e^{-i\bar{k}\bar{\xi}}}{\bar{\xi}} \left[ E(\eta_U) - E(\eta_L) \right] \right) \right\}_{\bar{\xi}_L}^{\bar{\xi}_U} \\ &\quad + \int_{\bar{\xi}_L}^{\bar{\xi}_U} \frac{1 + i\bar{k}\bar{\xi}}{\bar{\xi}^2} e^{-i\bar{k}\bar{\xi}} \left[ E(\bar{\eta}_U) - E(\bar{\eta}_L) \right] d\bar{\xi} \quad (96) \end{aligned}$$

To ensure proper values for cut boxes adopt the interpretation  $\sin(a) = 0$ , if  $a$  is imaginary.

For steady flow the  $\bar{\eta}$ -direction VIC is given by

$$V^S = -\frac{1}{\pi} \frac{\partial}{\partial \bar{\eta}} \int_{\bar{\xi}_L}^{\bar{\xi}_U} \int_{\bar{\eta}_L}^{\bar{\eta}_U} \frac{d\bar{\eta} d\bar{\xi}}{\sqrt{\bar{\xi}^2 - \bar{\eta}^2 - l^2}} \quad (97)$$

Interchanging the order of differentiation and integration yields

$$v^S = -\frac{1}{\pi} \int_{\xi_L}^{\xi_U} \left[ \frac{1}{\sqrt{\xi^2 - \eta_U^2 - l^2}} - \frac{1}{\sqrt{\xi^2 - \eta_L^2 - l^2}} \right] d\xi \quad (98)$$

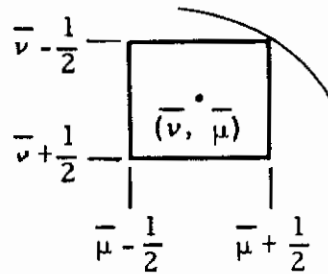
which may be further integrated to obtain

$$v^S = -\frac{1}{\pi} \left[ \cosh^{-1} \frac{\xi_U}{\sqrt{\left(\bar{\mu} + \frac{1}{2}\right)^2 + l^2}} - \cosh^{-1} \frac{\xi_L}{\sqrt{\left(\bar{\mu} + \frac{1}{2}\right)^2 + l^2}} \right. \\ \left. - \cosh^{-1} \frac{\xi_U}{\sqrt{\left(\bar{\mu} - \frac{1}{2}\right)^2 + l^2}} + \cosh^{-1} \frac{\xi_L}{\sqrt{\left(\bar{\mu} - \frac{1}{2}\right)^2 + l^2}} \right] \quad (99)$$

The proper values can be obtained from Equation 99 for the steady-state VIC for cut boxes if the following interpretation is adopted.

$$\cosh^{-1} \frac{a}{b} = 0 \text{ for } \left| \frac{a}{b} \right| \leq 1.0$$

In summary, the influence coefficients are written in their final form for an uncut box.



UNCUT BOX

$$\Omega = \sqrt{\xi^2 - l^2}$$

$$\xi_L = \bar{\nu} - \frac{1}{2}$$

$$\Phi = -\frac{1}{\pi} \int_{\xi_L}^{\bar{\nu}+1/2} e^{-ik\xi} \left[ \psi\left(\Omega, \bar{\mu} + \frac{1}{2}\right) - \psi\left(\Omega, \bar{\mu} - \frac{1}{2}\right) \right] d\xi \quad (100)$$

$$W = + \frac{l}{\pi} \left\{ \left( \frac{e^{-ik\bar{\xi}}}{\bar{\xi}} \left[ \psi\left(\Omega, \bar{\mu} + \frac{1}{2}\right) - \psi\left(\Omega, \bar{\mu} - \frac{1}{2}\right) \right] \right)_{\bar{\xi}_L}^{\bar{v} + \frac{1}{2}} + \int_{\bar{\xi}_L}^{\bar{v} + 1/2} \frac{1 - ik\bar{\xi}}{\bar{\xi}^2} e^{-ik\bar{\xi}} \left[ \psi\left(\Omega, \bar{\mu} + \frac{1}{2}\right) - \psi\left(\Omega, \bar{\mu} - \frac{1}{2}\right) \right] d\bar{\xi} \right\} \quad (101)$$

$$V = - \frac{M}{\pi k} \left\{ \left( \frac{e^{-ik\bar{\xi}}}{\bar{\xi}} \left[ E\left(\bar{\mu} + \frac{1}{2}\right) - E\left(\bar{\mu} - \frac{1}{2}\right) \right] \right)_{\bar{\xi}_L}^{\bar{v} + \frac{1}{2}} + \int_{\bar{\xi}_L}^{\bar{v} + 1/2} \frac{1 + ik\bar{\xi}}{\bar{\xi}^2} e^{-ik\bar{\xi}} \left[ E\left(\bar{\mu} + \frac{1}{2}\right) - E\left(\bar{\mu} - \frac{1}{2}\right) \right] d\bar{\xi} \right\} \quad (102)$$

$$V^S = - \frac{1}{\pi} \left[ \cosh^{-1} \frac{\bar{\xi}}{\sqrt{\left(\bar{\mu} + \frac{1}{2}\right)^2 + \ell^2}} - \cosh^{-1} \frac{\bar{\xi}}{\sqrt{\left(\bar{\mu} - \frac{1}{2}\right)^2 + \ell^2}} \right]_{\bar{\xi}_L}^{\bar{v} + \frac{1}{2}} \quad (103)$$

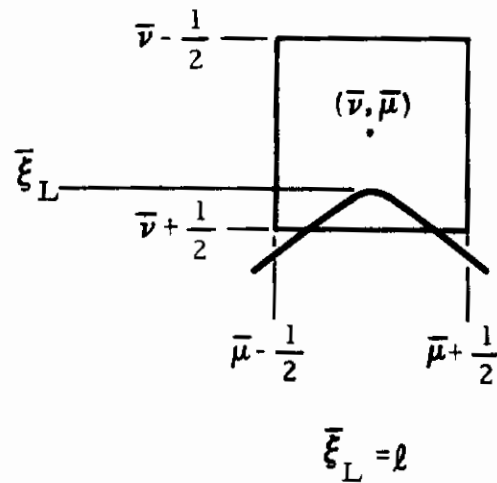
Equations 100, 101, 102, and 103 apply to all possible influence coefficients relating the velocity potential and its vertical and horizontal derivatives at a receiving point due to the area of a unit strength source box within the fore Mach cone from that point. The majority of those boxes have either no cut or at least one of the streamwise sides cut.

When one of the boxes contains the apex point of the Mach hyperbola, the lower limit of integration in Equations 100, 101 and 102 becomes the relative height,  $\ell$ , instead of either the box trailing edge or the most aft intersection point. The special treatment of the transcendental functions  $\sin(a)$ ,  $\sin^{-1}(a)$ , and  $\cosh^{-1}(a)$  when  $a$  is an invalid argument still applies to the apex box.



# Contrails

The influence coefficients for this box are somewhat simplified by the fact that the hyperbola does not intersect either streamwise side. The box and its influence coefficients are given below.



$$\Phi = - \int_{-1/2}^{1/2} e^{-i\bar{k}\bar{\xi}} J_0 \left( \frac{\bar{k}}{M} \sqrt{\bar{\xi}^2 - l^2} \right) d\bar{\xi} \quad (104)$$

$$W = l \left\{ \frac{e^{-i\bar{k}(\bar{\nu} + 1/2)}}{\bar{\nu} + 1/2} J_0 \left( \frac{\bar{k}}{M} \sqrt{(\bar{\nu} + 1/2)^2 - l^2} \right) \right. \quad (105)$$

$$\left. + \int_{-1/2}^{1/2} \frac{1 + i\bar{k}\bar{\xi}}{\bar{\xi}^2} e^{-i\bar{k}\bar{\xi}} J_0 \left( \frac{\bar{k}}{M} \sqrt{\bar{\xi}^2 - l^2} \right) d\bar{\xi} \right\} \quad (106)$$

In the limit as  $l \rightarrow 0$  these influence coefficients become those for the planar case of a box affecting itself. The coefficients are:

$$\Phi(0, 0) = - \int_0^{1/2} e^{-i\bar{k}\bar{\xi}} J_0 \left( \frac{\bar{k}}{M} \bar{\xi} \right) d\bar{\xi} \quad (107)$$

$$W(0, 0) = 1.0 \quad (108)$$

$$V(0, 0) = 0 \quad (109)$$

The expressions for the other planar influence coefficients can be obtained by setting  $l = 0$  in Equations 100, 101, 102, and 103.

The influence coefficients are evaluated numerically by a 5-point Gaussian quadrature technique similar to the method described in Part II of this report.

**APPENDIX II. SAMPLE INPUT FOR MACH BOX INTERFERENCE PROGRAM**

# FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. PROGRAMMER DATE PAGE of JOB NO.

NUMBER	IDENTIFICATION	DESCRIPTION	DO NOT KEY PUNCH
1		EMACH	- MACH NUMBER 1.0 < M <sub>0</sub> ≤ 3.5
13		AS	- SPEED OF SOUND L/T
25		CRØØT	- ROOT CHORD L
37			
49	73		
61	80		
1		SLEW	- WING L.E. SWEEP ANGLE 0° ≤ A < 90°
13		STEW	- WING T.E. SWEEP ANGLE -COS <sup>-1</sup> $\frac{1}{M}$ ≤ A ≤ COS <sup>-1</sup> $\frac{1}{M}$
25		YFL	- ROOT CHORD TO FOLD LINE - L
37		SLET	- TIP L.E. SWEEP ANGLE 0° ≤ A ≤ 90°
49	73	STET	- TIP T.E. SWEEP ANGLE -COS <sup>-1</sup> $\frac{1}{M}$ ≤ A ≤ COS <sup>-1</sup> $\frac{1}{M}$
61		YTIP	- ROOT CHORD TO TIP LINE - L
1		NBØX	- NUMBER OF CHORDWISE BOXES MAX 20
13		NFREQ	- NUMBER OF INPUT FREQUENCIES MAX 20
25		MØDES	- NUMBER OF MODE SHAPES MAX 10
37		LVØIC	- PRINT VPIC'S 0 - NO 1 - YES
49	73	LSSUL	- PRINT INITIAL H <sub>U</sub> , H <sub>L</sub> 0 - NO 1 - YES
61		LDØHI	- PRINT FINAL ΔØ'S AND H <sub>U</sub> , H <sub>L</sub> 0 - NO 1 - YES
1		MDEW	- WING MODE SHAPE INPUT 1 - DEFLECTIONS 1 - COEFFICIENTS
13		NPØLW	- INPUT OR FITTED POLYNOMIAL DEGREE MAX 5
25		NPØW	- NUMBER OF DEFLECTION POINTS MAX 100 0 - NO
37		LCØW	- PRINT POLYNOMIAL COEFFICIENTS 1 - YES
49	73		
61	80		

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# FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. PROGRAMMER DATE PAGE of JOB NO.

NUMBER	IDENTIFICATION	DESCRIPTION	DO NOT KEY PUNCH
1		MDET	TIP MODE SHAPE INPUT - 1 = DEFLECTIONS - 0 FOR NO TIP 1 = COEFFICIENTS
13		NPOLT	- INPUT OR FITTED POLYNOMIAL DEGREE MAX 5
25		NPOW	- NUMBER OF DEFLECTION POINTS MAX 100 0 = NO
37		LCOT	- PRINT POLYNOMIAL COEFFICIENTS 1 = YES
49	73		
61	80		
1		OMEGA (1)	ith FREQUENCY 1 = N FREQ - CPS
13		ANGS (1)	NUMBER OF FOLD ANGLES/ith FREQUENCY
25		OMEGA (1)	
37		ANGS (1)	
49	73		
61	80		
1		FANG	FOLD ANGLE FOR SET 1 0° ≤ γ ≤ 90°
13			CORRESPONDING TO OMEGA (1)
25			AND FIRST OF ANGS (1) VALUES.
37			
49	73		*This card must precede mode shape data even if single planar surface is being run.
61	80		
1		THW(1)	- 1.0 (THICKNESS MODE), 0.0 (NO DPHI), 1.0 (NONTHICKNESS)
13		COW (1,1)	COEFFICIENTS FOR WING MODE 1
25		COW (1,2)	
37		COW (1,3)	
49	73		
61	80		

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# FORTRAN FIXED ILC DIGIT DECIMAL DATA

DECK NO. _____		PROGRAMMER _____	DATE _____	PAGE _____ of _____	JOB NO. _____
NUMBER	IDENTIFICATION	DESCRIPTION	DO NOT KEY PUNCH		
1	1 . . 0	THW (I)	-1.0 (THICKNESS MODE), 0.0 (NO DPHI),		
13	0 . . 0		1.0 (NONTHICKNESS)		
25	1 . . 0	COW (2,1)	COEFFICIENTS FOR WING MODE I		
37	0 . . 0	COW (2,2)			
49	0 . . 0	COW (1,J)			
61					
1	1 . . 0	THT (1)	-1.0 (THICKNESS), 0.0 OR 1.0 (NONTHICKNESS MODE)		
13	1 . . 0	COT (1,1)	COEFFICIENTS FOR TIP MODE I		
25	0 . . 0	COT (1,2)			
37	0 . . 0	COT (1,3)			
49					
61					
1	1 . . 0	THT (I)	-1.0 (THICKNESS), 0.0 OR 1.0 (NONTHICKNESS MODE)		
13	0 . . 0	COT (2,1)	COEFFICIENTS FOR TIP MODE I		
25	1 . . 0	COT (2,2)			
37	0 . . 0	COT (1,J)			
49					
61					
1	3 0 . . 0	FANG	FOLD ANGLE FOR SET I 0° ≤ Y ≤ 90°		
13			CORRESPONDING TO ØMEGA (I)		
25			AND NEXT OF ANG (I) VALUES.		
37					
49					
61					

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# FORTRAN FIXED IO DIGIT DECIMAL DATA

DECK NO. \_\_\_\_\_ PROGRAMMER \_\_\_\_\_ DATE \_\_\_\_\_ PAGE \_\_\_\_\_ of \_\_\_\_\_ JOB NO. \_\_\_\_\_

NUMBER	IDENTIFICATION	DESCRIPTION	DO NOT KEY PUNCH
1 0 0	73 80	THW (1)	
13 1 0		CØW (1, 1)	
25 0 0		CØW (1, 2)	
37 0 0		CØW (1, 3)	
49 0 0	73 80		
61			
1 1 0	73 80	THW (1)	
13 0 0		CØW (2, 1)	
25 1 0		CØW (2, 2)	
37 0 0		CØW (1, J)	
49 0 0	73 80		
61			
1 0 0	73 80	THT (1)	
13 0 8 6 6 3		CØT (1, 1)	
25 0 0		CØT (1, 2)	
37 0 0		CØT (1, 3)	
49 0 0	73 80		
61			
1 1 0	73 80	THT (1)	
13 0 0		CØT (2, 1)	
25 0 8 6 6 3		CØT (2, 2)	
37 0 0		CØT (1, J)	
49 0 0	73 80		
61			

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# FORTRAN FIXED 16 DIGIT DECIMAL DATA

DECK NO. _____		PROGRAMMER _____	DATE _____	PAGE _____ of _____	JOB NO. _____
NUMBER		IDENTIFICATION	DESCRIPTION	DO NOT KEY PUNCH	
1	3 0 . 0		FANG		
13					
25					
37					
49		73		80	
61					
1	1 . 0		THW (1)		
13	1 . 0		CØW (1, 1)		
25	0 . 0		CØW (1, 2)		
37	0 . 0		CØW (1, 3)		
49		73		80	
61					
1	1 . 0		THW (1)		
13	0 . 0		CØW (2, 1)		
25	1 . 0		CØW (2, 2)		
37	0 . 0		CØW (1, J)		
49		73		80	
61					
1	1 . 0		THT (1)		
13	0 . 8 6 6 3		CØT (1, 1)		
25	0 . 0		CØT (1, 2)		
37	0 . 0		CØT (1, 3)		
49		73		80	
61					

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FORTRAN FIXED IO DIGIT DECIMAL DATA

DECK NO. \_\_\_\_\_ PROGRAMMER \_\_\_\_\_ DATE \_\_\_\_\_ PAGE \_\_\_\_\_ of \_\_\_\_\_ JOB NO. \_\_\_\_\_

NUMBER	IDENTIFICATION	DESCRIPTION	DO NOT KEY PUNCH
1 1 . . 0		THT (I)	
13 0 . . 0		COT (2, 1)	
25 0 . 8 6 6 3		COT (2, 2)	
37 0 . . 0		COT (I, J)	
49	73		80
61			
1			
13			
25			
37			
49	73		80
61			
1			
13			
25			
37			
49	73		80
61			
1			
13			
25			
37			
49	73		80
61			
1			
13			
25			
37			
49	73		80
61			

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*Contrails*

APPENDIX III. SAMPLE OUTPUT FOR MACH BOX INTERFERENCE PROGRAM

FORTAN IV PROGRAM FOR CALCULATING AERODYNAMICS OF INTERSECTING  
 SUPERSONIC LIFTING SURFACES BY THE SOURCE SUPERPOSITION METHOD.  
 THE MACH BOX TECHNIQUE IS APPLIED TO A WING WITH FOLDED TIPS.

