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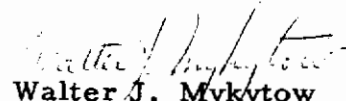
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## ABSTRACT

The fundamental equations of the transonic box method were derived, based on the representation of the velocity potential by a doublet distribution. They form the basis of a systematic method of treating an oscillating wing at  $M = 1$ , analogous to the supersonic Mach box method.

A digital computer program, written in Fortran IV, is presented. The program applies to a planar wing of polygonal planform, with a straight trailing edge, and as many as three sweep angles along the leading edge. For a maximum of ten modes of oscillation, the program computes the oscillatory potentials and pressures and a generalized force matrix.

Results obtained from the program are compared with existing theoretical and experimental values. Several possible extensions of the method are described.

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LIST OF SYMBOLS

<u>Symbol</u>	<u>Definition</u>
a	Local speed of sound; speed of sound at infinity
$a_{nm}$	Coefficient in the potential series
$A(i - i',  j - j' )$	Influence coefficient: the upwash at the center of $B_{ij}$ caused by a unit doublet distribution over $B_{i'j'}$
$A_{jr}$	Term in $\bar{\varphi}$ evaluated at $(x_j, y_j)$
AXY (I, J)	An integral over the wing planform
AY (J)	An integral along the trailing edge
b	Root chord length
B	Region composed of boxes, approximating S
$B_{ij}$	A box
$(B_{rr'})$	Matrix used in least squares surface fits
BXY	Part of AXY (I, J)
BY	Part of AY (J)
$\bar{C}_p$	Pressure coefficient
$(C'_r), (C''_r)$	Column matrices used in least squares surface fits
d	Dimensionless length of box side
$d_{nm}$	Coefficient in deflection polynomial
DA	The data array
f	Function which describes the wing deflection
F	Factor which gives $\bar{\varphi}$ the proper edge behavior

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<u>Symbol</u>	<u>Definition</u>
$\tilde{g}_u, \tilde{g}_l, \tilde{f}$	Functions used in the equations of upper and lower wing surfaces
$h_j$	Weight used in Gaussian quadrature
$i$	$\sqrt{-1}$
$i, j$	Indexes specifying box position
$I, J$	Indexes
$k$	Reduced frequency: $\omega b/U_\infty$
$l$	$kd$
$L_{ij}$	Generalized force coefficient
$M$	Mach number
$n, m$	Indexes equal to power of $x$ and power of $y^2$
NC	Number of coefficients
NP	Number of points
NS	Number of segments of leading edge given by the data
$p, q$	Integration variables
$Q$	Quantity minimized in least squares surface fits
$r$	Index
$S$	Function used in the equation of a surface
$S$	Region in the $xy$ -plane occupied by the wing; the area of this region
$t$	Time
$u, v$	Integration variables
$u_j$	Point used in Gaussian quadrature
$U_\infty$	Air speed of the wing; speed of flow at infinity

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<u>Symbol</u>	<u>Definition</u>
w	Upwash at $z = 0+$
W	The region of the xy-plane occupied by the wing's wake
$\tilde{x}, \tilde{y}, \tilde{z}$	Coordinates with dimensions of length
x, y, z	Dimensionless coordinates
$(x_i, y_j)$	Center of $B_{ij}$
$(x_j, y_j)$	Point at which a value of potential or deflection is given
$\left. \begin{array}{l} x_1, \dots, x_{NS} \\ y_0, \dots, y_{NS} \end{array} \right\}$	Coordinates of points on the leading edge given by the data
$x_0$	Function which describes the leading edge: $x = x_0(y)$
$y_{max}$	Value of y at the wing tip
$y_+, y_-$	Limits of integration
$\alpha_{nm}, \alpha_r$	Real part of $a_{nm}$
$\beta_r$	Imaginary part of $a_{nm}$
$\delta$	Constant factor in the deflection
$\Delta p_i$	Lifting pressure in the ith mode
$\xi, \eta$	Integration variables equivalent to x, y
$\nu$	Frequency
$\rho$	Density
$\rho$	Source or doublet strength
$\sigma$	The integral over x involved in BXY
$\Phi$	Velocity potential
$\phi$	Steady <del>perturbation</del> <sup>flow</sup> potential
$\phi$	Unsteady perturbation potential

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<u>Symbol</u>	<u>Definition</u>
$\bar{\varphi}$	Time independent factor of $\varphi$
$\bar{\varphi}_0$	Potential of a point source
$\bar{\varphi}_1$	Potential of a point doublet
$\bar{\varphi}_s$	Potential of a source distribution
$\bar{\varphi}_d$	Potential of a doublet distribution
$\bar{\varphi}_{ij}$	Value of $\bar{\varphi}$ in $B_{ij}$
$\bar{\varphi}'_j$	Real part of value of $\bar{\varphi}$ at $(x_j, y_j)$
$\Psi$	Upwash in the xy-plane caused by a point doublet
$\omega$	Angular frequency, $2\pi\nu$

## 1. INTRODUCTION

The transonic box program is designed to calculate the unsteady potentials for a given set of modes of wing oscillation and to compute the generalized forces. Pressure distributions may be obtained from the potentials.

A planar wing with a straight trailing edge is assumed. The oscillations are assumed to be symmetric in the spanwise coordinate  $y$ . None of these assumptions is necessary for the method. (See Section 5.)

The basic step in the box method is the solution of the system of simultaneous equations [Equation (33)] which determine a set of values of potential on the wing from a corresponding array of upwash values. A surface is fitted to these values, giving a functional representation of the potential that is used subsequently to find pressures and generalized forces.

The method used is suggested by the success of supersonic box methods (References 1 through 4). The potential is generated by a doublet distribution rather than by a source distribution because the latter method would involve diaphragm regions of infinite extent, whereas the doublet distribution is confined to the wing and its wake.

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## 2. THEORETICAL DEVELOPMENT OF THE METHOD

### 1. THE DIFFERENTIAL EQUATION

We consider an oscillating body moving at speed  $U_\infty$  through a nonviscous fluid. From the point of view of a moving coordinate system  $(\bar{x}, \bar{y}, \bar{z})$  in which the average position of the body is fixed, there is a flow past the body with velocity  $U_\infty$  at infinity. Assume that the flow is irrotational; then the velocity field of the flow is the gradient of a potential function  $\Phi$ , which satisfies the differential equation

$$\nabla^2 \Phi - \frac{1}{a^2} \left[ \Phi_{tt} + 2 \nabla \Phi \cdot \nabla \Phi_t + (\nabla \Phi \cdot \nabla) 1/2 (\nabla \Phi)^2 \right] = 0 \quad (1)$$

(See Reference 5, p. 193 where  $a$  is the local speed of sound.

Suppose that the flow is approximately uniform in the direction of the positive  $\bar{x}$ -axis. This may be true, for example, if the body is almost plane and the oscillations are small. Then  $\Phi$  may be broken up into several parts, as

$$\Phi = U_\infty (\bar{x} + \phi + \varphi) \quad (2)$$

where the first term gives a uniform flow, the second term gives the correction for a steady flow about the body, the third term gives the correction to this for the oscillating body, and  $\phi$  and  $\varphi$  are small.

To the first order,  $\phi$  and  $\varphi$  are different solutions of the same differential equation

$$(1 - M^2) \varphi_{\bar{x}\bar{x}} + \varphi_{\bar{y}\bar{y}} + \varphi_{\bar{z}\bar{z}} - \frac{2 U_\infty}{a^2} \varphi_{\bar{x}t} - \frac{1}{a^2} \varphi_{tt} = 0 \quad (3)$$

where  $M, a$  are the Mach number and speed of sound at infinity. (See Reference 5, p. 198.)  $\varphi$  is a periodic function of  $t$ . Since the differential equation is linear, we may put  $\varphi = \bar{\varphi}(x, y, z) e^{i\omega t}$ , where  $\omega$  is the angular frequency of oscillation. In terms of the nondimensional quantities,

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$$\begin{aligned}x &= \bar{x}/b \\y &= \bar{y}/b \\z &= \bar{z}/b \\k &= \omega b/U_\infty\end{aligned}$$

(b is a characteristic length of the body); Equation (3) becomes

$$(1-M^2)\bar{\varphi}_{xx} + \bar{\varphi}_{yy} + \bar{\varphi}_{zz} - 2iM^2k\bar{\varphi}_x + M^2k^2\bar{\varphi} = 0 \quad (4)$$

For  $M = 1$ , this reduces to

$$\bar{\varphi}_{yy} + \bar{\varphi}_{zz} - 2ik\bar{\varphi}_x + k^2\bar{\varphi} = 0 \quad (5)$$

the linearized transonic equation (see Reference 6, p. 7). It has been suggested by Landahl (Reference 6) that the proper equation to use instead of (4) is

$$\bar{\varphi}_{yy} + \bar{\varphi}_{zz} - 2iM^2k\bar{\varphi}_x + M^2k^2\bar{\varphi} = 0$$

if  $k \gg |M-1|$ . Comparison of this equation with (5) leads to a similarity rule for flows in the transonic range (see Reference 6, p. 18).

The range of validity of this equation is discussed by Landahl (Reference 6, Chapter 1). First, there is the requirement for linearization in any speed range, that the perturbation potential  $\phi + \varphi$  be small. This is not satisfied at the leading edge of a wing for any realistic cross-sectional shape; however, it may be satisfied over the rest of the wing, if the wing has small thickness, and the results on parts of the wing away from the leading edge are not much affected by the error there.

Another restriction peculiar to transonic speeds is associated with the absence of the term in  $\bar{\varphi}_{xx}$ . The actual flow has some variation in local Mach number which may influence the nature of the flow considerably if  $M$  is near 1. The presence of the term in  $\bar{\varphi}_x$  tends to reduce this influence, but for  $k$  small or zero, the equation is valid only for a highly swept wing with a pointed nose.

The difference of the local Mach number from the value 1 assumed in Equation (5) may come from two sources: (1) wing thickness, and (2) a change in the free stream Mach number. Thus, for any value of  $k$ , there are limits on the thickness ratio and the Mach number range, which increase with  $k$ . Estimates of these limits are not possible, because of the small amount of experimental data available.

## 2. BOUNDARY CONDITIONS

The solution of Equation (1) must give a velocity field which is such that a particle at the body surface moves along the moving surface. If the equation of the surface is

$$S(\tilde{x}, \tilde{y}, \tilde{z}, t) = 0$$

this equation must be satisfied when  $(\tilde{x}, \tilde{y}, \tilde{z})$  moves with the velocity  $\nabla\Phi$ . Differentiating with respect to  $t$  gives the condition

$$\nabla\Phi \cdot \nabla S + \frac{\partial S}{\partial t} = 0 \quad (6)$$

This determines the normal velocity at the surface.

Now suppose that the body (to be referred to henceforth as a wing) is almost planar, lying almost in the  $xy$ -plane (see Figure 1). For vertical

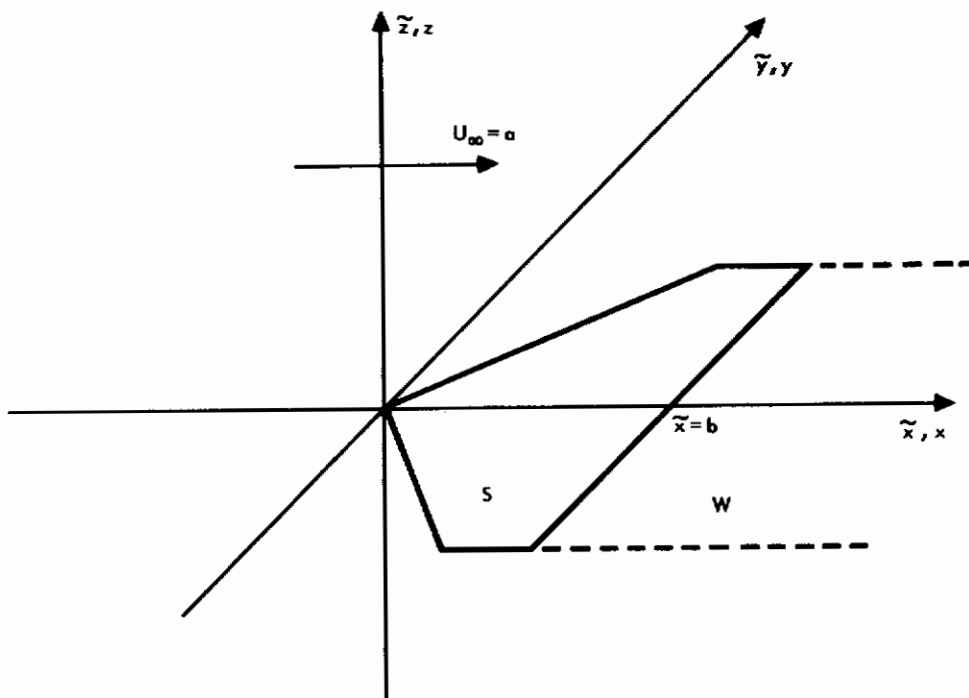


Figure 1. Coordinate Systems

oscillations of the body, the upper and lower surfaces may be represented by the equations

$$\tilde{z} = \tilde{g}_u(\tilde{x}, \tilde{y}) + e^{i\omega t} \tilde{f}(\tilde{x}, \tilde{y})$$

$$\tilde{z} = \tilde{g}_u(\tilde{x}, \tilde{y}) + e^{i\omega t} \tilde{f}(\tilde{x}, \tilde{y})$$

where the functions  $\tilde{g}_u, \tilde{g}_l$  are associated with the deviation of the shape of the body from planar, and  $\tilde{f}$  depends on the mode of oscillation. Then on the two surfaces, we may take

$$S = \tilde{z} - \tilde{g}_u - e^{i\omega t} \tilde{f}$$

$$S = \tilde{z} - \tilde{g}_l - e^{i\omega t} \tilde{f}$$

Use these expressions for S and Equation (2) in Equation (6). Neglecting terms that involve products of  $\varphi$  or  $\phi$  with  $\tilde{g}_u, \tilde{g}_l$ , or  $\tilde{f}$ , the resulting equation may be broken up into a steady part, which gives the boundary condition for  $\phi$ , and an unsteady part, which gives the boundary condition for  $\bar{\varphi}$ . The unsteady part is

$$\frac{\partial \bar{\varphi}}{\partial z} = \frac{\partial f}{\partial x} + ikf \quad (7)$$

where  $f = \tilde{f}/b$ . To the present degree of approximation, this condition should be applied at  $z = 0$ , over the region of the  $xy$ -plane on which the body projects.