FLIGHT CONDITIONS AND GEOMETRY

MACH NUMBER	SPEED OF SOUND	ROOT CHORD	REFERENCE AREA
2.00000	628.318 L/T	1.000 L	0.466 L**2
WING L.E. SWEEP	WING T.E. SWEEP	FOLD LINE SPAN	FOLD LINE CHORD
65.000 DEG	0. DEG	0.280 L	0.400 L
TIP L.E. SWEEP	TIP T.E. SWEEP	TIP LINE SPAN	TIP LINE CHORD
65.000 DEG	0. DEG	0.466 L	-0.000 L
CHORDWISE BOXES	BOX CHORD	SPANWISE BOXES	BOX SPAN
9	0.1077 L	7	0.0622 L

INNER AND OUTER WING AND TIP BOUNDARY BOXES

MIBW	MCBW	MIBT	MGBT	MCBTD	MCBWD
1	1	0	0	0	1
1	2	0	0	0	2
1	3	0	0	0	3
1	3	0	0	0	4
1	4	0	0	0	5
1	5	0	0	0	6
1	5	0	0	0	7
1	5	6	6	6	8
1	5	6	7	7	7
1	5	6	7	7	7



TRAILING EDGE COORDINATES

X/BGXL	Y/BGXW
9.286	0.
9.286	1.000
9.286	2.000
9.286	3.000
9.286	4.000
9.286	5.000
9.286	6.000

VELOCITY POTENTIAL DIFFERENCES (PHI UPPER - PHI LOWER) AND  
SOURCE STRENGTHS WITH INTERFERENCE EFFECTS FOR MODE 2

N	M	R DPHI(N,M)	I DPHI(N,M)	R SSU(N,M)	I SSU(N,M)	R SSL(N,M)	I SSL(N,M)
1	1	-6.2175E-02	-0.	1.0000E 00	0.	-1.0000E 00	-0.
2	1	-1.1002E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
2	2	-1.0042E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
3	1	-2.0699E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
3	2	-1.7366E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
3	3	-1.2699E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
4	1	-2.7113E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
4	2	-2.7322E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
4	3	-2.1979E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
4	4	0.	0.	-1.3901E 00	-0.	1.3901E 00	-0.
5	1	-3.6280E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
5	2	-3.4188E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
5	3	-2.3463E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
5	4	-1.4357E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
5	5	0.	0.	-2.2053E-01	-0.	2.2053E-01	-0.
6	1	-4.3055E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	2	-3.7060E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	3	-3.3971E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	4	-2.7576E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	5	-1.7214E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	6	0.	0.	-1.8792E-01	-0.	1.8792E-01	-0.
7	1	-4.1659E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	2	-4.5710E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	3	-4.2856E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	4	-3.8843E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	5	-3.0709E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	6	-1.8649E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	7	0.	0.	-1.2195E-01	-0.	1.2195E-01	-0.
7	8	0.	0.	0.	0.	0.	0.
7	9	0.	0.	0.	0.	0.	0.

VELOCITY POTENTIAL DIFFERENCES (PHI UPPER - PHI LOWER) AND  
SOURCE STRENGTHS WITH INTERFERENCE EFFECTS FOR MODE 2

N	M	R DPHI(N,M)	I DPHI(N,M)	R SSU(N,M)	I SSU(N,M)	R SSL(N,M)	I SSL(N,M)
8	1	-5.1052E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	2	-4.9452E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	3	-5.1930E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	4	-4.7656E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	5	-4.2609E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	6	-3.3582E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	7	-2.0125E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	8	0.	0.	-9.1364E-02	-0.	9.1364E-02	-0.
8	9	0.	0.	0.	0.	0.	0.
8	10	0.	0.	0.	0.	0.	0.
9	1	-5.8374E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	2	-5.8841E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	3	-5.5946E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	4	-5.7048E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	5	-5.2035E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	6	-4.6143E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	7	-3.6219E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	8	0.	0.	0.	0.	0.	0.
9	9	0.	0.	0.	0.	0.	0.
9	10	0.	0.	0.	0.	0.	0.

Contrails

MACH BOX INTERFERENCE PROGRAM OUTPUT

OSCILLATORY FREQUENCY (CPS) 0.  
 REDUCED FREQUENCY (BOX LENGTH) 0.  
 REDUCED FREQUENCY (ROOT CHORD) 0.  
 9 BOXES IN CHORD DIRECTION  
 FREE STREAM MACH NUMBER 2.000  
 TIP FOLD ANGLE (DEGREES) 0.

GENERALIZED FORCES

DPHI	DEFL	REAL PART	IMAG PART	ABS VALUE	PHASE ANGLE
1	1	-0.	-0.	0.	0. DEG
1	2	-0.	-0.	0.	0. DEG
2	1	-2.02671E 00	-0.	2.02671E 00	180.000 DEG
2	2	-1.30922E 00	-0.	1.30922E 00	180.000 DEG

Contrails

VELOCITY POTENTIAL DIFFERENCES (PHI UPPER - PHI LOWER) AND  
SOURCE STRENGTHS WITH INTERFERENCE EFFECTS FOR MODE 2

N	M	R DPHI(N,M)	I DPHI(N,M)	R SSU(N,M)	I SSU(N,M)	R SSL(N,M)	I SSL(N,M)
1	1	-6.2175E-02	-0.	1.0000E 00	0.	-1.0000E 00	-0.
2	1	-1.1002E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
2	2	-1.0042E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
3	1	-2.0699E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
3	2	-1.7366E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
3	3	-1.2699E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
4	1	-2.7113E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
4	2	-2.7322E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
4	3	-2.1979E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
4	4	0.	0.	-1.3901E 00	-0.	1.3901E 00	-0.
5	1	-3.6280E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
5	2	-3.4188E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
5	3	-2.3463E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
5	4	-1.4357E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
5	5	0.	0.	-2.2053E-01	-0.	2.2053E-01	-0.
6	1	-4.3055E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	2	-3.7060E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	3	-3.3971E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	4	-2.7576E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	5	-1.7214E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
6	6	0.	0.	-1.8792E-01	-0.	1.8792E-01	-0.
7	1	-4.1659E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	2	-4.5710E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	3	-4.2856E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	4	-3.8843E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	5	-3.0709E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
7	6	0.	0.	-1.9995E 00	-0.	-6.2407E-01	-0.
7	7	0.	0.	-1.2195E-01	-0.	1.2195E-01	-0.
7	8	0.	0.	0.	0.	0.	0.
7	9	-1.8835E-01	-0.	3.0294E 00	0.	-3.0294E 00	-0.

Contrails

VELOCITY POTENTIAL DIFFERENCES (PHI UPPER - PHI LOWER) AND  
SOURCE STRENGTHS WITH INTERFERENCE EFFECTS FOR MODE 2

N	M	R DPHI(N,M)	I DPHI(N,M)	R SSU(N,M)	I SSU(N,M)	R SSL(N,M)	I SSL(N,M)
8	1	-5.1052E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	2	-4.9452E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	3	-5.1930E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	4	-4.7656E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
8	5	-4.3328E-01	-0.	1.3477E 00	0.	-1.0000E 00	-0.
8	6	0.	0.	-2.0928E 00	-0.	3.9366E-01	-0.
8	7	0.	0.	-3.9155E-01	-0.	4.0362E-02	-0.
8	8	0.	0.	-9.1364E-02	-0.	9.1364E-02	-0.
8	9	-3.1177E-01	-0.	2.6830E 00	0.	-2.6830E 00	-0.
8	10	-2.0306E-01	-0.	1.4022E 00	0.	-1.4022E 00	0.
9	1	-5.8374E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	2	-5.8841E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	3	-5.5946E-01	-0.	1.0000E 00	0.	-1.0000E 00	-0.
9	4	-5.7903E-01	-0.	1.1121E 00	0.	-1.0000E 00	-0.
9	5	-4.9924E-01	-0.	1.2430E 00	0.	-1.0000E 00	0.
9	6	0.	0.	-2.1980E 00	-0.	-1.4318E-01	-0.
9	7	0.	0.	-6.2728E-01	-0.	-2.0469E-01	0.
9	8	0.	0.	0.	0.	0.	0.
9	9	-4.1930E-01	-0.	2.8195E 00	0.	-2.8195E 00	-0.
9	10	-3.3742E-01	-0.	1.4604E 00	0.	-1.4604E 00	0.



Contrails

MACH BOX INTERFERENCE PROGRAM OUTPUT

	OSCILLATORY FREQUENCY (CPS)	0.	9 BOXES IN CHORD DIRECTION
	REDUCED FREQUENCY (BOX LENGTH)	0.	FREE STREAM MACH NUMBER 2.000
	REDUCED FREQUENCY (ROOT CHORD)	0.	TIP FOLD ANGLE (DEGREES) 30.000

GENERALIZED FORCES

OPHI	DEFL	REAL PART	IMAG PART	ABS VALUE	PHASE ANGLE
1	1	-0.	-0.	0.	0. DEG
1	2	-0.	-0.	0.	0. DEG
2	1	-1.87613E 00	0.	1.87613E 00	180.000 DEG
2	2	-1.17817E 00	0.	1.17817E 00	180.000 DEG

VELOCITY POTENTIAL DIFFERENCES (PHI UPPER - PHI LOWER) AND  
SOURCE STRENGTHS WITH INTERFERENCE EFFECTS FOR MODE 1

N	M	R DPHI(N,M)	I DPHI(N,M)	R SSU(N,M)	I SSU(N,M)	R SSL(N,M)	I SSL(N,M)
1	1	-5.5791E-04	-3.1081E-02	0.	5.0000E-01	-0.	-5.0000E-01
2	1	-2.0870E-03	-5.4939E-02	0.	5.0000E-01	-0.	-5.0000E-01
2	2	-2.0219E-03	-5.0141E-02	0.	5.0000E-01	-0.	-5.0000E-01
3	1	-6.4435E-03	-1.0316E-01	0.	5.0000E-01	-0.	-5.0000E-01
3	2	-5.2994E-03	-8.6545E-02	0.	5.0000E-01	-0.	-5.0000E-01
3	3	-4.0007E-03	-6.3264E-02	0.	5.0000E-01	-0.	-5.0000E-01
4	1	-1.1349E-02	-1.3476E-01	0.	5.0000E-01	-0.	-5.0000E-01
4	2	-1.1673E-02	-1.3578E-01	0.	5.0000E-01	-0.	-5.0000E-01
4	3	-9.3184E-03	-1.0921E-01	0.	5.0000E-01	-0.	-5.0000E-01
4	4	0.	0.	-8.1243E-02	-6.8868E-01	8.1243E-02	6.8868E-01
5	1	-1.9702E-02	-1.7963E-01	0.	5.0000E-01	-0.	-5.0000E-01
5	2	-1.8309E-02	-1.6927E-01	0.	5.0000E-01	-0.	-5.0000E-01
5	3	-1.1136E-02	-1.1632E-01	0.	5.0000E-01	-0.	-5.0000E-01
5	4	-6.4766E-03	-7.1176E-02	0.	5.0000E-01	-0.	-5.0000E-01
5	5	0.	0.	-2.9390E-02	-1.0643E-01	2.9390E-02	1.0643E-01
6	1	-2.8031E-02	-2.1219E-01	0.	5.0000E-01	-0.	-5.0000E-01
6	2	-2.1745E-02	-1.8302E-01	0.	5.0000E-01	-0.	-5.0000E-01
6	3	-1.9912E-02	-1.6773E-01	0.	5.0000E-01	-0.	-5.0000E-01
6	4	-1.6520E-02	-1.3609E-01	0.	5.0000E-01	-0.	-5.0000E-01
6	5	-9.7556E-03	-8.4994E-02	0.	5.0000E-01	-0.	-5.0000E-01
6	6	0.	0.	-2.6957E-02	-8.9705E-02	2.6957E-02	8.9705E-02
7	1	-2.5458E-02	-2.0545E-01	0.	5.0000E-01	-0.	-5.0000E-01
7	2	-3.2195E-02	-2.2456E-01	0.	5.0000E-01	-0.	-5.0000E-01
7	3	-3.0388E-02	-2.1050E-01	0.	5.0000E-01	-0.	-5.0000E-01
7	4	-2.7375E-02	-1.9077E-01	0.	5.0000E-01	-0.	-5.0000E-01
7	5	-2.2021E-02	-1.5070E-01	0.	5.0000E-01	-0.	-5.0000E-01
7	6	0.	0.	-1.7538E-01	-9.7694E-01	-2.1059E-02	-3.0891E-01
7	7	0.	0.	-2.2307E-02	-5.6588E-02	2.2307E-02	5.6588E-02
7	8	0.	0.	0.	0.	0.	0.
7	9	-1.4103E-02	-9.2313E-02	2.0016E-01	1.4887E 00	-2.0016E-01	-1.4887E 00

VELOCITY POTENTIAL DIFFERENCES (PHI UPPER - PHI LOWER) AND  
SOURCE STRENGTHS WITH INTERFERENCE EFFECTS FOR MODE I

N	M	R DPHI(N,M)	I DPHI(N,M)	R SSU(N,M)	I SSU(N,M)	R SSL(N,M)	I SSL(N,M)
8	1	-3.8351E-02	-2.5017E-01	0.	5.0000E-01	-0.	-5.0000E-01
8	2	-3.6414E-02	-2.4249E-01	0.	5.0000E-01	-0.	-5.0000E-01
8	3	-4.2690E-02	-2.5353E-01	0.	5.0000E-01	-0.	-5.0000E-01
8	4	-3.9303E-02	-2.3258E-01	0.	5.0000E-01	-0.	-5.0000E-01
8	5	-3.6795E-02	-2.1119E-01	3.3287E-02	6.6974E-01	-0.	-5.0000E-01
8	6	0.	0.	-2.1562E-01	-1.0124E 00	8.4990E-02	1.8635E-01
8	7	0.	0.	-6.3473E-02	-1.8435E-01	3.2249E-02	1.3231E-02
8	8	0.	0.	-1.9658E-02	-4.1084E-02	1.9658E-02	4.1084E-02
8	9	-2.7680E-02	-1.5158E-01	1.9491E-01	1.3092E 00	-1.9491E-01	-1.3092E 00
8	10	-1.6987E-02	-9.8980E-02	6.8009E-02	6.9002E-01	-6.8009E-02	-6.9002E-01
9	1	-5.0137E-02	-2.8434E-01	0.	5.0000E-01	-0.	-5.0000E-01
9	2	-5.1254E-02	-2.8642E-01	0.	5.0000E-01	-0.	-5.0000E-01
9	3	-4.7701E-02	-2.7254E-01	0.	5.0000E-01	-0.	-5.0000E-01
9	4	-5.5190E-02	-2.8038E-01	1.4519E-02	5.5368E-01	-0.	-5.0000E-01
9	5	-4.7753E-02	-2.4151E-01	2.2102E-02	6.1819E-01	-0.	-5.0000E-01
9	6	0.	0.	-2.2390E-01	-1.0615E 00	4.8558E-02	-8.0207E-02
9	7	0.	0.	-9.7709E-02	-2.9466E-01	3.8760E-02	-1.1195E-01
9	8	0.	0.	0.	0.	0.	0.
9	9	-4.0363E-02	-2.0277E-01	2.0292E-01	1.3756E 00	-2.0292E-01	-1.3756E 00
9	10	-3.2927E-02	-1.6301E-01	8.4994E-02	7.1431E-01	-8.4994E-02	-7.1431E-01