### 3. THE BOUNDARY VALUE PROBLEM FOR $\bar{\varphi}$

In linearized theory, a disturbance of a flow at Mach 1 does not have any influence upstream. Consequently,

$$\bar{\varphi}(x, y, z) = 0, \quad x < 0 \quad (8)$$

if the body lies in the region  $x \geq 0$ . This is one of the conditions  $\bar{\varphi}$  must satisfy.

$\bar{\varphi}$  is a solution of Equation (5) in all space outside S and W, the regions in the  $xy$ -plane occupied by the wing and its wake (see Figure 1). In general,  $\bar{\varphi}$  is discontinuous in these regions. A boundary condition on W is obtained by equating the pressures above and below the surface of the wake. From the linearized form of the pressure coefficient,

$$\bar{C}_p = -2(\bar{\varphi}_x + ik\bar{\varphi})$$

(see Reference 6, p. 15) we get



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$$\left[ \bar{\varphi}_x(x, y, z) + ik\bar{\varphi}(x, y, z) \right] \Big|_{z=0-}^{0+} = 0, \quad (x, y) \text{ in } W \quad (9)$$

This condition, plus Equation (7) applied on the two sides of S, plus Equation (8), determine  $\bar{\varphi}$  as a solution of Equation (5).

The conditions satisfied by  $\bar{\varphi}(x, y, z)$  are satisfied also by  $-\bar{\varphi}(x, y, -z)$ . Hence,  $\bar{\varphi}$  is an odd function of  $z$ . This implies that  $\bar{C}_p$  is zero in the wake. In the half space  $z > 0$ ,  $\bar{\varphi}$  is a solution of Equation (5), which satisfies Equation (8) and the boundary conditions

$$\bar{\varphi}_z(x, y, 0+) = \frac{\partial f}{\partial x} + ikf, \quad (x, y) \text{ in } S \quad (10)$$

$$\bar{\varphi}_x(x, y, 0+) + ik\bar{\varphi}(x, y, 0+) = 0, \quad (x, y) \text{ in } W \quad (11)$$

$$\bar{\varphi}(x, y, 0+) = 0, \quad (x, y) \text{ not in } S + W \quad (12)$$

Such a solution may be built up from a doublet distribution over  $S + W$  or a source distribution over the half plane  $z = 0, x > 0$ .

#### 4. BASIC SOURCE AND DOUBLET SOLUTIONS OF THE DIFFERENTIAL EQUATION (See References 7, 8, and 9.)

The solution of Equation (5) which represents a point source at the origin is

$$\bar{\varphi}_0(x, y, z) = \begin{cases} 0, & x \leq 0 \\ -\frac{1}{2\pi} \frac{1}{x} e^{-\frac{1}{2} ik \left( x + \frac{y^2 + z^2}{x} \right)}, & x > 0 \end{cases} \quad (13)$$

(See Reference 9.) The potential of a point doublet oriented parallel to the  $z$ -axis is obtained by differentiation, as

$$\bar{\varphi}_1(x, y, z) = \frac{\partial \bar{\varphi}_0}{\partial z} = \begin{cases} 0, & x \leq 0 \\ \frac{ikz}{2\pi} \frac{1}{x} e^{-\frac{1}{2} ik \left( x + \frac{y^2 + z^2}{x} \right)}, & x > 0 \end{cases} \quad (14)$$

It is easily verified that these functions satisfy Equation (5) for  $x \neq 0$ . They are poorly behaved at  $x = 0$  for real values of  $k$ .

To improve the behavior of  $\bar{\varphi}_0$  and  $\bar{\varphi}_1$  at  $x = 0$ , assume that  $k$  has a small negative imaginary part. This causes  $\bar{\varphi}_0$  and  $\bar{\varphi}_1$  to approach zero exponentially as  $x \rightarrow 0+$ , except at the origin. All partial derivatives of all orders have the same property. Thus,  $\bar{\varphi}_0$  and  $\bar{\varphi}_1$  are solutions of Equation (5) everywhere except at  $(0, 0, 0)$ . In the final formulas to be obtained, the imaginary part of  $k$  can be put equal to zero.

Solutions of Equation (5) for  $z > 0$  which satisfy Equation (8) are given for a distribution of sources as

$$\bar{\varphi}_s(x, y, z) = \iint_{\xi > 0} \rho(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, z) d\xi d\eta \quad (15)$$

and for a distribution of doublets as

$$\bar{\varphi}_d(x, y, z) = \iint_{\xi > 0} \rho(\xi, \eta) \bar{\varphi}_1(x-\xi, y-\eta, z) d\xi d\eta \quad (16)$$

where, to be completely general,  $\rho(\xi, \eta)$  may be any function such that the integrals exist. From the form of  $\bar{\varphi}_0$  and  $\bar{\varphi}_1$ , the region of integration may be restricted to the plane strip  $0 < \xi < x$ . It is shown in Appendix I that these functions satisfy the following boundary conditions for  $z = 0, x > 0$ :

$$\bar{\varphi}_{sz}(x, y, 0+) = \rho(x, y) \quad (17)$$

$$\bar{\varphi}_d(x, y, 0+) = \rho(x, y) \quad (18)$$

(in fact, if the same function  $\rho$  is used in both integrals,  $\bar{\varphi}_d = \partial\bar{\varphi}_s/\partial z$ ).

## 5. THE DETERMINATION OF $\bar{\varphi}$ BY A SOURCE DISTRIBUTION

One method of attack on the problem of finding  $\bar{\varphi}$  is to set  $\bar{\varphi} = \bar{\varphi}_s$ . Then, in terms of the upwash

$$w(x, y) = \bar{\varphi}_z(x, y, 0+) \quad (19)$$

we have from Equations (17) and (15)

$$\bar{\varphi}(x, y, z) = \iint_{\xi > 0} w(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, z) d\xi d\eta \quad (20)$$

for  $z \geq 0$ .

The values of  $w$  on  $S$  are known by Equation (10). Elsewhere,  $w$  is unknown, and it must be chosen so that the boundary conditions (11) and (12) are satisfied. We may take the limit as  $z \rightarrow 0+$  in Equation (20) by taking the limit under the integral sign:

$$\bar{\varphi}(x, y, 0+) = \iint_{\xi > 0} w(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, 0) d\xi d\eta \quad (21)$$

From Equations (11) and (12) are obtained the system of integral equations

$$\iint_{\xi > 0} w(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, 0) d\xi d\eta = 0, \quad (x, y) \text{ not in } S + W \quad (22)$$

$$\left(\frac{\partial}{\partial x} + ik\right) \iint_{\xi > 0} w(\xi, \eta) \bar{\varphi}_0(x-\xi, y-\eta, 0) d\xi d\eta = 0, \quad (x, y) \text{ in } W \quad (23)$$

Solution of Equations (22) and (23), followed by evaluation of  $\bar{\varphi}$  according to Equation (21), would yield the values of  $\bar{\varphi}$  on  $S$ , from which pressures and forces can be computed.

A box method based on a source distribution, described briefly in Reference 9, has been used by Weatherill at the Boeing Company. Some of his preliminary results are given in Reference 9.

## 6. THE DETERMINATION OF $\bar{\varphi}$ BY A DOUBLET DISTRIBUTION

If we set  $\bar{\varphi} = \bar{\varphi}_d$ , then by Equations (18), (16), and (12)

$$\bar{\varphi}(x, y, z) = \iint_{S+W} \bar{\varphi}(\xi, \eta, 0+) \bar{\varphi}_1(x-\xi, y-\eta, z) d\xi d\eta \quad (24)$$

In terms of

$$\psi(x, y) = \lim_{z \rightarrow 0} \frac{1}{z} \bar{\varphi}_1(x, y, z) = \begin{cases} 0, & x \leq 0 \\ \frac{ik}{2\pi} \frac{1}{x} e^{-\frac{1}{2} ik \left(x + \frac{y^2}{x}\right)}, & x > 0 \end{cases} \quad (25)$$

the normal derivative of  $\bar{\varphi}$  at  $z = 0$  is given by a singular integral:

$$w(x, y) = \iint_{S+W} \bar{\varphi}(\xi, \eta, 0+) \psi(x-\xi, y-\eta) d\xi d\eta \quad (26)$$

The values of  $\bar{\varphi}(\xi, \eta, 0+)$  must be determined then from

$$\iint_{S+W} \bar{\varphi}(\xi, \eta, 0+) \psi(x-\xi, y-\eta) d\xi d\eta = w(x, y), \quad (x, y) \text{ in } S \quad (27)$$

$$\left( \frac{\partial}{\partial x} + ik \right) \bar{\varphi}(x, y, 0+) = 0, \quad (x, y) \text{ in } W \quad (28)$$

## 7. A COMPARISON OF THE METHODS

Except for the singularity of the integral in Equation (27), all points of difference are in favor of solving the problem by doublets. There are these points:

- a. The region of integration in the source method extends theoretically to  $\pm\infty$  in  $\eta$ ; even practically, the region must be extended an extreme distance. In the doublet method, the region is restricted to  $S + W$ . This distinction is not so great for supersonic flows. There, the region of influence of the wing is swept back along Mach lines, and the set of points in this region that influences the wing is bounded (see Reference 10).
- b. After the unknown function under the integral sign is known, the source method requires an extra step — the evaluation of  $\bar{\varphi}$  on the wing from Equation (21).
- c. If values in the wake must be considered, the condition in the wake for the source method, Equation (23), is more complicated than the corresponding condition, Equation (28), for the doublet method.

The doublet method was used because of point a.

## 8. THE ADVANTAGE OF A STRAIGHT TRAILING EDGE

Suppose the wing has a straight trailing edge perpendicular to the direction of flow ( $x = \text{constant}$  along the edge); then the wing is not influenced by the wake. This is reflected in the equations by the fact that the integrands are zero when  $\xi > x$ . Hence, in either method, for the determination of  $\bar{\varphi}$  on the wing, the condition in  $W$  need not be used.

## 9. THE DOUBLET BOX METHOD

Consider a flow at Mach 1 past an oscillating wing with its nose at the origin, lying approximately in the  $xy$ -plane, with  $x = 1$  along the trailing edge. The value of the unsteady potential  $\bar{\varphi}$  on the wing may be found by solution of Equation (27), which may be written as

$$\iint_S \bar{\varphi}(\xi, \eta, 0+) \psi(x-\xi, y-\eta) d\xi d\eta = w(x, y), \quad (x, y) \text{ in } S \quad (29)$$

To get an approximate solution of this equation, let the  $xy$ -plane be covered with a grid of square boxes with sides of length  $d$ , so that box edges lie along the coordinate axes (see Figure 2). Let the region  $B$  be composed of all boxes whose centers lie in  $S$ ;  $B$  is an approximation to  $S$  by boxes. Let  $i, j$  be box indexes in the  $x$ - and  $y$ -directions. Approximate  $\bar{\varphi}$  by a constant value  $\bar{\varphi}_{ij}$  in the  $(i, j)$ -th box  $B_{ij}$ . Impose the condition of Equation (7) at the center  $(x_i, y_j)$  of each box  $B_{ij}$  in  $B$ , with the region of integration replaced by  $B$ . Then Equation (29) gives a system of linear algebraic equations for the  $\bar{\varphi}_{ij}$ 's:

$$\sum_{i', j'} \bar{\varphi}_{i'j'} \iint_{B_{i'j'}} \psi(x_i - \xi, y_j - \eta) d\xi d\eta = w(x_i, y_j) \quad (30)$$

Examination of the integral in Equation (30) shows that it depends on  $i, j, i', j'$  only via  $i-i', |j-j'|$ . The notation

$$A(i-i', |j-j'|) = \iint_{B_{i'j'}} \psi(x_i - \xi, y_j - \eta) d\xi d\eta \quad (31)$$

is introduced. Formulas for the evaluation of this quantity are given in Appendix II.

Segregating the terms with  $i'=i$  on the left, Equation (30) becomes

$$\sum_{j'} A(0, |j-j'|) \bar{\varphi}_{ij'} = w(x_i, y_j) - \sum_{i' < i} \sum_{j'} A(i-i', |j-j'|) \bar{\varphi}_{i'j'} \quad (32)$$

For fixed  $i$  and varying  $j$ , this is a smaller system of equations that may be solved for each consecutive value of  $i$ .