VELOCITY POTENTIAL DIFFERENCES (PHI UPPER - PHI LOWER) AND  
SOURCE STRENGTHS WITH INTERFERENCE EFFECTS FOR MODE 2

N	M	R DPHI(N,M)	I DPHI(N,M)	R SSU(N,M)	I SSU(N,M)	R SSL(N,M)	I SSL(N,M)
1	1	-6.2191E-02	-5.5772E-04	1.0000E 00	2.6922E-02	-1.0000E 00	-2.6922E-02
2	1	-1.1005E-01	-2.1313E-03	1.0000E 00	8.0767E-02	-1.0000E 00	-8.0767E-02
2	2	-1.0045E-01	-2.0031E-03	1.0000E 00	8.0767E-02	-1.0000E 00	-8.0767E-02
3	1	-2.0727E-01	-6.0365E-03	1.0000E 00	1.3461E-01	-1.0000E 00	-1.3461E-01
3	2	-1.7382E-01	-5.3773E-03	1.0000E 00	1.3461E-01	-1.0000E 00	-1.3461E-01
3	3	-1.2702E-01	-4.1519E-03	1.0000E 00	1.3461E-01	-1.0000E 00	-1.3461E-01
4	1	-2.7179E-01	-1.1739E-02	1.0000E 00	1.8846E-01	-1.0000E 00	-1.8846E-01
4	2	-2.7388E-01	-1.1207E-02	1.0000E 00	1.8846E-01	-1.0000E 00	-1.8846E-01
4	3	-2.2015E-01	-9.2931E-03	1.0000E 00	1.8846E-01	-1.0000E 00	-1.8846E-01
4	4	0.	0.	-1.3886E 00	3.6504E-02	1.3886E 00	-3.6504E-02
5	1	-3.6439E-01	-1.9110E-02	1.0000E 00	2.4230E-01	-1.0000E 00	-2.4230E-01
5	2	-3.4323E-01	-1.8790E-02	1.0000E 00	2.4230E-01	-1.0000E 00	-2.4230E-01
5	3	-2.3567E-01	-1.8492E-02	1.0000E 00	2.4230E-01	-1.0000E 00	-2.4230E-01
5	4	-1.4409E-01	-1.3895E-02	1.0000E 00	2.4230E-01	-1.0000E 00	-2.4230E-01
5	5	0.	0.	-2.1681E-01	4.3481E-02	2.1681E-01	-4.3481E-02
6	1	-4.3326E-01	-2.8715E-02	1.0000E 00	2.9615E-01	-1.0000E 00	-2.9615E-01
6	2	-3.7351E-01	-3.3454E-02	1.0000E 00	2.9615E-01	-1.0000E 00	-2.9615E-01
6	3	-3.4209E-01	-3.1188E-02	1.0000E 00	2.9615E-01	-1.0000E 00	-2.9615E-01
6	4	-2.7734E-01	-2.3899E-02	1.0000E 00	2.9615E-01	-1.0000E 00	-2.9615E-01
6	5	-1.7295E-01	-1.7733E-02	1.0000E 00	2.9615E-01	-1.0000E 00	-2.9615E-01
6	6	0.	0.	-1.8313E-01	3.9905E-02	1.8313E-01	-3.9905E-02
7	1	-4.2254E-01	-5.5862E-02	1.0000E 00	3.4999E-01	-1.0000E 00	-3.4999E-01
7	2	-4.6203E-01	-4.6288E-02	1.0000E 00	3.4999E-01	-1.0000E 00	-3.4999E-01
7	3	-4.3280E-01	-4.2439E-02	1.0000E 00	3.4999E-01	-1.0000E 00	-3.4999E-01
7	4	-3.9182E-01	-3.9292E-02	1.0000E 00	3.4999E-01	-1.0000E 00	-3.4999E-01
7	5	-3.0925E-01	-2.9960E-02	1.0000E 00	3.4999E-01	-1.0000E 00	-3.4999E-01
7	6	0.	0.	-2.0108E 00	-5.9613E-02	-6.3261E-01	-2.0052E-01
7	7	0.	0.	-1.1625E-01	3.6067E-02	1.1625E-01	-3.6067E-02
7	8	0.	0.	0.	0.	0.	0.
7	9	-1.8978E-01	-1.8675E-02	3.0466E 00	3.5512E-01	-3.0466E 00	-3.5512E-01

VELOCITY POTENTIAL DIFFERENCES (PHI UPPER - PHI LOWER) AND  
SOURCE STRENGTHS WITH INTERFERENCE EFFECTS FOR MODE 2

N	N	R DPHI(N,M)	I DPHI(N,M)	R SSU(N,M)	I SSU(N,M)	R SSL(N,M)	I SSL(N,M)
8	1	-5.1958E-01	-7.0037E-02	1.0000E 00	4.0384E-01	-1.0000E 00	-4.0384E-01
8	2	-5.0334E-01	-6.9919E-02	1.0000E 00	4.0384E-01	-1.0000E 00	-4.0384E-01
8	3	-5.2620E-01	-5.7545E-02	1.0000E 00	4.0384E-01	-1.0000E 00	-4.0384E-01
8	4	-4.8228E-01	-5.2450E-02	1.0000E 00	4.0384E-01	-1.0000E 00	-4.0384E-01
8	5	-4.3801E-01	-4.5725E-02	1.3527E 00	4.2362E-01	-1.0000E 00	-4.0384E-01
8	6	0.	0.	-2.1076E 00	-7.1716E-02	3.9800E-01	-2.0703E-01
8	7	0.	0.	-3.8929E-01	6.0494E-02	3.4936E-02	-8.5057E-02
8	8	0.	0.	-8.4881E-02	3.3157E-02	8.4881E-02	-3.3157E-02
8	9	-3.1456E-01	-3.0362E-02	2.6968E 00	4.1413E-01	-2.6968E 00	-4.1413E-01
8	10	-2.0470E-01	-2.2627E-02	1.4023E 00	3.1490E-01	-1.4023E 00	-3.1490E-01
9	1	-5.9661E-01	-8.7007E-02	1.0000E 00	4.5768E-01	-1.0000E 00	-4.5768E-01
9	2	-6.0111E-01	-8.5535E-02	1.0000E 00	4.5768E-01	-1.0000E 00	-4.5768E-01
9	3	-5.7153E-01	-8.4815E-02	1.0000E 00	4.5768E-01	-1.0000E 00	-4.5768E-01
9	4	-5.8817E-01	-6.7913E-02	1.1135E 00	4.5604E-01	-1.0000E 00	-4.5768E-01
9	5	-5.0639E-01	-5.9909E-02	1.2473E 00	4.9233E-01	-1.0000E 00	-4.5768E-01
9	6	0.	0.	-2.2299E 00	-1.8398E-01	-1.4943E-01	-2.9759E-01
9	7	0.	0.	-6.2593E-01	6.8099E-02	-2.1256E-01	-2.1521E-01
9	8	0.	0.	0.	0.	0.	0.
9	9	-4.2509E-01	-5.0567E-02	2.8561E 00	5.7705E-01	-2.8561E 00	-5.7705E-01
9	10	-3.4082E-01	-3.7447E-02	1.4614E 00	3.5701E-01	-1.4614E 00	-3.5701E-01

Contrails

MACH BOX INTERFERENCE PROGRAM OUTPUT

OSCILLATORY FREQUENCY (CPS) 100.00000 9 BOXES IN CHORD DIRECTION  
 REDUCED FREQUENCY (BOX LENGTH) 0.05384 FREE STREAM MACH NUMBER 2.000  
 REDUCED FREQUENCY (ROOT CHORD) 0.50000 TIP FOLD ANGLE (DEGREES) 30.000

GENERALIZED FORCES

DPHI	DEFL	REAL PART	IMAG PART	ABS VALUE	PHASE ANGLE
1	1	-1.68557E-02	-9.30777E-01	9.30929E-01	-91.037 DEG
1	2	-1.23360E-02	-5.83169E-01	5.83299E-01	-91.212 DEG
2	1	-1.88000E 00	-6.25115E-01	1.98121E 00	-161.608 DEG
2	2	-1.18191E 00	-4.56494E-01	1.26700E 00	-158.882 DEG



*Contrails*

APPENDIX IV. PROGRAM LISTINGS

MACH BOX INTERFERENCE PROGRAM

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MBX
      MAIN PROGRAM
COMMON/A/DPHI(20,60,2),SSU(20,60,2),SSL(20,60,2),THW(10),THT(10)
COMMON/B/MIHW(20),MUBW(20),MIBT(20),MORT(20),MORTD(20),MORWD(20)
COMMON/C/BARNU,BARNU,EL,CFR,CFI,CNR,CNI,EKBAR,EKM,P(10),W(10)
COMMON/D/XTE(4),YTE(4),QR(21,2),QI(21,2),COW(10,21),COT(10,21)
COMMON/E/BOXL,BOXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBOX,CAPPHI(20,20,2)
COMMON/F/MODE,KODE,NP,MP,MFL,NMAX,NBW,NBT,NPINT,NPFLET,NPOLW,NPOLTMBOX0070
COMMON/G/SLEW,STEW,YFL,SLET,STET,YTIP,CROOT,CFL,CTIP,MDET,MTIP
      DIMENSION
1XINW(100),YINW(100),ZINW(10,100),XINT(100),YINT(100),ZINT(10,100),
ZOMEGA(20),ANGS(20),Q(10,10,2)
      EQUIVALENCE
1(DPHI,XINW,XINT),(DPHI(101),YINW,YINT),(DPHI(201),ZINW,ZINT)
2345 READ (5,31) EMACH,AS,CROOT
      IF(EMACH.GT.1.C.AND.EMACH.LE.3.5) GO TO 2348
      WRITE (6,2347)
2347 FORMAT(1H1,20X,41HMACH NUMBER OUT OF LIMITS,PROGRAM STOPPED)
      CALL EXIT
2348 READ (5,31) SLEW,STEW,YFL,SLET,STET,YTIP
      31 FORMAT(6E12.5)
      READ (5,10) NBOX,NFREQ,MODES,LVPIC,LSSJL,LDPHI
      READ (5,10) MDEW,NPOLW,NPOW,LCOW
      READ (5,10) MDET,NPOLT,NPOT,LCOT
      10 FORMAT(6I12)
      CALL HOUNDS(ERROR)
      IF(ERROR)63,98,63
      63 WRITE (6,74) EMACH
      74 FORMAT(51H1 YOUR CONFIGURATION HAD A SUBSONIC TRAILING EDGE,/ 2X,MBX00280
161H10 MAKE THAT EDGE SONIC THE MACH NUMBER HAS BEEN INCREASED TO//MBX00290
22CX, 8HEMACH = , F8.5)
      98 CONTINUE
      AREA = YFL*(CROOT - CTIP) + YTIP*(CFL + CTIP)
      QF = -4.0 * BOXL * BOXW / AREA
      WRITE (6,309)
      309 FORMAT(1H1//////// 21X,63HFORTRAN IV PROGRAM FOR CALCULATING
1AERODYNAMICS OF INTERSECTING /1HC,20X,63HSUPERSONIC LIFTING SURFACMBX00360
2ES BY THE SOURCE SUPERPOSITION METHOD. /1HC,20X,61HTHE MACH BOX TEMMBX00370
3CHNIQUE IS APPLIED TO A WING WITH FOLDED TIPS. )
      WRITE (6,310) EMACH,AS,CROOT,AREA,SLEW,STEW,YFL,CFL,
MBX000020
MBX000030
MBX000040
MBX000050
MBX000060
MBX000070
MBX000080
MBX000090
MBX000100
MBX000110
MBX000120
MBX000130
MBX000140
MBX000150
MBX000160
MBX000170
MBX000180
MBX000190
MBX000200
MBX000210
MBX000220
MBX000230
MBX000240
MBX000250
MBX000260
MBX000270
MBX000280
MBX000290
MBX000300
MBX000310
MBX000320
MBX000330
MBX000340
MBX000350
MBX000360
MBX000370
MBX000380
MBX000390

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MACH BOX INTERFERENCE PROGRAM

MBOX MAIN PROGRAM

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15 LET,SIET,YTIP,CTIP,NBOX,BOXL,MTIP,BOXW
310 FORMAT(1H-//,37X,30HFLIGHT CONDITIONS AND GEOMETRY,/1H-/
1 11X,11HMACH NUMBER,11X,14HSPEED OF SOUND,13X,10HROOT CHORD,12X, MBX00400
2 14HREFERENCE AREA//F20.5,F22.3,4H L/T,F21.3,2H L,F21.3,5H L*2/ MBX00420
3 1H-,8X,15HWING L.E. SWEEP,9X,15HWING T.E. SWEEP,9X,14HFOLD LINE MBX00430
4 SPAN,10X,15HFOLD LINE CHORD//F18.3,4H DEG,F20.3,4H DEG,F21.3, MBX00440
5 2H L,F21.3,2H L/1H-,9X,14HTIP L.E. SWEEP,10X,14HTIP T.E. SWEEP, MBX00450
6 10X,13HTIP LINE SPAN,11X,14HTIP LINE CHORD//F18.3,4H DEG,F20.3, MBX00460
7 4H DEG,F21.3,2H L,F21.3,2H L/1H-,8X,15HCHORDWISE BOXES,12X, 9HB0X MBX00470
8 CHORD,12X,14HSPANWISE BOXES,14X,8HB0X SPAN//I17,F26.4,2H L,I19, MBX00480
9 F27.4,2H L) MBX00490
WRITE (6,311) (MIBW(I),MOBW(I),MIBT(I),MOBT(I),MOBTD(I),
1MOBWD(I),I = 1,NBOX) MBX00500
311 FORMAT(1H1,31X,43HINNER AND OUTER WING AND TIP BOUNDARY BOXES/1H-/MBX00510
112X,4HMIBW,11X,4HMOBW,11X,4HMIBT,11X,4HMOBT,11X,5HMOBTD,10X, MBX00520
25HMOBWD /1H0/(6I15)) MBX00530
WRITE (6,303) (XTE(I),YTE(I),I = 1,MTIP) MBX00540
303 FORMAT(1H1,39X,25HTRAILING EDGE COORDINATES/1H-/43X, 6HX/BOXL, 8X, MBX00550
1 6HY/BOXW /1H0/(35X,2F14.3)) MBX00560
NMAX = 1 MBX00570
DO 123 I = 1,NBOX MBX00580
123 NMAX = MAX0(MOBWD(I),NMAX) MBX00590
READ (5,31) (OMEGA(I),ANGS(I),I=1,NFREQ) MBX00600
DO 820 L=1,NFREQ MBX00610
EK = OMEGA(L) * BOXL / (EMACH * AS) * 6.2831853 MBX00620
EKRR = EK * (EMACH**2) / ((EMACH**2) - 1.0) MBX00630
EKKEF = EK * CR00T/BOXL MBX00640
EKM = EKBAR/EMACH MBX00650
FAN = ANGS(L) MBX00660
DO 140 I1 = 1,20 MBX00670
DO 140 I2 = 1,20 MBX00680
CAPPHI(I1,I2,1) = 0.0 MBX00690
CAPPHI(I1,I2,2) = 0.0 MBX00700
CALL CAPHI MBX00710
IF(LVPIC)198,202,198 MBX00720
198 WRITE (6,111) MBX00730
111 FORMAT(1H1,27X,48HPLANAR VELOCITY POTENTIAL INFLUENCE COEFFICIENTS MBX00740
MBX00750

```

MACH BOX INTERFERENCE PROGRAM