# Contrails

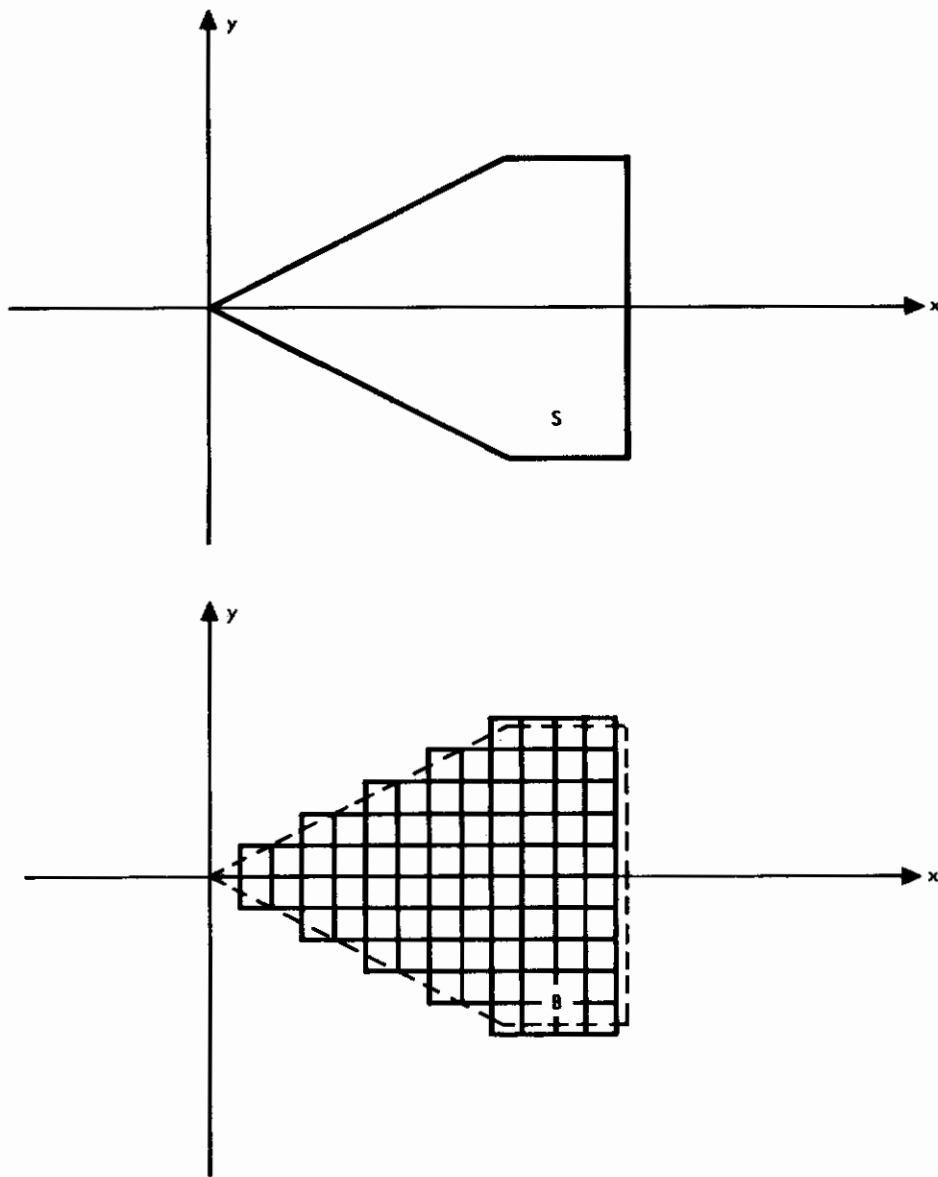


Figure 2. Approximation of the Wing by Region B

Now suppose the wing is symmetric about the x-axis; then only modes of oscillation that are symmetric or antisymmetric in y need be treated. Consider a symmetric mode.  $\bar{\varphi}_{ij}$  will have the same value at corresponding boxes across the x-axis. This may be used to reduce the range of the sums in Equation (32) and the range of j. Let  $j = 1$  in the row of boxes in which  $0 < y < d$ . Then, combining terms for symmetrically placed boxes,

$$\sum_{j' \geq 1} [A(0, |j-j'|) + A(0, j+j'-1)] \bar{\varphi}_{ij'} \quad (33)$$

$$= w(x_i, y_j) - \sum_{i' < i} \sum_{j' \geq 1} [A(i-i', |j-j'|) + A(i-i', j+j'-1)] \bar{\varphi}_{i'j'}, \quad j \geq 1$$

The equations for  $j \leq 0$  are implied by those with  $j \geq 1$ . Thus, the size of the system has been reduced by a factor of 2.

For antisymmetric modes, Equation (33) applies, with the sums of values of A replaced by differences.

## 10. EXTENSIONS OF THE METHOD

The computer program discussed in Section 3 has some restrictions that are not inherent in the box method, such as the requirement of a straight trailing edge. Some possible modifications that extend the applicability of the program will now be described.

To modify the program for modes antisymmetric in y, it is only necessary to change some of the signs in Equation (33), as indicated in the discussion above, and replace even powers of y by odd powers in the formulas used for deflection and potential.

To deal with a more general trailing edge, it is necessary to use the values of  $\bar{\psi}$  in the wake. For fixed y, if  $x = x_T$  at the trailing edge, Equation (28) may be integrated to give

$$\bar{\psi}(x, y, 0+) = e^{-ik(x - x_T)} (x_T, y, 0+)$$

in W. In addition to the set of boxes B on the wing, a corresponding set of boxes  $B_W$  on W must be considered. After finding a value  $\bar{\psi}_{ij}$  in a box of B along the trailing edge, the formula above may be used to find values in the boxes directly downstream. If the ith row of boxes includes boxes of  $B_W$ , to the right side of Equation (33) must be added the contribution of all boxes  $B_{i'j'}$  in  $B_W$  with  $i' < i$ . The computer program must also be modified in several other respects, to take into account the more general wing shape.



# Contrails

A wing that consists of several almost planar sections in different planes, such as a wing with folded tips, may also be handled by the doublet box method. Equation (33) applies, if  $\bar{\psi}_{ij}$  is interpreted as one-half of the discontinuity in  $\bar{\psi}$  between the upper and lower surfaces. The influence coefficients involved are given by a more general formula (not given in this report), allowing for out-of-plane influence of the doublets. Formulas analogous to those of Appendix II may be developed, which are not much more complicated. The main effect of this extension on the computer program would be a greater number of distinct values of the influence coefficients, so that it would not be possible to store them all in an array in core unless the limit on the number of boxes in each direction were considerably reduced.

Rectangular boxes, not necessarily square, may be used. Let the boxes have sides of length  $d_1$  chordwise and  $d_2$  spanwise.

If

$$l_1 = kd_1, \quad l_2 = kd_2^2/d_1$$

the formula for the influence coefficients, Equation (39) in Appendix II, must be replaced by

$$A(n, m) = \frac{ik}{2\pi} \iint_{\substack{|v-m| < 1/2 \\ |u-n| < 1/2 \\ u > 0}} \frac{dudv}{u^2} e^{-1/2 i (l_1 u + l_2 v^2/u)}$$

This may be evaluated by the methods of Appendix II. Except for this difference, the method is essentially the same. The best choice of box shape probably depends on the aspect ratio of the wing.



## 3. DESCRIPTION OF THE COMPUTER PROGRAM

### 1. COORDINATE SYSTEMS

An initial coordinate system  $(\tilde{x}, \tilde{y}, \tilde{z})$  is assumed, with the  $\tilde{x}$ -axis in the direction of the flow. The undisturbed position of the wing is in a region  $S$  in the  $\tilde{x}\tilde{y}$ -plane, with the  $\tilde{x}$ -axis along the center line and the origin at the nose (see Figure 1). This coordinate system is used in the data.