```

MBØX      MAIN PROGRAM
1/1H-,9X,1H1,5X,1HJ,10X,5HBARNU,5X,5HBARMU,10X,16HREAL CAPPHI(I,J) MBXØØ0770
2,6X,16HIMAG CAPPHI(I,J),/1HØ) MBXØØ0780
ØØ 2ØØ I = 1,NBØX MBXØØ0790
ØØ 2ØØ J = 1,I MBXØØ0800
BARNU = I - 1 MBXØØ0810
BARMU = J - 1 MBXØØ0820
2ØØ WRITE (6,5) I,J,BARNU,BARMU,CAPPHI(I,J,1),CAPPHI(I,J,2) MBXØØ0830
5 FORMAT( 5X,2I6,4X,2F10.1,4X,1P1E20.4,1P1E22.4) MBXØØ0840
2Ø2 IF(L.GT.1 .AND. ANGSL(L).LT.1.Ø) GØ TØ 95 MBXØØ0850
2Ø READ (5,31) FANG MBXØØ0860
CFA = CØSD(FANG) MBXØØ0870
SFA = SIND(ABS(FANG)) MBXØØ0880
IF(MDEW)3Ø,6Ø,6Ø MBXØØ0890
3Ø ØØ 36 IX = 1,1ØØ MBXØØ0900
XINW(IX) = Ø.Ø MBXØØ0910
YINW(IX) = Ø.Ø MBXØØ0920
ØØ 36 IZ = 1,1Ø MBXØØ0930
36 ZINW(IZ,IX) = Ø.Ø MBXØØ0940
READ (5,31) (XINW(J),YINW(J),J=1,NPØW) MBXØØ0950
ØØ 5Ø I = 1,MØDES MBXØØ0960
5Ø READ (5,31) THW(I),(ZINW(I,J),J=1,NPØW) MBXØØ0970
CALL LSSUR(XINW,YINW,ZINW,NPØW,MØDES,NPØLW+1,NCW,CØW) MBXØØ0980
GØ TØ 81 MBXØØ0990
6Ø NCW = ((NPØLW + 3)*NPØLW + 2)/2 MBXØ1000
7Ø ØØ 8Ø I = 1,MØDES MBXØ1010
8Ø READ (5,31) THW(I),(CØW(I,J),J=1,NCW) MBXØ1020
81 IF(LCØW)84,85,84 MBXØ1030
84 WRITE (6,77Ø) MBXØ1040
77Ø FORMAT(1H1,32X,39HWING DEFLECTION POLYNØMIAL CØEFFICIENTS/1HØ,11X, MBXØ1050
138HZ(MØDE) = SUM(A(MØDE,T)*X**(R-S)*Y**S) /16X, MBXØ1060
232HWHERE T = (R**2 + R + 2*S + 2)/2 /1H-, 4X, 4HMØDE, 3X,25HA(MØDE MBXØ1070
3,T),T = 1,2,--,NCW ) MBXØ1080
ØØ 771 I = 1,MØDES MBXØ1090
771 WRITE (6,772) I,(CØW(I,J),J=1,NCW) MBXØ1100
772 FORMAT(1HØ, 17,1P6E15.4 /( 8X,1P6E15.4)) MBXØ1110
85 IF(MØET)9Ø,1Ø8,11Ø MBXØ1120
1Ø8 NPØLT = -1 MBXØ1130

```

MACH BOX INTERFERENCE PROGRAM

MBX MAIN PROGRAM

```

90 DO 96 IX = 1,100
  XINT(IX) = 0.0
  YINT(IX) = 0.0
  DO 96 IZ = 1,10
    ZINT(IZ,IX) = 0.0
  READ (5,31) (XINT(J),YINT(J),J=1,NPOT)
  DO 105 I = 1,MODES
    READ (5,31) TH(I),(ZINT(I,J),J=1,NPOT)
    CALL LSSUR(XINT,YINT,ZINT,NPOT,MODES,NPOT+1,NCT,COT)
  GO TO 800
110 NCT = ((NPOT + 3)*NPOT + 2)/2
120 DO 130 I = 1,MODES
130 READ (5,31) TH(I),(COT(I,J),J=1,NCT)
800 IF(LCOT)94,95,94
  94 WRITE (6,780)
780 FORMAT(1H1,33X,38HTIP DEFLECTION POLYNOMIAL COEFFICIENTS /1H0,11X,
138H2(MODE) = SUM(A(MODE,T)*X**(R-S)*Y**S) /16X,
232HWHERE T = (R**2 + R + 2*S + 2)/2 /1H-, 4X, 4HMODE, 3X,25HA(MODEM
3,T),T = 1,2,--,NCT )
  DO 781 I = 1,MODES
781 WRITE (6,772) I,(COT(I,J),J=1,NCT)
95 DO 801 I1 = 1,10
  DO 801 I2 = 1,10
    Q(I1,I2,1)=0.0
    Q(I1,I2,2)=0.0
  DO 810 MODE=1,MODES
    IF(THW(MODE))807,810,807
807 MODE = MODE
    DO 2 L1=1,20
    DO 2 L2=1,60
    DO 2 L3=1,2
      DPHI(L1,L2,L3) = 0.0
      SSU(L1,L2,L3)=0.0
      SSL(L1,L2,L3)=0.0
    CALL SOURCE
    IF(LSSUL)150,160,150

```

MBX01140  
 MBX01150  
 MBX01160  
 MBX01170  
 MBX01180  
 MBX01190  
 MBX01200  
 MBX01210  
 MBX01220  
 MBX01230  
 MBX01240  
 MBX01250  
 MBX01260  
 MBX01270  
 MBX01280  
 MBX01290  
 MBX01300  
 MBX01310  
 MBX01320  
 MBX01330  
 MBX01340  
 MBX01350  
 MBX01360  
 MBX01370  
 MBX01380  
 MBX01390  
 MBX01400  
 MBX01410  
 MBX01420  
 MBX01430  
 MBX01440  
 MBX01450  
 MBX01460  
 MBX01470  
 MBX01480  
 MBX01490  
 MBX01500

MACH BOX INTERFERENCE PROGRAM

MBOX MAIN PROGRAM

```

150 WRITE (6,4) MODE
4  FORMAT(1H1,32X,34HINITIAL SOURCE STRENGTHS FOR MODE ,I2 / 1H-,9X,
1 1HN, 5X,1HM, 6X,13HREAL SSU(N,M) 5X,13HIMAG SSU(N,M), 5X,
2 13HREAL SSL(N,M), 5X,13HIMAG SSL(N,M) /1H0)
DO 100 N = 1,NBOX
MPEND = MGBWD(N)
IF(N.GE.NPFLET) MPEND = NMAX+MGBTD(N)-MFL
DO 100 M = 1,MPEND
100 WRITE (6,19) N,M,SSU(N,M,1),SSU(N,M,2),SSL(N,M,1),SSL(N,M,2)
19  FORMAT( 5X,21b, 1P4E18.4 )
160 DO 1605 I2 = 1,2
DO 1605 I1 = 1,21
QR(I1,I2) = 0.0
1605 QI(I1,I2) = 0.0
CALL CODE
IF(LDPhi)73,83,73
73 WRITE (6,76) MODE
76  FORMAT(1H1,23X,58HVELOCITY POTENTIAL DIFFERENCES (PHI UPPER - PHI
1  LOWER) AND /26X,51HSOURCE STRENGTHS WITH INTERFERENCE EFFECTS FOR
2  MODE,13/1H-,6X,1HN,3X,1HM,5X,11HR DPHI(N,M),3X,11HI DPHI(N,M),4X,
3  10HR SSU(N,M),4X,10HI SSU(N,M),4X,10HR SSL(N,M),4X,10HI SSL(N,M)
4  /1H0)
DO 75 N = 1,NBOX
MPEND = MGBWD(N)
IF(N.GE.NPFLET) MPEND = NMAX+MGBTD(N)-MFL
DO 75 M = 1,MPEND
75 WRITE (6,72) N,M,DPHI(N,M,1),DPHI(N,M,2),SSU(N,M,1),
1  SSU(N,M,2),SSL(N,M,1),SSL(N,M,2)
72  FORMAT( 4X,214,2X,1P6E14.4)
83  CONTINUE
NC = MAXO(NCW,NCI)
DO 750 J = 1,MODES
DO 750 N = 1,NC
Q(MODE,J,1) = Q(MODE,J,1) + QR(N,1)*COW(J,N) + QR(N,2)*COT(J,N)
750 Q(MODE,J,2) = Q(MODE,J,2) + QI(N,1)*COW(J,N) + QI(N,2)*COT(J,N)
810 CONTINUE
WRITE (6,2001) OMEGA(L),NBOX,EK,EMACH,EKREF,FANG

```



MACH BOX INTERFERENCE PROGRAM

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MBUX          MAIN PROGRAM

2001 FORMAT(1H1,33X,37HMACH BOX INTERFERENCE PROGRAM OUTPUT ,//
1 1H0,8X,27HOSCILLATORY FREQUENCY (CPS),F12.5,14X,I2, 25H BOXES INMRX01890
2 CHORD DIRECTION /1H0,8X,30HREDUCED FREQUENCY (BOX LENGTH),F9.5, MBX01900
3 14X,23HFREE STREAM MACH NUMBER, F9.3,1H0,8X,30HREDUCED FREQUENCYMBX01910
4 (ROOT CHORD),F9.5,14X,24HTIP FOLD ANGLE (DEGREES), F8.3 /1H-/42X,MRX01920
5 18HGENERALIZED FORCES /1H-,6X,4HDPHI,3X,4HDEFL,10X,9HREAL PART, MBX01930
6 10X,9HIMAG PART,10X,9HABS VALUE,10X,11HPHASE ANGLE //)
      DO 2010 I = 1,M0DES
      DO 2010 J = 1,M0DES
      QREAL = QF * Q(I,J,1)
      QIMAG = QF * Q(I,J,2)
      QABS = SQRT(QREAL**2 + QIMAG**2)
      QANGLE = 0.0
      IF(QABS.NE.0.0) QANGLE = ATAND(QIMAG,QREAL)
2010 WRITE (6,2007) I,J,QREAL,QIMAG,QABS,QANGLE
2007 FORMAT(1H0,2X,2I7,2X,1P3E19.5,5P1F16.3,4H DEG)
      IF(FAN - 1.0)820,820,1211
1211 FAN = FAN - 1.0
      GO TO 20
820 CONTINUE
      GO TO 2345
      END
MBX01880
MBX01890
MBX01900
MBX01910
MBX01920
MBX01930
MBX01940
MBX01950
MBX01960
MBX01970
MBX01980
MBX01990
MBX02000
MBX02010
MBX02020
MBX02030
MBX02040
MBX02050
MBX02060
MBX02070
MBX02080
MBX02090

```

MACH BOX INTERFERENCE PROGRAM

```

BOUND      WING AND TIP BOUNDARYS

SUBROUTINE BOUNDS(ERROR)
COMMON/B/MIBW(20),MOBW(20),MIBT(20),MGBT(20),MORTD(20),MOBWD(20)
COMMON/D/XTE(40),YTE(40),QR(21,2),QI(21,2),COW(10,21),CGT(10,21)
COMMON/E/BOXL,BOXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBOX,CAPPHI(20,20,2),MBX02140
COMMON/F/MODE,KODE,NP,MP,MFL,NMAX,NBW,NBT,NPINT,NPFLET,NPQLW,NPQLT,MBXC2150
COMMON/G/SLEW,STEW,YFL,SLET,STET,YTIP,CROOT,CFL,CTIP,MDET,MTIP
TANDF(A) = SIND(A)/COSD(A)
ERROR = 0.0
DO 1070 I = 1,40
  XTE(I) = 0.0
  YTE(I) = 0.0
1070 DO 1080 I=1,20
  MIBW(I)=0
  MOBW(I)=0
  MOBWD(I)=0
  MIBT(I)=0
  MGBT(I)=0
  MORTD(I)=0
  5 BETA = SQRT(EMACH**2 - 1.0)
  PGWL = TANDF(SLEW) / BETA
  PGWT = TANDF(STEW) / BETA
  IF(1.0 - ABS(PGWT))10,20,20
10 EMACH = SQRT(1.0 + (PGWT * BETA)**2)
  ERROR = 1.0
  GO TO 5
20 PGTL = TANDF(SLET) / BETA
  PGTT = TANDF(STET) / BETA
  IF(1.0 - ABS(PGTT))30,40,40
30 EMACH = SQRT(1.0 + (PGTT * BETA**2))
  ERROR = 1.0
  GO TO 5
40 IF(MDET.EQ. 0 ) YTIP = YFL
  DY = YTIP - YFL
  XFL = YFL * PGWL * BETA
  XFT = YFL * PGWT * BETA + CROOT
  CFL = XFT - XFL
  XTL = XFL + DY * PGTL * BETA

```

MACH BOX INTERFERENCE PROGRAM

BOUND WING AND TIP BOUNDARIES

```

XIT = XFT + DY * PGIT * BETA
CTIP = XIT - XTL
CREF = AMAX1(CROOT,XFT,XIT)
70 BOXL = CREF / FLOAT(NBOX)
BOXW = BOXL / BETA
MFL = 1.0 + YFL / BOXW
YNFL = FLOAT(MFL) - 0.5
BOXW = YFL / YNFL
BOXL = BOXW * BETA
IF(20.4*BOXL - CREF)80,90,90
80 NBOX = NBOX - 1
GO TO 70
90 MTIP = 1.0 + YTIP/BOXW
IF(60-2*MTIP+MFL-2-IFIX (CTIP/BOXL))80,100,100
100 XNRT = CROOT / BOXL
XNFL = XFL / BOXL
XNFT = XFT / BOXL
NBOX = CREF/BOXL + 0.5001
NPROOT = XNRT + 0.5001
NPFLEW = XNFL - PGWL/2.0 + 1.5001
NPFTEW = XNFT - PGWT/2.0 + 0.5001
IF(NPFTEW - NPFLEW)110,120,120
110 NPFTEW = NPFLEW
120 NPFLTE = XNFL + PGIL/2.0 + 1.5001
NPFTET = XNFT + PGIT/2.0 + 0.5001
130 NPIPLE = XNFL + PGIL*(FLOAT(MTIP-1) - YNFL) + 1.5001
NPITPE = XNFT + PGIT*(FLOAT(MTIP-1) - YNFL) + 0.5001
IF(NPITPE - NPIPLE)140,150,150
140 MTIP = MTIP - 1
GO TO 130
150 NBW = MAX0(NPROOT,NPFTEW)
NBT = MAX0(NPFTET,NPITPE)
MFLT = MFL + 1
DO 220 NP = 1,NBW
MIBW(NP) = 1
IF(NP - NPROOT)170,170,160
160 MIBW(NP) = (FLOAT(NP) - 0.5 - XNRT)/PGWT + 2.0001

```

MBX02480  
 MBX02490  
 MBX02500  
 MBX02510  
 MBX02520  
 MBX02530  
 MBX02540  
 MBX02550  
 MBX02560  
 MBX02570  
 MBX02580  
 MBX02590  
 MBX02600  
 MBX02610  
 MBX02620  
 MBX02630  
 MBX02640  
 MBX02650  
 MBX02660  
 MBX02670  
 MBX02680  
 MBX02690  
 MBX02700  
 MBX02710  
 MBX02720  
 MBX02730  
 MBX02740  
 MBX02750  
 MBX02760  
 MBX02770  
 MBX02780  
 MBX02790  
 MBX02800  
 MBX02810  
 MBX02820  
 MBX02830  
 MBX02840