In the program, a dimensionless coordinate system  $(x, y, z)$  is used, based on the root chord length  $b$ :

$$x = \tilde{x}/b$$

$$y = \tilde{y}/b$$

$$z = \tilde{z}/b$$

### 2. WING GEOMETRY

The wing is symmetric, with trailing edge  $\tilde{x} = b$ . To complete its description, the portion of the leading edge on which  $\tilde{y} > 0$  must be specified. This is done by giving the coordinates of the end points of  $NS$  line segments along the edge ( $1 \leq NS \leq 3$ ), beginning at a point at which  $\tilde{y} = 0$ :  $(0, \tilde{y}_0), (\tilde{x}_1, \tilde{y}_1), \dots, (\tilde{x}_{NS}, \tilde{y}_{NS})$ . The edge of  $S$  includes the polygonal line through these points. If  $\tilde{y}_0 > 0$ , it also includes the line from the origin to  $(0, \tilde{y}_0)$ . If  $\tilde{x}_{NS} < b$ , it includes the line from  $(\tilde{x}_{NS}, \tilde{y}_{NS})$  to  $(b, \tilde{y}_{NS})$ . (See Figure 3.)

Leading edges of fairly general shape may be approximated by such polygonal lines.

### 3. THE DEFLECTION DATA

A mode is specified by the vertical deflection function  $f(\tilde{x}, \tilde{y})$  in terms of which the equation for the instantaneous position of the planform is

$$\tilde{z} = \text{Re} [\delta \cdot e^{i\omega t} f(\tilde{x}, \tilde{y})]$$

where  $\delta$  is a constant.

In the program,  $f$  is assumed to be a polynomial in  $\tilde{x}$  and  $\tilde{y}^2$ . The data may give either the coefficients of this polynomial, or values of  $f$  at a

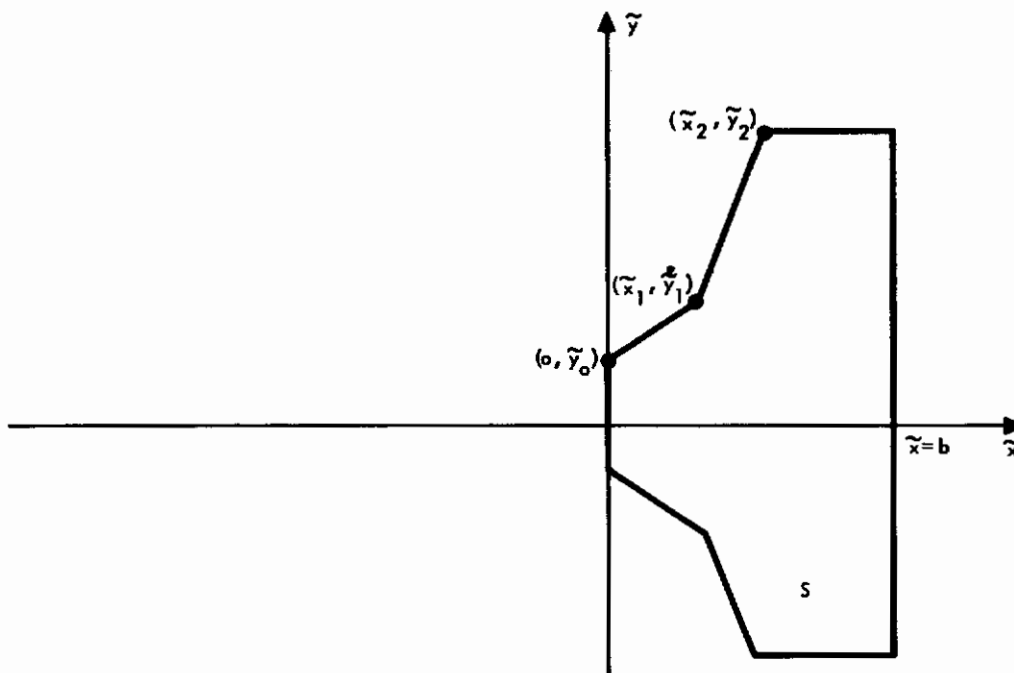


Figure 3. Wing Geometry (NS = 2)

set of points on the wing. In the latter case, a polynomial is fitted to the given values by a least square error technique.

#### 4. LEAST SQUARE SURFACE FITS

The problem involved here is the approximation of a function of  $x$  and  $y$  by an expression of the form

$$\bar{\varphi}(x, y) = \sum_{n, m} a_{nm} x^n y^{2m} F(x, y)$$

when a set of values of the function is known. This arises in the program in two places. The subroutine DRED fits a representation of the deflection of this type with  $F = 1$  to the given deflection values. In the subroutine BXP, such a fit is made for the potential, with

$$F(x, y) = \left\{ \begin{array}{l} \sqrt{x^2 - x_0(y)^2} \\ \text{or} \\ \sqrt{x - x_0(y)} \end{array} \right\} \cdot \left\{ \begin{array}{l} \sqrt{1 - y^2/y_{\max}^2} \\ \text{or} \\ 1 \end{array} \right\}$$

# Contrails

( $x = x_0(y)$  is the equation of the leading edge) depending on the wing shape. This factor approximates the proper behavior of  $\bar{\varphi}$  at the edges. The factor  $\sqrt{x^2 - x_0(y)^2}$  is used for a pointed nose ( $\tilde{y}_0 = 0$ ), and  $\sqrt{x - x_0(y)}$  for an unswept nose ( $\tilde{y}_0 > 0$ ). The factor  $\sqrt{1 - y^2/y_{\max}^2}$  is included if the planform has a side edge along which  $y = y_{\max}$ .

The factor  $F(x, y)$  is real, so the values of  $\bar{\varphi}$  have real and imaginary parts that involve only the corresponding parts of the  $a_{nm}$ 's. Hence, these real and imaginary parts may be handled separately, reducing the problem from one in complex numbers to one in real numbers.

Let  $\alpha_{nm} = \text{Re}[a_{nm}]$ , and let the real parts of given values of the function at data points be  $\bar{\varphi}'_j$  at  $(x_j, y_j)$ ,  $j = 1, \dots, NP$ . Then for the real parts we wish to have

$$\sum_{n,m} \alpha_{nm} x_j^n y_j^{2m} F(x_j, y_j) \cong \bar{\varphi}'_j, \quad j = 1, \dots, NP$$

The least squares method minimizes

$$Q = \sum_j \left[ \sum_{n,m} \alpha_{nm} x_j^n y_j^{2m} F(x_j, y_j) - \bar{\varphi}'_j \right]^2$$

(See Reference 11, Chapter 16.)

For condensed notation, let  $r$  be a single index over the pairs  $(n, m)$ , let  $\alpha_{nm} = \alpha_r$ , and  $x_j^n y_j^{2m} F(x_j, y_j) = A_{jr}$ . Then

$$Q = \sum_j \left[ \sum_r \alpha_r A_{jr} - \bar{\varphi}'_j \right]^2$$

Let the range of  $r$  be from 1 to  $NC \leq NP$ .