MACH BOX INTERFERENCE PROGRAM

```

BOUND      WING AND TIP BOUNDARYS

170 IF(NP - NPFTW)190,190,180
180 M8BW(NP) = (FLGAT(NP) - 0.5 - XNRT)/PGWT + 1.0001
GO TO 220

190 M8BW(NP) = MFL
M8BWD(NP) = MAXO(NP,NP+MFL-NPFLEW)
IF(NP - NPFLEW)200,210,210
200 M8BW(NP) = (FLGAT(NP) - 0.5)/PGWL + 1.0001
M8BWD(NP) = MAXO(NP,M8BW(NP))
210 MP = MFL + NPFTW - NP
IF(MDET)215,220,215
215 MP = MTIP + NPTPE - NP
220 M8BWD(NP) = MINO(M8BWD(NP),MP)
DO 225 MP = 1,MFL
YTE(MP) = MP - 1
225 XTE(MP) = XNRT + FLGAT(MP-1)*PGWT
IF(MDET)230,230,230
230 DO 300 NP = NPFLET,NBT
MIRT(NP) = MFLT
IF(NP - NPFTET)250,250,240
240 MIRT(NP) = YNFL + (FLGAT(NP) - XNFT - 0.5)/PGTT + 2.0001
250 IF(NP - NPTPTE)270,270,260
260 M8BT(NP) = YNFL + (FLGAT(NP) - XNFT - 0.5)/PGTT + 1.0001
GO TO 300
270 M8BT(NP) = MTIP
MP = NP - NPFLET + MFLT
M8BTD(NP) = MAXO(MP,NP+MTIP-NPTPLE)
IF(NP - NPTPLE)280,290,290
280 M8BT(NP) = YNFL + (FLGAT(NP) - XNFL - 0.5)/PGTL + 1.0001
M8BTD(NP) = MAXO(MP,M8BT(NP))
290 M8BTD(NP) = MINO(M8BTD(NP),MTIP+NPTPTE - NP)
M8BWD(NP) = MAXO(M8BTD(NP),M8BWD(NP))
300 CONTINUE
DO 310 MP = MFLT,MTIP
YTE(MP) = MP - 1
310 XTE(MP) = XNFT + (FLGAT(MP-MFL) - 0.5)*PGTT
320 RETURN
END

```

MBX02850  
 MBX02860  
 MBX02870  
 MBX02880  
 MBX02890  
 MBX02900  
 MBX02910  
 MBX02920  
 MBX02930  
 MBX02940  
 MBX02950  
 MBX02960  
 MBX02970  
 MBX02980  
 MBX02990  
 MBX03000  
 MBX03010  
 MBX03020  
 MBX03030  
 MBX03040  
 MBX03050  
 MBX03060  
 MBX03070  
 MBX03080  
 MBX03090  
 MBX03100  
 MBX03110  
 MBX03120  
 MBX03130  
 MBX03140  
 MBX03150  
 MBX03160  
 MBX03170  
 MBX03180  
 MBX03190  
 MBX03200  
 MBX03210

MACH BOX INTERFERENCE PROGRAM

```

CAPHI      IN-PLANE VPICS

C          INPLANE VELOCITY POTENTIAL INFLUENCE COEFFICIENTS.
SUBROUTINE CAPHI
COMMON/C/BARNU,BARMU,EL,CFR,CFI,CNR,CNI,EKBAR,EKM,P(10),W(10)
COMMON/E/BOXL,BOXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBOX,CAPPHI(20,20,2)
IF(EKBAR)99,20,30
20 CAPPHI(1,1,1) = -0.5
GO TO 50
30 DO 40 I = 1,5
J = 6 - I
ARG = EKBAR * P(J) / 2.0
ARGM = ARG / EMACH
BSL = ZJ(ARGM)
CAPPHI(1,1,1) = CAPPHI(1,1,1) - COS(ARG)*BSL*W(I)/2.0
CAPPHI(1,1,2) = CAPPHI(1,1,2) + SIN(ARG)*BSL*W(I)/2.0
40 CONTINUE
CAPPHI(1,1,1) IS NOW COMPLETE
50 EL = 0.0
DO 60 NP = 2,NBOX
BARNU=NP-1
DO 60 MP = 1,NP
BARMU=MP-1
CALL VPIC
CAPPHI(NP,MP,1) = CFR
CAPPHI(NP,MP,2) = CFI
60 CONTINUE
99 RETURN
END

```

MBX03230  
 MBX03240  
 MBX03250  
 MBX03260  
 MBX03270  
 MBX03280  
 MBX03290  
 MBX03300  
 MBX03310  
 MBX03320  
 MBX03330  
 MBX03340  
 MBX03350  
 MBX03360  
 MBX03370  
 MBX03380  
 MBX03390  
 MBX03400  
 MBX03410  
 MBX03420  
 MBX03430  
 MBX03440  
 MBX03450  
 MBX03460  
 MBX03470  
 MBX03480  
 MBX03490

MACH BOX INTERFERENCE PROGRAM

```

SOURCE      INITIAL SOURCE STRENGTHS

C          SETS UP ORIGINAL SOURCE STRENGTH ARRAY FOR ONE MODE
C          WITH SSU(REAL)= DZDX = - SSL(REAL) ON SURFACE AND ZERO ON DIAFRAM
C          SSU(IMAG)= EK*Z = - SSL(IMAG) ON SURFACE AND ZERO ON DIAFRAM
C          SUBROUTINE SOURCE
C          COMMON/A/DPHI(20,60,2),SSU(20,60,2),SSL(20,60,2),THW(10),THT(10)
C          COMMON/B/MIBW(20),MOBW(20),MIBT(20),MIBD(20),MOBWD(20)
C          COMMON/D/XIE(40),YTE(40),QR(21,2),QI(21,2),COW(10,21),COT(10,21)
C          COMMON/E/BOXL,BOXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBOX,CAPPHI(20,20,2)
C          COMMON/F/MODE,KODE,NP,MP,MFL,NMAX,NBW,NBT,NPINT,NPFLET,NPOLW,NPOLT
C          DO 90 NP = 1,NBW
C          MPL = MIBW(NP)
C          MPR = MOBW(NP)
C          DO 90 MP = MPL,MPR
C          X =(FLOAT(NP) - 0.5)* BOXL
C          Y =(FLOAT(MP-1))*BOXW
C          Z = COW(MODE,1)
C          DZDX = 0.0
C          I = 1
C          J = 0
C          DO 70 N = 1,NPOLW
C          DO 70 M = J,N
C          NM = N-M
C          FN = NM
C          I = I + 1
C          YM = COW(MODE,I)
C          IF(M)30,40,30
C          30 YM = Y**M * YM
C          40 XM = 1.0
C          IF(NM)50,60,50
C          50 XM = X**NM
C          DZDX = DZDX + FN * XM/X * YM
C          Z = Z + XM * YM
C          60 Z = Z + XM * YM
C          70 CONTINUE
C          SSU(NP,MP,1) = DZDX
C          SSL(NP,MP,1) = -DZDX * SIGN(1.0,THW(MODE))
C          SSU(NP,MP,2) = EK * Z / BOXL
C          SSL(NP,MP,2) = -EK * Z * SIGN(1.0,THW(MODE))/BOXL

```

MBX03510  
 MBX03520  
 MBX03530  
 MBX03540  
 MBX03550  
 MBX03560  
 MBX03570  
 MBX03580  
 MBX03590  
 MBX03600  
 MBX03610  
 MBX03620  
 MBX03630  
 MBX03640  
 MBX03650  
 MBX03660  
 MBX03670  
 MBX03680  
 MBX03690  
 MBX03700  
 MBX03710  
 MBX03720  
 MBX03730  
 MBX03740  
 MBX03750  
 MBX03760  
 MBX03770  
 MBX03780  
 MBX03790  
 MBX03800  
 MBX03810  
 MBX03820  
 MBX03830  
 MBX03840  
 MBX03850  
 MBX03860  
 MBX03870



MACH BOX INTERFERENCE PROGRAM

SOURCE INITIAL SOURCE STRENGTHS

```

90 CONTINUE
  IF(NPOLT + 1)1000,1000,110
110 DO 200 NP = NPFLT, NBT
    MPL = MIBT(NP)
    MPR = MIBT(NP)
    DO 200 MP = MPL,MPR
      X = (FLGAT(NP) - 0.5)*BOXL
      Y = (FLGAT(MP-1))*BOXW
      MPN = MP + NMAX - MFL
      IF(ABS(FANG).LE. 5.0) MPN = MP
      Z = COT(MODE,1)
      DZDX = 0.0
      I = 1
      DO 170 N = 1,NPOLT
        DO 170 M = J,N
          NM = N - M
          FN = NM
          I = I + 1
          YM = COT(MODE,I)
          IF(M)130,140,130
130 YM = Y**M * YM
140 XM = 1.0
          IF(NM)150,160,150
150 XM = X**NM
          DZDX = DZDX + FN * XM/X * YM
160 Z = Z + XM * YM
170 CONTINUE
          SSU(NP,MPN,1) = DZDX
          SSL(NP,MPN,1) = -DZDX * SIGN(1.0,THT(MODE))
          SSU(NP,MPN,2) = EK * Z / BOXL
          SSL(NP,MPN,2) = -EK * Z * SIGN(1.0,THT(MODE)) / BOXL
200 CONTINUE
1000 RETURN
      END

```

MBX03880  
 MBX03890  
 MBX03900  
 MBX03910  
 MBX03920  
 MBX03930  
 MBX03940  
 MBX03950  
 MBX03960  
 MBX03970  
 MBX03980  
 MBX03990  
 MBX04000  
 MBX04010  
 MBX04020  
 MBX04030  
 MBX04040  
 MBX04050  
 MBX04060  
 MBX04070  
 MBX04080  
 MBX04090  
 MBX04100  
 MBX04110  
 MBX04120  
 MBX04130  
 MBX04140  
 MBX04150  
 MBX04160  
 MBX04170  
 MBX04180  
 MBX04190  
 MBX04200  
 MBX04210



MACH BOX INTERFERENCE PROGRAM

```

CJDE          BOX CODE
C IDENTIFIES BOX TYPE FOR LOGICAL PURPOSES WITH VARIOUS KODE VALUES
C KODE = 1 , WING IN NON-INTERFERENCE REGION
C KODE = 2 , WING DIAFRAM IN NON-INTERFERENCE REGION
C KODE = 3 , WING IN INTERFERENCE REGION
C KODE = 4 , TIP IN INTERFERENCE REGION
C KODE = 5 , TIP DIAFRAM IN INTERFERENCE REGION
C KODE = 6 , WING DIAFRAM IN INTERFERENCE REGION
C SUBROUTINE CODE
C COMMON/B/MIBW(20),MGBW(20),MIBT(20),MGBT(20),MGBTD(20),MGBWD(20)
C COMMON/E/BOXL,BOXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBOX,CAPPHI(20,20,2)
C COMMON/F/MODE,KODE,NP,MP,MFL,NMAX,NBW,NBT,NPINT,NPFLET,NPQLT,NPQLW,NPQLT
D3 1000 NP = 1,NBOX
NP = NP
MIW = MIBW(NP)
MOW = MGBW(NP)
MIT = MIBT(NP)
MOT = MGBT(NP)
MTO = MGBTD(NP)
MWD = MGBWD(NP)
NPINT = NP - NPFLET + 1
IF(NPINT)10,10,30
10 MPL = MIW
MPR = MOW
KODE = 1
CALL SSPHI(MPL,MPR)
MPL = MPR + 1
IF(MPL - MWD)20,20,1000
20 MPR = MWD
KODE = 2
CALL SSPHI(MPL,MPR)
GO TO 1000
30 IF(NPINT - MFL)40,40,50
40 MPL = MIW
MPR = MFL - NPINT + 1
KODE = 1
CALL SSPHI(MPL,MPR)
IF(NPINT - 1)1000,60,50

```

MBX04230  
 MBX04240  
 MBX04250  
 MBX04260  
 MBX04270  
 MBX04280  
 MBX04290  
 MBX04300  
 MBX04310  
 MBX04320  
 MBX04330  
 MBX04340  
 MBX04350  
 MBX04360  
 MBX04370  
 MBX04380  
 MBX04390  
 MBX04400  
 MBX04410  
 MBX04420  
 MBX04430  
 MBX04440  
 MBX04450  
 MBX04460  
 MBX04470  
 MBX04480  
 MBX04490  
 MBX04500  
 MBX04510  
 MBX04520  
 MBX04530  
 MBX04540  
 MBX04550  
 MBX04560  
 MBX04570  
 MBX04580  
 MBX04590

MACH BOX INTERFERENCE PROGRAM

CODE	BOX CODE		
50	MPL = MFL - NPINT + 2		MBX04600
	MPL = MAX0 (MPL,MIW)		MBX04610
	MPR = MOW		MBX04620
	KODE = 3		MBX04630
	IF(ABS(FANG).LE. 5.0) KODE = 1		MBX04640
55	CALL SSPHI(MPL,MPR)		MBX04650
60	IF(MFL + 1 - MIT)70,70,80		MBX04660
70	MPL = MIT		MBX04670
	MPR = MGT		MBX04680
	KODE = 1		MBX04690
	IF(ABS(FANG).LE. 5.0) GO TO 75		MBX04700
72	KODE = 6		MBX04710
	CALL SSPHI(MPL,MPR)		MBX04720
	KODE = 4		MBX04730
75	CALL SSPHI(MPL,MPR)		MBX04740
80	IF(MGT - MTD)90,100,100		MBX04750
90	MPL = MGT + 1		MBX04760
	MPR = MTD		MBX04770
	KODE = 2		MBX04780
	IF(ABS(FANG).LE. 5.0) GO TO 95		MBX04790
92	KODE = 6		MBX04800
	CALL SSPHI(MPL,MPR)		MBX04810
	KODE = 5		MBX04820
95	CALL SSPHI(MPL,MPR)		MBX04830
100	IF(MTD - MWD)110,1000,1000		MBX04840
110	MPL = MTD + 1		MBX04850
	MPR = MWD		MBX04860
	KODE = 2		MBX04870
	CALL SSPHI(MPL,MPR)		MBX04880
1000	CONTINUE		MBX04890
	RETURN		MBX04900
	END		MBX04910

MACH BOX INTERFERENCE PROGRAM

```

SSPHI      LOGICAL FLOW

C          TO ACCOUNT FOR LAG IN SIGNALS BETWEEN PLANES THAT INTERSECT AT AN MBX04930
C          ACUTE ANGLE THAT RESULTS FROM THE FINITE GRID SIZE, CERTAIN NOTED MBX04940
C          MODIFICATIONS TO THE EQUATIONS FOR THE WING DIAFRAM INTERFERENCE MBX04950
C          REGION (KODE = 6) WERE MADE TO SSN, PHIP, AND PHIN. MBX04960
C          SUBROUTINE SSPHI(MPL,MPR) MBX04970
C          COMMON/A/DPHI(20,60,2),SSU(20,60,2),THW(10),THT(10) MBX04980
C          COMMON/E/BOXL,BOXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBOX,CAPPHI(20,20,2) MBX04990
C          COMMON/F/MODE,KODE,NP,MP,MFL,NMAX,NBW,NBT,NPINT,NPFLET,NPOLW,NPOLI MBX05000
C          DO 100 MP = MPL,MPR MBX05010
C          MP = MP MBX05020
C          MM = MP MBX05030
C          GO TO (10,20,10, 5,15,20),KODE MBX05040
C          5 MM = MP + NMAX - MFL MBX05050
C          10 IF(DPHI(NP,MM,1))100,25,100 MBX05060
C          15 MM = MP + NMAX -MFL MBX05070
C          20 IF(SSU(NP,MM,1))100,25,100 MBX05080
C          25 GO TO (55,55,50,50,55,55),KODE MBX05090
C          50 CALL SSN MBX05100
C          55 CALL PHIP MBX05110
C          GO TO (70,65,60,70,65,65),KODE MBX05120
C          60 CALL PHIN MBX05130
C          GO TO (70,65,70,70,65,65),KODE MBX05140
C          65 CALL SSN MBX05150
C          GO TO 100 MBX05160
C          70 DPHI(NP,MM,1) = DPHI(NP,MM,1) * BOXW MBX05170
C          DPHI(NP,MM,2) = DPHI(NP,MM,2) * BOXW MBX05180
C          CALL GENFOR MBX05190
C          100 CONTINUE MBX05200
C          RETURN MBX05210
C          END MBX05220