To minimize  $Q$ , we set

$$\frac{\partial Q}{\partial \alpha_r} = 0, \quad r = 1, \dots, NC$$

This leads to the system of equations

$$\sum_{r'} \left( \sum_j A_{jr} A_{jr'} \right) \alpha_{r'} = \sum_j A_{jr} \bar{\varphi}_j, \quad r = 1, \dots, NC \quad (34)$$

Put

$$\left. \begin{aligned} \sum_j A_{jr} A_{jr'} &= B_{rr'} \\ \sum_j A_{jr} \bar{\varphi}_j &= C'_r \end{aligned} \right\} \quad (35)$$

Then Equation (34) reduces to

$$\sum_{r'} B_{rr'} \alpha_{r'} = C'_r, \quad r = 1, \dots, NC \quad (36)$$

The matrices  $(B_{rr'})$  and  $(C'_r)$  must be set up to solve Equation (36). It is not necessary, however, to set up the matrix  $(A_{jr})$ . Only one row of  $(A_{jr})$  is needed at a time. This is fortunate, because the program allows  $(A_{jr})$  to become as large as 2500 x 20. For each value of  $j$ , the  $j$ th row of  $(A_{jr})$  is computed, and from this the  $j$ th terms in the sums in Equation (35) are formed and added in.

In the complex case, there is a corresponding system of equations for the imaginary parts:

$$\sum_{r'} B_{rr'} \beta_{r'} = C''_r$$

The two systems of equations are solved together by the subroutine XSIMEQ, which allows for more than one set of values on the right.

## 5. GENERALIZED FORCES

The generalized force coefficient  $L_{ij}$  is defined (Reference 6) by

$$L_{ij} = \frac{1}{1/2 \rho U_\infty^2 S} \iint_S \Delta p_i(x, y) f_j(x, y) dx dy$$

# Contrails

where  $\Delta p_i$  is the lifting pressure difference in the  $i$ th mode, and  $f_j$  is the deflection function in the  $j$ th mode. In terms of the potential  $\bar{\varphi}(x, y)$  on the upper surface,

$$\Delta p_i = 2\rho U_\infty^2 (\bar{\varphi}_x + i k \bar{\varphi})$$

$$L_{ij} = \frac{4}{S} \iint_S (\bar{\varphi}_x + i k \bar{\varphi}) f_j \, dx \, dy$$

After integration by parts,

$$L_{ij} = \frac{4}{S} \left\{ \int_{x=1} \bar{\varphi} f_j \, dy + \iint_S \bar{\varphi} \left( i k f_j - \frac{\partial f_j}{\partial x} \right) \, dx \, dy \right\} \quad (37)$$

In Equation (37) insert the series

$$\bar{\varphi} = \sum_{n, m} a_{nm} x^n y^{2m} F(x, y)$$

$$f_j = \sum_{n', m'} d_{n'm'} x^{n'} y^{2m'}$$

The result is

$$\begin{aligned} L_{ij} = & \frac{8}{S} \sum_{n', m'} d_{n'm'} \sum_{n, m} a_{nm} \left[ \frac{1}{2} \int_{x=1} y^{2m+2m'} F(1, y) \, dy \right. \\ & + i k \cdot \frac{1}{2} \iint_S x^{n+n'} y^{2m+2m'} F(x, y) \, dx \, dy \\ & \left. - n' \cdot \frac{1}{2} \iint_S x^{n+n'-1} y^{2m+2m'} F(x, y) \, dx \, dy \right] \end{aligned}$$

The integrals in this expression depend only on the wing shape. They are computed by the subroutine FØRCI before the work on the individual modes begins. During the work on the  $i$ th mode, the sum over  $n$  and  $m$  is

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performed in the last part of the subroutine BØXP, for each set of values of  $n'$  and  $m'$ . The sum over  $n'$  and  $m'$  and multiplication by  $8/S$  is performed in the last part of the main program.

## 6. THE USE OF GAUSSIAN QUADRATURE IN THE EVALUATION OF GENERALIZED FORCES

Gaussian quadrature is an approximation of the form

$$\int_a^b f(u) du \cong \sum_{j=1}^N h_j f(u_j)$$

exact for polynomials of degree  $\leq 2N - 1$ . (See Reference 11, Chapter 7.) This formula is used with  $(a, b) = (0, 1)$ ,  $N = 6$ . The values of the  $h_j$ 's and  $u_j$ 's for this case were obtained from values listed in Reference 11. They are given as  $H(1), \dots, H(6)$ ,  $U(1), \dots, U(6)$  in the subroutine SECT.

Subroutine FØRCI finds the values of

$$AXY(I, J) = \frac{1}{2} \iint_S x^{I-1} y^{2J-2} F(x, y) dx dy$$

and

$$AY(J) = \frac{1}{2} \int_{x=1} y^{2J-2} F(1, y) dy$$

for  $I, J = 1, \dots, 9$ . To do this, the contributions to the integrals from each section of wing behind a straight piece of leading edge are calculated separately in SECT.

The form of  $F(x, y)$  is

$$F(x, y) = \left\{ \begin{array}{l} \sqrt{x - x_o(y)} \\ \text{or} \\ \sqrt{x^2 - x_o(y)^2} \end{array} \right\} \cdot \left\{ \begin{array}{l} \sqrt{1 - y^2/y_{\max}^2} \\ \text{or} \\ 1 \end{array} \right\}$$

depending on the wing shape. We have integrals that behave like square roots at the leading edge. The integrals over one wing section are of the form

$$BXY = \int_{y_-}^{y_+} dy \int_{x_o(y)}^1 dx x^{I-1} y^{2J-2} F(x, y)$$

and

$$BY = \int_{y_-}^{y_+} dy y^{2J-2} F(1, y)$$

In BXY, the chordwise integral is evaluated first at each value of  $y$  at which it will be needed. The new variable

$$u = \sqrt{x - x_o(y)} / \sqrt{1 - x_o(y)} \quad (38)$$

is introduced. Then

$$\int_{x_o(y)}^1 dx x^{I-1} y^{2J-2} F(x, y) = \int_0^1 du \cdot 2 \left[ 1 - x_o(y) \right] x^{I-1} y^{2J-2} F(x, y)$$

The integrand, as a function of  $u$ , is well-behaved at the leading edge. It is approximated by

$$\sigma(y) = \sum_{i=1}^6 h_i \cdot 2 \left[ 1 - x_o(y) \right] x_i^{I-1} y^{2J} F(x_i, y)$$

where  $x_i$  is computed from the value of  $u_i$  according to Equation (38).

In the  $y$ -integration in BXY and BY, the integrand approaches zero as  $y \rightarrow y_{\max}$  like  $\sqrt{1 - y/y_{\max}}$  or  $(1 - y/y_{\max})^{3/2}$ . Accordingly, the change of variable

$$y = \begin{cases} y_+ - (y_+ - y_-)v, & y_+ < y_{\max} \\ y_+ - (y_+ - y_-)v^2, & y_+ = y_{\max} \end{cases}$$

is used, which makes the interval of integration  $0 < v < 1$  and removes the square root behavior in the last section of the wing. This leads to the formulas

$$\begin{aligned}
 BXY &= (y_+ - y_-) \sum_{j=1}^6 h_j \sigma(y_j) \cdot \begin{cases} 1, & y_+ < y_{\max} \\ 2u_j, & y_+ = y_{\max} \end{cases} \\
 BY &= (y_+ - y_-) \sum_{j=1}^6 h_j y_j^{2J-2} F(1, y_j) \cdot \begin{cases} 1, & y_+ < y_{\max} \\ 2u_j, & y_+ = y_{\max} \end{cases}
 \end{aligned}$$

## 7. LEADING EDGE CORRECTION

The value of potential found for each box from Equation (33) is taken to be the value of  $\bar{\varphi}$  at the box center. Thus, the values obtained are in error only by virtue of the error introduced in the values of upwash when the actual distribution of potential in a box is replaced by this constant value. This error is especially important in the first row of boxes, for a wing with an unswept leading edge. The major effect is on the upwash values in that row.