```

MACH BOX INTERFERENCE PROGRAM

SSN CORRECTS THE DIAFRAM AND INTERFERENCE SOURCE STRENGTHS

```

SUBROUTINE SSN
COMMON/A/DPHI(20,60,2),SSU(20,60,2),SSL(20,60,2),THW(10),THT(10)
COMMON/B/MIBW(20),MGBW(20),MIBT(20),MGBT(20),MGBT(20),MGBWD(20)
COMMON/C/BARNU,BARMU,EL,CFR,CFI,CNR,CNI,EKBAR,EKM,P(10),W(10)
COMMON/E/BOXL,BOXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBOX,CAPPHI(20,20,2)
COMMON/F/MODE,KODE,NP,MP,MFL,NMAX,NBW,NBT,NPINT,NPFLET,NPOLW,NPOLT
DSSR = 0.0
DSSI = 0.0
GO TO (100,90,30,10,90,30),KODE
CALCULATE INDUCED VELOCITY SUM ON TIP PLANE FROM WING PLANE
DO 23 IC=2,NP
  BARNU = IC - 1
  NU = NP - IC + 1
  YMTFL = FLOAT(MP-1) - YNFL
  XCLB = FLOAT(MFL) + YMTFL * CFA
  IND = 1
  IF(XCLB - (BARNU + 0.5))11,12,12
11 IND = 2
12 JCL = 1
  JCR = MGBW(NU)
  DO 23 JC=JCL,JCR
    YMWFL = YNFL - FLOAT(JC-1)
    BARMU = YMTFL * CFA + YMWFL
    EL = YMTFL * SFA
    DO 19 K = 1,IND
      IF((BARNU+0.5)**2 - (ABS(BARMU)-0.5)**2 - EL**2 - 1.E-6)20,20,15
20 IF(ABS(BARMU).GE.0.5 .OR.ABS(EL).GE.BARNU+0.5) GO TO 22
15 CALL VIC
    CALL KOMPLX(DSSR,DSSI,CNR,CNI,SSU(NU,JC,1),SSU(NU,JC,2))
22 IF(JC - 1)23,23,19
19 BARMU = XCLB - 0.5 + FLOAT(JC-1)
23 CONTINUE
GO TO 70
C CALCULATE INDUCED VELOCITY SUM ON WING PLANE FROM TIP PLANE
30 IF(NPINT-1) 50, 50,31
31 NPXCL = NPINT - MFL
  IND = 1

```

MBX05240  
 MBX05250  
 MBX05260  
 MBX05270  
 MBX05280  
 MBX05290  
 MBX05300  
 MBX05310  
 MBX05320  
 MBX05330  
 MBX05340  
 MBX05350  
 MBX05360  
 MBX05370  
 MBX05380  
 MBX05390  
 MBX05400  
 MBX05410  
 MBX05420  
 MBX05430  
 MBX05440  
 MBX05450  
 MBX05460  
 MBX05470  
 MBX05480  
 MBX05490  
 MBX05500  
 MBX05510  
 MBX05520  
 MBX05530  
 MBX05540  
 MBX05550  
 MBX05560  
 MBX05570  
 MBX05580  
 MBX05590  
 MBX05600

MACH BOX INTERFERENCE PROGRAM

SSN CORRECTS THE DIAFRAM AND INTERFERENCE SOURCE STRENGTHS

```

33 IF(MP - NPXCL)33,33,34
34 IND= 2
34 DO 43 IC = 2,NPINT
  BARNU = IC - 1
  NU = NP - IC + 1
  JCL = MFL + 1
  JCR = MBTD(NU)
  DO 43 JC=JCL,JCR
    YMWFL =-FLOAT(MP-1) + YNFL
    YMTFL =-YNFL + FLOAT(JC-1)
    TRE = 1 - KODE/6
    SIX = KODE/3 - 1
    MPN = JC + NMAX - MFL
    DO 43 K = 1,IND
      BARMU = YMWFL * CFA      + YMTFL
      EL = YMWFL * SFA
      IF((BARNU+0.5)**2 - (ABS(BARMU)-0.5)**2 - EL**2 - 1.E-6)40,40,35
40 IF(ABS(BARMU).GE.0.5 .OR.ABS(EL).GE.BARNU+0.5) GO TO 43
35 CALL VIC
    SSR = TRE * SSU(NU,MPN,1) + SIX * SSL(NU,MPN,1)
    SSI = TRE * SSU(NU,MPN,2) + SIX * SSL(NU,MPN,2)
    CALL KOMPLX(DSSR,DSSI,CNR ,CNI ,SSR,SSI)
43 YMWFL = FLOAT(MP-1) + YNFL
50 GO TO (100,90,70,70,90,90),KODE
C  CORRECT SURFACE SOURCE STRENGTH FOR INTERFERENCE EFFECTS
70 TRE = 1 - KODE/4
  SIX = KODE/2 - 1
  MPN = MP + NMAX - MFL
  MTM = IFIX(TRE+0.5)*MP + IFIX(SIX+0.5)*MPN
  SSU(NP,MTM,1) = SSU(NP,MTM,1) - DSSR
  SSU(NP,MTM,2) = SSU(NP,MTM,2) - DSSI
  SSL(NP,MTM,1) = SSL(NP,MTM,1) + SIX * DSSR
  SSL(NP,MTM,2) = SSL(NP,MTM,2) + SIX * DSSI
  IF(KODE.NE.4) GO TO 100
  YT = FLOAT(MP - MFL) - 0.5
  EL = YT*SFA
  IF(ABS(EL).GE.0.5) GO TO 100

```

MBX05610  
 MBX05620  
 MBX05630  
 MBX05640  
 MBX05650  
 MBX05660  
 MBX05670  
 MBX05680  
 MBX05690  
 MBX05700  
 MBX05710  
 MBX05720  
 MBX05730  
 MBX05740  
 MBX05750  
 MBX05760  
 MBX05770  
 MBX05780  
 MBX05790  
 MBX05800  
 MBX05810  
 MBX05820  
 MBX05830  
 MBX05840  
 MBX05850  
 MBX05860  
 MBX05870  
 MBX05880  
 MBX05890  
 MBX05900  
 MBX05910  
 MBX05920  
 MBX05930  
 MBX05940  
 MBX05950  
 MBX05960  
 MBX05970

MACH BOX INTERFERENCE PROGRAM

SSN CORRECTS THE DIAFRAM AND INTERFERENCE SOURCE STRENGTHS

```

      BARNU = 0.0
      BARMU = YT*CFA - YT
      CALL COREC
      GO TO 100
C
      CALCULATE THE DIAFRAM SOURCE STRENGTH
      90 TRE = 1 - KODE/6 - 2/KODE
      SIX = KODE/3 - 1 + 4/KODE
      MPN = MP + NMAX - MFL
      MTM = IFIX(SIX*0.5)*MP + IFIX(TRE*0.5)*MPN
      DEN = CAPPHI(1,1)**2 + CAPPHI(1,2)**2
      DELSSR = 0.0
      DELSSI = 0.0
      BR = 0.5 * DPHI(NP,MTM,1)
      BI = 0.5 * DPHI(NP,MTM,2)
      CR = CAPPHI(1,1)/DEN
      CI = -CAPPHI(1,2)/DEN
      CALL KOMPLX(DELSSR,DELSSI,BR,BI,CR,CI)
      INTERFERENCE CORRECTIONS TO WING DIAFRAM BOXES ARE MADE BY ADDING
      C THE INDUCED VELOCITY TO THE LOWER SIDE WHEN SATISFYING DPHI = 0.
      C THIS IS EQUIVALENT TO EQ 83 WHEN THE DPHI SUM IS MODIFIED
      80 SSU(NP,MTM,1) = -DELSSR
      SSU(NP,MTM,2) = -DELSSI
      SSL(NP,MTM,1) = DELSSR - FLOAT(KODE/6) * DSSR
      SSL(NP,MTM,2) = DELSSI - FLOAT(KODE/6) * DSSI
      IF(KODE.NE.5) GO TO 99
      YT = FLOAT(MP - MFL) - 0.5
      EL = YT*SFA
      IF(ABS(EL).GE.0.5) GO TO 99
      BARNU = 0.0
      BARMU = YT*CFA - YT
      CALL COREC
      99 DPHI(NP,MTM,1) = 0.0
      DPHI(NP,MTM,2) = 0.0
      100 RETURN
      END

```

MBX05980  
 MBX05990  
 MBX06000  
 MBX06010  
 MBX06020  
 MBX06030  
 MBX06040  
 MBX06050  
 MBX06060  
 MBX06070  
 MBX06080  
 MBX06090  
 MBX06100  
 MBX06110  
 MBX06120  
 MBX06130  
 MBX06140  
 MBX06150  
 MBX06160  
 MBX06170  
 MBX06180  
 MBX06190  
 MBX06200  
 MBX06210  
 MBX06220  
 MBX06230  
 MBX06240  
 MBX06250  
 MBX06260  
 MBX06270  
 MBX06280  
 MBX06290  
 MBX06300  
 MBX06310  
 MBX06320



MACH BOX INTERFERENCE PROGRAM

PHIP CALCULATES PLANAR PORTION OF THE VELOCITY POTENTIAL

```

SUBROUTINE PHIP
COMMON/A/DPHI(20,60,2),SSU(20,60,2),SSL(20,60,2),THW(10),THT(10)
COMMON/B/MIBW(20),MGBW(20),MIBT(20),MGBT(20),MGBTD(20),MGBWD(20)
COMMON/E/BGXL,BGXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBQX,CAPPHI(20,20,2)
COMMON/F/MODE,KODE,NP,MP,MFL,NMAX,NBW,NBT,NPINT,NPFLET,NPQLW,NPQLT
IT = KODE/4 - KODE/6
IW = 1 - IT
N = NP*IW + NPINT*IT
M = MP + (NMAX-MFL)*IT
ML = 1 + MFL*IT
DO 20 IC = 1,N
  NU = NP - IC + 1
  MR = MGBWD(NU)*IW + MGBTD(NU)*IT
  JCL = MAXO(ML,MP-IC+1)
  JCR = MINO(MR,MP+IC-1)
  DO 20 JC = JCL,JCR
    MU = JC + (NMAX-MFL)*IT
    MUBAR = IABS(MP-JC) + 1
    DSSR = SSU(NU,MU,1) - SSL(NU,MU,1)
    DSSI = SSU(NU,MU,2) - SSL(NU,MU,2)
    IF(KODE.NE.6) GO TO 15
    THE LOWER WING AND UPPER WING DIAFRAM BOXES ARE USED TO MODIFY
    THE PLANAR PORTION OF DPHI FOR THE WING DIAFRAM BOXES.
    DSSR = -2.0*SSL(NU,MU,1)
    DSSI = -2.0*SSL(NU,MU,2)
    IF(JC.LE.MFL)GO TO 15
    DSSR = 2.0*SSU(NU,MU,1)
    DSSI = 2.0*SSU(NU,MU,2)
    15 DPR = CAPPHI(IC,MUBAR,1)
    DPI = CAPPHI(IC,MUBAR,2)
    IF(JC.EQ.1.OR .JC.GT.1+IC-MP) GO TO 20
    MUBAR = MP + JC - 1
    DPR = CAPPHI(IC,MUBAR,1) + DPR
    DPI = CAPPHI(IC,MUBAR,2) + DPI
    20 CALL KOMPLX(DPHI(NP,M,1),DPHI(NP,M,2),DPR,DPI,DSSR,DSSI)
    100 RETURN
END

```

MBX06340  
 MBX06350  
 MBX06360  
 MBX06370  
 MBX06380  
 MBX06390  
 MBX06400  
 MBX06410  
 MBX06420  
 MBX06430  
 MBX06440  
 MBX06450  
 MBX06460  
 MBX06470  
 MBX06480  
 MBX06490  
 MBX06500  
 MBX06510  
 MBX06520  
 MBX06530  
 MBX06540  
 MBX06550  
 MBX06560  
 MBX06570  
 MBX06580  
 MBX06590  
 MBX06600  
 MBX06610  
 MBX06620  
 MBX06630  
 MBX06640  
 MBX06650  
 MBX06660  
 MBX06670  
 MBX06680  
 MBX06690  
 MBX06700



MACH BOX INTERFERENCE PROGRAM

PHIN CALCULATES NONPLANAR PORTION OF THE VELOCITY POTENTIAL

```

SUBROUTINE PHIN
COMMON/A/DPHI(20,60,2),SSU(20,60,2),SSL(20,60,2),THW(10),THT(10)
COMMON/B/MIBW(20),MGBW(20),MIBT(20),MGBT(20),MGBTD(20),MGBWD(20)
COMMON/C/BARNU,BARMU,EL,CFR,CFI,CNR,CNI,EKBAR,EKM,P(10),W(10)
COMMON/E/BGXL,BGXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBOX,CAPPHI(20,20,2)
COMMON/F/MODE,KODE,NP,MP,MFL,NMAX,NBW,NBT,NPINT,NPFLET,NPOLW,NPOLTMBOX
THE MODIFIED WING DIAFRAM DPHI SUM CONTAINS NO OUT-OF-PLANE TERMS
IF(NPINT.LE.1)RETURN
NPXCL = NPINT - MFL
IND = 1
IF(MP.LE.NPXCL) IND = 2
DO 80 IC = 2,NPINT
  BARNU = IC - 1
  NU = NP - IC + 1
  JCL = MFL + 1
  JCR = MGBTD(NU)
  DO 80 JC = JCL,JCR
    YMWFL = YNFL - FLOAT(MP-1)
    YMTFL = FLOAT(JC-1) - YNFL
    MM = JC + NMAX - MFL
    DO 80 K = 1,IND
      BARMU = YMWFL * CFA + YMTFL
      EL = YMWFL * SFA
      IF((BARNU+0.5)**2 - (ABS(BARMU)-0.5)**2 - EL**2.LE.1.E-6)GO TO 80
      CALL VPIC
      CALL KEMPLX(DPHI(NP,MP,1),DPHI(NP,MP,2),CFR,CFI,SSU(NU,MM,1),
1      SSU(NU,MM,2))
      80 YMWFL = YNFL + FLOAT(MP-1)
100 RETURN
END

```

C

MACH BOX INTERFERENCE PROGRAM

GFOR GENERALIZED FORCES

```

SUBROUTINE GENFOR
COMMON/A/DPHI(20,60,2),SSU(20,60,2),SSL(20,60,2),THW(10),THT(10)
COMMON/D/XTE(40),YTE(40),QR(21,2),QI(21,2),CGW(10,21),COT(10,21)
COMMON/E/BOXL,BOXW,EMACH,EK,YNFL,FANG,CFA,SFA,NBOX,CAPPHI(20,20,2)
COMMON/F/MODE,KODE,NP,MP,MFL,NMAX,NBW,NBT,NPINT,NPFLET,NPQLW,NPQLT
CON = 1.0
K = 1
IF(MP-1) 20,10,20
10 CON = 0.5
20 WGT = XTE(MP) - FLOAT(NP-1)
IF(WGT -1.5)30,40,40
30 CON = CON * WGT
K = 2
40 NS = 1
MM = MP
NPQL = NPQLW
GO TO (60,200,60,50,200,200),KODE
50 NS = 2
MM = MP + NMAX - MFL
NPQL = NPQLT
60 CONTINUE
X = (FLOAT(NP) - 0.5)*BOXL
Y = FLOAT(MP-1)*BOXW
DPH1R = DPHI(NP,MM,1)
DPH1I = DPHI(NP,MM,2)
I = 0
DZ = 0.0
J = 0
DO 80 N = J,NPQL
DO 80 M = J,N
NM = N - M
FN = NM
I = I + 1
YM = 1.0
IF(M)62,64,62
62 YM = Y**M
54 XM = 1.0

```

MACH BOX INTERFERENCE PROGRAM

GFOR GENERALIZED FORCES

```

        IF(NM)66,68,66
        66 XM = X**NM
           DZ = FN * XM/X * YM
        68 Z = XM * YM
           CALL KOMPLX(QR(I,NS),QI(I,NS),DPHIR*CON,DPHII*CON,DZ,-EK*Z/BOXL)
        80 CONTINUE
           GO TO (200,90),K
        90 CON = 1.0
           IF(MP - 1) 110,100,110
        100 CON = 0.5
        110 DX = XTE(MP) + 2.5 - FLOAT(NP)
           X = XTE(MP) *BOXL
           Y = YTE(MP) *BOXL
           A1 = DPHI(NP-2,MM,1)
           IF(A1.EQ.0.0) GO TO 190
           A2 = -1.5*A1 + 2.0*DPHI(NP-1,MM,1) - 0.5*DPHIR
           A3 = 0.5*A1 - DPHI(NP-1,MM,1) + 0.5*DPHIR
           DPHIR = A1 + A2*DX + A3*DX*DX
           A1 = DPHI(NP-2,MM,2)
           A2 = -1.5*A1 + 2.0*DPHI(NP-1,MM,2) - 0.5*DPHII
           A3 = 0.5*A1 - DPHI(NP-1,MM,2) + 0.5*DPHII
           DPHII = A1 + A2*DX + A3*DX*DX
        190 I = 0
           DO 180 N = J,NP0L
           DO 180 M = J,N
             NM = N - M
             I = I+1
             YM = 1.0
             IF(M)120,130,120
        120 YM = Y**M
        130 Z = X**NM * YM
           QR(I,NS) = QR(I,NS) - DPHIR * Z * CON / BOXL
           QI(I,NS) = QI(I,NS) - DPHII * Z * CON / BOXL
        180 CONTINUE
        200 CONTINUE
           RETURN
           END
MBX07400
MBX07410
MBX07420
MBX07430
MBX07440
MBX07450
MBX07460
MBX07470
MBX07480
MBX07490
MBX07500
MBX07510
MBX07520
MBX07530
MBX07540
MBX07550
MBX07560
MBX07570
MBX07580
MBX07590
MBX07600
MBX07610
MBX07620
MBX07630
MBX07640
MBX07650
MBX07660
MBX07670
MBX07680
MBX07690
MBX07700
MBX07710
MBX07720
MBX07730
MBX07740
MBX07750
MBX07760

```

MACH BOX INTERFERENCE PROGRAM

VPIC VELOCITY POTENTIAL INFLUENCE COEFFS.

C VELOCITY POTENTIAL INFLUENCE COEFFICIENT AT A POINT  
C DUE TO A RECTANGULAR SOURCE SHEET OF UNIT STRENGTH THAT LIES  
C AT LEAST PARTIALLY WITHIN THE FORE MACH CONE FROM THAT POINT.  
C THE VPIC MAY BE FOR A POINT THAT IS EITHER COPLANAR OR  
C NONCOPLANAR WITH THE SOURCE SHEET.  
C

SUBROUTINE VPIC

COMMON/C/BARNU,BARMU,EL,CFR,CFI,CNR,CNI,EKBAR,EKM,P(10),W(10)

CFR = 0.0

CFI = 0.0

ELS = EL\*\*2

ETAR = ABS(BARMU) - 0.5

ETAL = ETAR + 1.0

XR = SQRT(ETAR\*\*2 + ELS)

XL = SQRT(ETAL\*\*2 + ELS)

XU = BARNU + 0.5

XB = BARNU - 0.5

XIR = AMAX1(XB,AMINI(XU,XR))

XIL = AMAX1(XB,AMINI(XU,XL))

XILO = XIR

IF(ETAR.LT.0.0) XILO = AMAX1(ABS(EL),XB)

DO 42 K = 1,3

GO TO (6,8,10),K

6 XI = XILO

DX = XIR - XILO

IF(DX)42,42,12

8 XI = XIR

DX = XIL - XIR

IF(DX)42,42,12

10 XI = XIL

DX = XU - XIL

IF(DX)42,42,12

12 DO 40 I = 1,5

X = XI + P(I)\*DX

GC = DX\*W(I)

Y = SQRT(ABS(X\*\*2 - ELS))

WR = EKBAR \* X

MBX07780  
MBX07790  
MBX07800  
MBX07810  
MBX07820  
MBX07830  
MBX07840  
MBX07850  
MBX07860  
MBX07870  
MBX07880  
MBX07890  
MBX07900  
MBX07910  
MBX07920  
MBX07930  
MBX07940  
MBX07950  
MBX07960  
MBX07970  
MBX07980  
MBX07990  
MBX08000  
MBX08010  
MBX08020  
MBX08030  
MBX08040  
MBX08050  
MBX08060  
MBX08070  
MBX08080  
MBX08090  
MBX08100  
MBX08110  
MBX08120  
MBX08130  
MBX08140

MACH BOX INTERFERENCE PROGRAM

VPIC VELOCITY POTENTIAL INFLUENCE COEFFS.

```

IF(K.EQ.1) GO TO 16
DV = GC * GO(Y,ETAR,ETAL,EKM)/3.14159265
GO TO 18
16 DV = - GC * ZJ(Y*EKM)
18 CFR = CFR + COS(WR)*DV
40 CFI = CFI - SIN(WR)*DV
42 CONTINUE
RETURN
END

```

MBX08150  
MBX08160  
MBX08170  
MBX08180  
MBX08190  
MBX08200  
MBX08210  
MBX08220  
MBX08230

MACH BOX INTERFERENCE PROGRAM

NVIC VELOCITY INFLUENCE COEFFS.

NORMAL, HORIZONTAL, AND VERTICAL VELOCITY INFLUENCE COEFFICIENTS  
AT A POINT DUE TO AN OUT-OF-PLANE UNIT STRENGTH RECTANGULAR  
SOURCE SHEET THAT LIES AT LEAST PARTIALLY WITHIN THE FORE MACH  
CONE FROM THAT POINT.

SUBROUTINE VIC

COMMON/C/BARNU, BARMU, EL, CFR, CFI, CNR, CNI, EKBAR, EKM, P(10), W(10)  
COMMON/E/BOXL, BOXW, EMACH, EK, YNFL, FANG, CFA, SFA, NBOX, CAPPHI(20, 20, 2)

VR = 0.0

VI = 0.0

HR = 0.0

HI = 0.0

B = SIGN(1.0, BARMU)

ELS = EL\*2

ETAR = ABS(BARMU) - 0.5

ETAL = ETAR + 1.0

ETRS = ETAR\*2

ETLS = ETAL\*2

XR = SQRT(ETRS + ELS)

XL = SQRT(ETLS + ELS)

XU = BARNU + 0.5

XB = BARNU - 0.5

XIR = AMAX1(XB, AMINI(XU, XR))

XIL = AMAX1(XB, AMINI(XU, XL))

XILO = XIR

IF(ETAR.LT.0.0) XILO = AMAX1(ABS(EL), XB)

XI = XU

WR = -1.0

DO 10 I = 1, 2

YS = ABS(XI\*2 - ELS)

Y = SQRT(YS)

S = EKBAR \* XI

C = COS(S)

S = SIN(S)

DV = GO(Y, ETAR, ETAL, EKM)/XI

VR = VR + WR\*C\*DV

VI = VI - WR\*S\*DV

MBX08250  
MBX08260  
MBX08270  
MBX08280  
MBX08290  
MBX08300  
MBX08310  
MBX08320  
MBX08330  
MBX08340  
MBX08350  
MBX08360  
MBX08370  
MBX08380  
MBX08390  
MBX08400  
MBX08410  
MBX08420  
MBX08430  
MBX08440  
MBX08450  
MBX08460  
MBX08470  
MBX08480  
MBX08490  
MBX08500  
MBX08510  
MBX08520  
MBX08530  
MBX08540  
MBX08550  
MBX08560  
MBX08570  
MBX08580  
MBX08590  
MBX08600  
MBX08610

MACH BOX INTERFERENCE PROGRAM

NVIC VELOCITY INFLUENCE COEFFS.

```

DH = F0(YS,ETRS,ETLS,EKM)/XI
HR = HR + WR*C*DH
HI = HI - WR*S*DH
XI = XIL0
IF(XI.NE.XB) GO TO 12
10 WR = -WR
12 DO 28 J = 1,3
   GO TO(14,16,18),J
14 DX = XIR - XI
   IF(DX)28,28,20
16 XI = XIR
   DX = XIL - XI
   IF(DX)28,28,20
18 XI = XIL
   DX = XU - XI
   IF(DX)28,28,20
20 DO 26 I = 1,5
   X = XI + DX*P(I)
   GC = DX*W(I)/X**2
   YS = ABS(X**2 - ELS)
   Y = SQR(YS)
   WR = X*EKBAR
   C = COS(WR)
   S = SIN(WR)
   IF(J.EQ.1) GO TO 22
   DV = GC*G0(Y,ETAR,ETAL,EKM)
   DH = GC*F0(YS,ETRS,ETLS,EKM)
   HR = HR - (C + WR*S)*DH
   HI = HI + (S - WR*C)*DH
   GO TO 24
22 DV = -3.14159265 *GC*ZJ(EKM*Y)
24 VR = VR - (C + WR*S)*DV
26 VI = VI + (S - WR*C)*DV
28 CONTINUE
   VR = VR*EL
   VI = VI*EL
   IF(EKM.EQ.0.0)GO TO 30

```

MBX08620  
 MBX08630  
 MBX08640  
 MBX08650  
 MBX08660  
 MBX08670  
 MBX08680  
 MBX08690  
 MBX08700  
 MBX08710  
 MBX08720  
 MBX08730  
 MBX08740  
 MBX08750  
 MBX08760  
 MBX08770  
 MBX08780  
 MBX08790  
 MBX08800  
 MBX08810  
 MBX08820  
 MBX08830  
 MBX08840  
 MBX08850  
 MBX08860  
 MBX08870  
 MBX08880  
 MBX08890  
 MBX08900  
 MBX08910  
 MBX08920  
 MBX08930  
 MBX08940  
 MBX08950  
 MBX08960  
 MBX08970  
 MBX08980



MACH BOX INTERFERENCE PROGRAM

NVIC VELOCITY INFLUENCE COEFFS.

```

HR = HR/EKM
HI = HI/EKM
GO TO 36
30 FTOP = 1.0
FBOT = 1.0
IF( XL.LT.XU) FTOP =(XU/ XL) + SQRT((XU/ XL)**2 - 1.0)
IF( XR.LT.XU) FBOT =(XU/ XR) + SQRT((XU/ XR)**2 - 1.0)
IF( XL.LT.XB) FBOT = FBOT*((XB/ XL) + SQRT((XB/ XL)**2 - 1.0))
IF( XR.LT.XB) FTOP = FTOP*((XB/ XR) + SQRT((XB/ XR)**2 - 1.0))
HR = ALOG(FTOP/FBOT)
36 CNR = (VR*CFA + B*HR*SFA)/3.1415926
CNI = (VI*CFA + B*HI*SFA)/3.1415926
RETURN
END

```

MBX08990  
 MBX09000  
 MBX09010  
 MBX09020  
 MBX09030  
 MBX09040  
 MBX09050  
 MBX09060  
 MBX09070  
 MBX09080  
 MBX09090  
 MBX09100  
 MBX09110  
 MBX09120

MACH BOX INTERFERENCE PROGRAM

VVEL VERTICAL VELOCITY AND VELOCITY POTENTIAL INTEGRAND

```

FUNCTION G0(R,ETAR,ETAL,EKM)
DIMENSION BSL(23),S(2),C(2),AS(2),SO(2)
S(1)=ETAR
S(2)=ETAL
DO 6 I=1,2
IF (ABS(S(I)).GE.R ) GO TO 4
S(I) = S(I)/R
C(I)=SQRT(1.0-S(I)**2)
AS(I)=2.0*ATAN(S(I)/(1.0+C(I)))
S(I)=2.0*S(I)*C(I)
C(I)=2.0*C(I)**2-1.0
GO TO 6
4 AS(I)=SIGN(1.57079633,S(I))
S(I)=0.0
6 SO(I)=0.0
GO=AS(1)-AS(2)
IF (ABS(G0).LE.1.E-5) GO TO 86
ARG=EKM*R
IF (ARG.EQ.0.0) RETURN
CALL BSLS(BSL,ARG,N)
GO=BSL(1)*GO
F=1.0
FI=1.0
DO 100 J=1,N
GO=BSL(J+1)*(S(1)-S(2))/FI-GO
S(I) IS SIN(2*J*AS(I))
DO 90 I=1,2
S4=2.0*S(I)*C(I)-SO(I)
SO(I)=S(I)
90 S(I)=S4
F=-F
100 FI=FI+1.0
IF (F.LT.0.0) GO=-GO
RETURN
86 GO=0.0
RETURN
END

```

MBX09140  
 MBX09150  
 MBX09160  
 MBX09170  
 MBX09180  
 MBX09190  
 MBX09200  
 MBX09210  
 MBX09220  
 MBX09230  
 MBX09240  
 MBX09250  
 MBX09260  
 MBX09270  
 MBX09280  
 MBX09290  
 MBX09300  
 MBX09310  
 MBX09320  
 MBX09330  
 MBX09340  
 MBX09350  
 MBX09360  
 MBX09370  
 MBX09380  
 MBX09390  
 MBX09400  
 MBX09410  
 MBX09420  
 MBX09430  
 MBX09440  
 MBX09450  
 MBX09460  
 MBX09470  
 MBX09480  
 MBX09490  
 MBX09500

MACH BOX INTERFERENCE PROGRAM

```

HVEL      HORIZONTAL VELOCITY INTEGRAND
C
C      HORIZONTAL VELOCITY INTEGRAND FUNCTION
FUNCTION F0(RS,ETRS,ETLS,EKM)
F0 = 0.0
IF(EKM.LE.0.0) RETURN
FR = RS - ETRS
FL = RS - ETLS
IF(FR.GT.0.0) F0 = SIN(EKM*SQRT(FR))
IF(FL.GT.0.0) F0 = F0 - SIN(EKM*SQRT(FL))
RETURN
END

```

MBX09520  
 MBX09530  
 MBX09540  
 MBX09550  
 MBX09560  
 MBX09570  
 MBX09580  
 MBX09590  
 MBX09600  
 MBX09610  
 MBX09620

MACH BOX INTERFERENCE PROGRAM

ZJ ZERO ORDER BESSEL FUNCTION

```

FUNCTION ZJ(ARG)
  A = -(ARG/2.0)**2
  ZJ = 1.0
  PF = 1.0
  AV = 1.0
  DO 30 K = 1,20
    AN = AN * A/PF**2
    PF = PF + 1.0
    IF (ABS(AN) - 1.E- 5) 40,40,30
  30 ZJ = ZJ + AN
  40 RETURN
  END

```

MBX09640  
 MBX09650  
 MBX09660  
 MBX09670  
 MBX09680  
 MBX09690  
 MBX09700  
 MBX09710  
 MBX09720  
 MBX09730  
 MBX09740  
 MBX09750

MACH BOX INTERFERENCE PROGRAM

```

BSLS      EVEN BESSEL FUNCTIONS

SUBROUTINE BSL(BSL,ARG,N)
DIMENSION BSL(23)
DO 1000 I=1,23
1000 BSL(I)=0.0
IF(ABS(ARG) - 0.10)70,5,5
5 N = MIN1(20.0,(ARG + 10.0))
ARG = ARG ** 2
20 AN = 2 * N + 4
BSL(N+2) = (4.0 * AN * (AN-1.0)/ARGS - (AN-1.0)/AN) * 0.3E-30
PF = 0.0
DO 50 I=0,N
M = N-I+1
AN = 2*M+1
BSL(M) = (4.0*(AN-1.0)*(AN-2.0)/ARGS-(AN-2.0)/AN-1.0)*BSL(M+1)
X -(AN-2.0)/AN*BSL(M+2)
50 PF = PF + 2.0*BSL(M+1)
PF = PF + BSL(I)
AN = 0.0
IF(ABS(PF).GT.1.0) AN = ABS(PF)*1.E-10
M = N+2
DO 60 I = 1,M
IF(AN - ABS(BSL(I)))54,56,56
54 BSL(I) = BSL(I) / PF
GO TO 60
56 BSL(I) = 0.0
60 CONTINUE
DO 180 I = 1,M
J = M - I + 1
IF(ABS(BSL(J)) - 1.E-7)180,180,120
180 N = J - 1
GO TO 120
70 BSL(2) = 0.125 * ARG**2
BSL(1) = 1.0 - 2.0*BSL(2)
N = 2
120 RETURN
END

```

MBX09770  
 MBX09780  
 MBX09790  
 MBX09800  
 MBX09810  
 MBX09820  
 MBX09830  
 MBX09840  
 MBX09850  
 MBX09860  
 MBX09870  
 MBX09880  
 MBX09890  
 MBX09900  
 MBX09910  
 MBX09920  
 MBX09930  
 MBX09940  
 MBX09950  
 MBX09960  
 MBX09970  
 MBX09980  
 MBX09990  
 MBX10000  
 MBX10010  
 MBX10020  
 MBX10030  
 MBX10040  
 MBX10050  
 MBX10060  
 MBX10070  
 MBX10080  
 MBX10090  
 MBX10100  
 MBX10110  
 MBX10120

127

# MACH BOX INTERFERENCE PROGRAM

## GAUSS GAUSSIAN INTEGRATION CONSTANTS

```

BLOCK DATA
COMMON/C/BARNU,BARMU,EL,CFR,CFI,CNR,CNI,EKBAR,EKM,P(10),W(10)
C P(1) IS = 0.5 + ITH ZERO OF LEGENDRE POLYNOMIAL
DATA P/ 0.953089923,0.769234655,0.5,0.230765345,0.046910077,
1 0.0,0.0,0.0,0.0,0.0 /
C W(1) IS THE ITH WEIGHT COEFFICIENT FOR GAUSSIAN QUADRATURE
DATA W/ 0.118463442,0.239314335,0.284444444,0.239314335,
1 0.118463442,0.0,0.0,0.0,0.0,0.0,0.0 /
END

```

MBX10370  
 MBX10380  
 MBX10390  
 MBX10400  
 MBX10410  
 MBX10420  
 MBX10430  
 MBX10440  
 MBX10450

128

# MACH BOX INTERFERENCE PROGRAM

## KPLX COMPLEX MULTIPLICATION

```

SUBROUTINE KOMPLX(AR,AI,BR,BI,CR,CI)
AR=AR+BR*CR-BI*CI
AI=AI+BR*CI+BI*CR
RETURN
END

```

MBX10470  
 MBX10480  
 MBX10490  
 MBX10500  
 MBX10510



MACH BOX INTERFERENCE PROGRAM

LSSUR LEAST SQUARES SURFACE FIT

```

SUBROUTINE LSSUR(X2,Y2,Z2,N,NZ,NZ,ND,NG,A)
COMMON/A/X2(100),Y2(100),Z2(10,100),C(21,100),P1(100,21),
      I      X1F1(21,21),PTP(21,21),PTPINV(21,21),F1(21)
      DIMENSION A(10,21),X1F1C(21,100),X1M(21,21)
      EQUIVALENCE (P1,X1F1C),(PTP,X1M)
      FIND CENTROID X1,Y1
      SUM1=0.0
      SUM2=0.0
      DO 5 I=1,N
      SUM1=SUM1+X2(I)
      SUM2=SUM2+Y2(I)
      FN=N
      XI=SUM1/FN
      YI=SUM2/FN
      SHIFT TO X1,Y1
      DO 7 I=1,N
      X2(I)=X2(I)-XI
      Y2(I)=Y2(I)-YI
      CONTINUE
      NORMALIZE ON LARGEST LENGTH
      BIG=0.0
      DO 9 I=1,N
      BIG = AMAX1(BIG,ABS(X2(I)),ABS (Y2(I)))
      CONTINUE
      DO 16 I=1,N
      X2(I)=X2(I)/BIG
      Y2(I)=Y2(I)/BIG
      CONTINUE
      COMPUTE P MATRIX
      DO 50 L=1,N
      NG=1
      P1(L,1)=1.0
      IF(NG - 2) 60,20,20
      CONTINUE
      DO 40 I=2,ND
      IM1=I-1
      DO 30 K=1,IM1

```

MBX10760  
 MBX10770  
 MBX10780  
 MBX10790  
 MBX10800  
 MBX10810  
 MBX10820  
 MBX10830  
 MBX10840  
 MBX10850  
 MBX10860  
 MBX10870  
 MBX10880  
 MBX10890  
 MBX10900  
 MBX10910  
 MBX10920  
 MBX10930  
 MBX10940  
 MBX10950  
 MBX10960  
 MBX10970  
 MBX10980  
 MBX10990  
 MBX11000  
 MBX11010  
 MBX11020  
 MBX11030  
 MBX11040  
 MBX11050  
 MBX11060  
 MBX11070  
 MBX11080  
 MBX11090  
 MBX11100  
 MBX11110  
 MBX11120

MACH BOX INTERFERENCE PROGRAM

LSSUR LEAST SQUARES SURFACE FIT

```

NG=NG+1
NGMI=NG+1-I
PI(L,NG)=PI(L,NGMI)*X2(L)
30 CONTINUE
NG=NG+1
PI(L,NG)=PI(L,NGMI)*Y2(L)
40 CONTINUE
50 CONTINUE
60 CONTINUE
DO 250 I=1,NG
DO 250 J=1,NG
PTP(I,J)=0.0
DO 230 K=1,N
PTP(I,J)=PTP(I,J)+PI(K,I)*PI(K,J)
230 CONTINUE
250 CONTINUE
CALL MATINV(PTP,PTPINV,NG)
C=(PTPINV)(PT)
DO 130 I=1,NG
DO 130 J=1,N
C(I,J)=0.0
DO 130 K=1,NG
C(I,J)=C(I,J)+PTPINV(I,K)*PI(K,J)
130 CONTINUE
C X1 MATRIX
C INITIALIZE ZEROS AND ONES ON DIAGONAL
DO 610 I=1,21
DO 610 J=1,21
X1M(I,J)=0.0
610 DO 620 I=1,21
X1M(I,I)=1.0
620 C INITIALIZE COLUMN INDEX AND SIGN
NG=1
SIGNR=1.0
C INDEX DEGREE OF FIRST ROW
DO 760 LI=1,ND
SIGNR=-SIGNR

```

MACH BOX INTERFERENCE PROGRAM

LSSUR LEAST SQUARES SURFACE FIT

```

L1P1=L1+1
DO 750 L2=1,L1P1
NG=NG+1
NEXPYC=L2-1
NEXPXC=L1-NEXPYC
FEXPXC=NEXPXC
FEXPYC=NEXPYC
XIM(1,NG)=X1**NEXPXC*Y1**NEXPYC*SIGNR
INDEX DOWN COLUMN THRU (L1-1)DEGREE
L1M1=L1-1
NG1=1
SIGNC=1.0
DO 740 K1=1,L1M1
SIGNC=-SIGNC
INDEX TERMS OF DEGREE(IN COLUMN)
K1P1=K1+1
DO 730 K2=1,K1P1
NG1=NG1+1
NEXPYR=K2-1
NEXPXR=K1-NEXPYR
SUBTRACT EXPONENTS
NXMX=NEXPXC-NEXPXR
NYMY=NEXPYC-NEXPYR
IF(NXMX)730,660,660
IF(NYMY)730,670,670
660 FX=1.0
670 FY=1.0
IF(NXMX)700,700,680
680 FXMX=NXMX
DO 690 K3=1,NXMX
FK3=K3
FX=FX*(FEXPXC-(FK3-1.0))/FK3
CONTINUE
690
700 IF(NYMY)720,720,705
705 FYMY=NYMY
DO 710 K3=1,NYMY
FK3=K3

```

MACH BOX INTERFERENCE PROGRAM

LSSUR LEAST SQUARES SURFACE FIT

```

710  CONTINUE
720  X1M(NG1,NG)=X1**NXMNX*Y1**NYMNY*FX*FY*SIGNR*SIGNC
730  CONTINUE
740  CONTINUE
750  CONTINUE
760  CONTINUE
    C
    F1 MATRIX
    F1(1)=1.0
    L=1
    DO 4010 I=1,ND
    IPI=I+1
    FACTOR=L.O/HIG **I
    DO 4010 K=I,IPI
    L=L+1
    F1(L)=FACTOR
4010 CONTINUE
    C
    X1F1={X1M}*(F1)
    DO 4020 J=1,NG
    DO 4020 I=1,NG
    X1F1(I,J)=X1M(I,J)*F1(J)
4020 CONTINUE
    C
    X1F1C={X1F1}(C)
    DO 4030 I=1,NG
    DO 4030 J=1,N
    X1F1C(I,J)=0.0
    DO 4030 K=1,NG
    X1F1C(I,J)=X1F1C(I,J)+X1F1(I,K)*C(K,J)
    C
    A={X1F1C}(Z)
    DO 4040 I=1,NG
    DO 4040 J=1,NZ
    A(J,I) = 0.0
    DO 4040 K=1,N
    A(J,I) = A(J,I) + X1F1C(I,K)*Z2(K,J)
4040 A(J,I) = A(J,I) + X1F1C(I,K)*Z2(K,J)
    RETURN
    END

```

MBX11870  
 MBX11880  
 MBX11890  
 MBX11900  
 MBX11910  
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 MBX12100  
 MBX12110  
 MBX12120  
 MBX12130  
 MBX12140  
 MBX12150  
 MBX12160  
 MBX12170  
 MBX12180  
 MBX12190  
 MBX12200  
 MBX12210  
 MBX12220

MACH BOX INTERFERENCE PROGRAM

```

MINV      MATRIX INVERSION
SUBROUTINE MATINV(PTP,PTPINV,NG)
MATRIX INVERSION BY GAUSSIAN ELIMINATION
DIMENSION PTP(NG,NG),PTPINV(NG,NG)
SET-UP IDENTITY MATRIX
DO 310 I=1,NG
DO 310 J=1,NG
PTPINV(I,J)=0.0
DO 320 I=1,NG
PTPINV(I,I)=1.0
ENTER OVERALL ELIMINATION LOOP
DO 460 I=1,NG
IPI=I+1
CHECK FOR ZERO IN PIVOTAL ELEMENT
IF(PTP(I,I))355,325,355
IF(IPI-NG)326,326,331
SEARCH FOR NON-ZERO IN SAME COLUMN
DO 330 K = IPI,NG
IF(PTP(K,I))335,330,335
CONTINUE
NON-ZERO PIVOT CANNOT BE FOUND
WRITE (6,332)
FORMAT(1HX,23HDETERMINANT IS SINGULAR)
CALL EXIT
CONTINUE
INTERCHANGE I AND K ROWS IN BOTH MATRICES
DO 340 L = 1,NG
HOLD=PTP(I,L)
PTP(I,L)=PTP(K,L)
PTP(K,L)=HOLD
CONTINUE
DO 350 L=1,NG
HOLD=PTPINV(I,L)
PTPINV(I,L)=PTPINV(K,L)
PTPINV(K,L)=HOLD
CONTINUE
CONTINUE
DIVIDE ROW BY PIVOT

```

MBX12240  
 MBX12250  
 MBX12260  
 MBX12270  
 MBX12280  
 MBX12290  
 MBX12300  
 MBX12310  
 MBX12320  
 MBX12330  
 MBX12340  
 MBX12350  
 MBX12360  
 MBX12370  
 MBX12380  
 MBX12390  
 MBX12400  
 MBX12410  
 MBX12420  
 MBX12430  
 MBX12440  
 MBX12450  
 MBX12460  
 MBX12470  
 MBX12480  
 MBX12490  
 MBX12500  
 MBX12510  
 MBX12520  
 MBX12530  
 MBX12540  
 MBX12550  
 MBX12560  
 MBX12570  
 MBX12580  
 MBX12590  
 MBX12600

MACH BOX INTERFERENCE PROGRAM

```

MINV      MATRIX INVERSION

      PIVOT=PTP(I,I)
      PTP MATRIX
      DO 360 L=I,NG
      PTP(I,L)=PTP(I,L)/PIVOT
      PTPINV MATRIX
      DO 370 L=I,NG
      PTPINV(I,L)=PTPINV(I,L)/PIVOT
      ARE WE AT LAST ELEMENT
      IF(I-NG)375,47C,470
      CONTINUE
      ZERO COLUMN AND ROW OF PTP SIMULTANEOUSLY
      DO 450 K=IPI,NG
      FIND ZEROING FACTOR FOR COLUMN
      C0=PTP(K,I)
      PTP(K,I)=0.0
      MULTIPLY ROW BY ZEROING FACTOR
      DO 390 L=IPI,NG
      PTP(K,L)=PTP(K,L)-C0*PTP(I,L)
      DO 400 L=I,NG
      PTPINV(K,L)=PTPINV(K,L)-C0*PTPINV(I,L)
      CONTINUE
      CONTINUE
      CONTINUE
      J1=NG+1
      J1=J1-1
      IF(J1-1)488,488,474
      J1M1=J1-1
      DO 477 I1=1,J1M1
      C0=PTP(I1,J1)
      DO 477 J2=1,NG
      PTP(I1,J2)=PTP(I1,J2)-C0*PTP(I1,J2)
      PTPINV(I1,J2)=PTPINV(I1,J2)-C0*PTPINV(I1,J2)
      CONTINUE
      GO TO 472
      CONTINUE
      RETURN
      END

```

MBX12610  
 MBX12620  
 MBX12630  
 MBX12640  
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 MBX12670  
 MBX12680  
 MBX12690  
 MBX12700  
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UNCLASSIFIED

Security Classification

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13. ABSTRACT Ashley's approach to mutual interference theory by source superposition methods has been applied to the prediction of supersonic air loads on intersecting thin lifting surfaces in steady of oscillatory motion. Steady loading is regarded as the special case of zero frequency of oscillation. Each surface may be oscillating in a mode of rigid or elastic vibration or linear combinations thereof. Evvard's diaphragm concept has been extended to treat the out-of-plane interference problem. As a result, any leading or side edge on any of the intersecting surfaces may be subsonic. The study reported herein has lead to the formulation of a method by which diaphragm regions can be selected that eliminate the need for calculating out-of-plane velocity potentials. Based on mutual interference theory, the method requires only the calculation of out-of-plane velocities so that tangential flow conditions may be met. The Mach-box method has been used to obtain an approximate solution to the problem. In following the aerodynamic influence coefficient procedure of Zartarian and Hsu, each surface and diaphragm is overlaid with a grid of rectangular Mach boxes, the diagonals of which are parallel to the Mach lines. For purposes of calculating induced velocities and velocity potentials, the source strength over the area of each box is assumed to be constant. Out-of-plane velocity influence coefficients are calculated for the center of each box in the interference region, and the tangential flow condition is satisfied there. For purposes of calculating generalized forces, the resulting velocity potential over the area of each box is also assumed to be constant and equal to the value at the center of the box.		

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