To estimate this error, consider the two-dimensional case, in which  $\bar{\varphi}$  is independent of  $y$ . In Equation (26), the expression for upwash due to a doublet distribution, integrate by parts over  $\xi$ , then integrate over  $\eta$ . The result is

$$\begin{aligned}
 \bar{w}(x, y) &= \frac{ik}{2\pi} \iint_{0 < \xi < x} \frac{d\xi d\eta}{(x - \xi)^2} \bar{\varphi}(\xi, \eta, 0+) e^{-\frac{1}{2} ik \left( x - \xi + \frac{(y - \eta)^2}{x - \xi} \right)} \\
 &= \frac{1}{\pi} \iint_{0 < \xi < x} \frac{d\xi d\eta}{(y - \eta)^2} \left( \bar{\varphi}_\xi + \frac{1}{2} ik \bar{\varphi} \right) e^{-\frac{1}{2} ik \left( x - \xi + \frac{(y - \eta)^2}{x - \xi} \right)} \\
 &= -\sqrt{\frac{2ik}{\pi}} \int_0^x \frac{d\xi}{\sqrt{x - \xi}} e^{-\frac{1}{2} ik(x - \xi)} \left( \bar{\varphi}_\xi + \frac{1}{2} ik \bar{\varphi} \right)
 \end{aligned}$$

For  $x = \frac{1}{2} d$ , if  $kd$  is small,



$$\bar{w}\left(\frac{1}{2}d, y\right) \cong -\sqrt{\frac{2ik}{\pi}} \int_0^{\frac{1}{2}d} \frac{d\xi}{\sqrt{\frac{1}{2}d - \xi}} \bar{\varphi}_{\xi}(\xi, \eta, 0+)$$

The correct leading edge behavior is possessed by the expression  $\bar{\varphi} = C\sqrt{\xi}$ . We have

$$\bar{w}\left(\frac{1}{2}d, y\right) \Big|_{\bar{\varphi} = C\sqrt{\xi}} = -\sqrt{\frac{2ik}{\pi}} \frac{C}{2} \int_0^{\frac{1}{2}d} \frac{d\xi}{\sqrt{\xi\left(\frac{1}{2}d - \xi\right)}} = -\sqrt{\frac{2ik}{\pi}} C \frac{\pi}{2}.$$

If  $\bar{\varphi}$  is constant on  $0 < \xi < d$ , and has the value  $C\sqrt{1/2}d$ , then  $\bar{\varphi}_{\xi}$ , in the above integral, can be expressed in terms of a delta function:

$$\bar{\varphi}_{\xi} = C\sqrt{\frac{1}{2}d} \delta(\xi)$$

Accordingly,

$$\bar{w}\left(\frac{1}{2}d, y\right) \Big|_{\bar{\varphi} = C\sqrt{\frac{1}{2}d}} = -\sqrt{\frac{2ik}{\pi}} C.$$

Note that the latter value is smaller than the value of  $\bar{w}$  evaluated for  $\bar{\varphi} = C\sqrt{\xi}$  by the factor  $2/\pi$ . This implies that the values of potential found for the first row of boxes will be more accurate if the upwashes in that row are multiplied by  $2/\pi$ .

## 8. THE FORM OF OUTPUT

The viewpoint is taken that calculation of the generalized forces is the basic purpose of the program. They are always printed out. There are other outputs that will be printed if the appropriate data signal is given. Each of the following is printed if the data item specified in parentheses is non-zero:

- a. The coefficients of the deflection polynomial, if it has been computed as a fit to given values of deflection, DA(87)
- b. The upwash array, DA(88)
- c. The potential array, DA(89)
- d. The coefficients of the potential series, DA(90)

- e. Values of pressure and potential at the box centers, computed from the series, DA(91).

## 9. THE DATA SUBROUTINE DATRD

This subroutine reads all data items into the array DA. Punched cards used for data are considered to contain six fields of length 12 as indicated in the sample data sheets. The first field contains information for DATRD. Ending in column 12 is an integer giving the location in the data array for the entry in the second field. The following fields go into consecutive locations, if the data are numeric. Floating point numbers should be written with decimal points, and fixed point numbers adjusted to the right end of the field.

The word ALPHA in columns 2 through 6 indicates that the data on the card are alphanumeric. These are stored in DA in a different way, taking up ten locations per card. The data may be printed later, just as they appeared on the card.

On a numeric card, if a field is blank the corresponding location in DA is unchanged. This is not true for an ALPHA card.

A minus sign in column 1 indicates the last card to be read at the time. DATRD reads cards until this minus sign is encountered, then returns to the main program.

## 10. A NOTE ON THE USE OF TAPES

In writing of this program, the following tape numbers have been used: output tape, number 6; input tape, number 5; and tape simulated by an internal file, number 99.

The tape numbers 5 and 99 appear in the subroutine DATRD. Elsewhere only the output tape number is used. It occurs in the main program, and in the subroutines SHAPE, DRED, BØXP and BØXPØ.

## 11. USE OF THE PROGRAM FOR FIXED WING AND MODES AT VARIOUS FREQUENCIES

If a non-zero quantity is entered in the appropriate location in the data array, (DA26), it indicates that a wing shape and set of modes to be used are the same already used for another frequency. Then quantities that depend only on wing shape and deflection data will not be computed, but will be taken from the permanent arrays in which they were stored in the previous case. The number of boxes along the root chord, DA(27), may not be changed when this is done.

When this option is exercised, all work for the present frequency will be carried out after reading one set of data, which need only include the frequency and the indicator, DA (26). Titles for the individual modes are not printed.

## 12. DESCRIPTION OF THE DATA ARRAY

All data are entered into the array DA, dimensioned for 700, as described in Paragraph 9. The layout of the array is as follows:

1 - 10	Title
13 - 22	Mode Title
23 :	Frequency (cycles per unit of time), $\nu$
24 :	Root chord length, b
25 :	Speed of sound, a
26 :	Indicator for new frequency (See Paragraph 11)
27 :	Number of boxes along root chord
28 :	Number of modes
29 :	Number of sections of leading edge to be given (See Paragraph 2)
30 - 36	Coordinates of points on the leading edge (See Paragraph 2)
39 :	Indicator to suppress calculation of potential for a mode
46 - 70	Coefficients of the deflection polynomial (See Paragraph 3)
87 - 91	Output indicators (See Paragraph 8)
98 :	Number of points at which deflections are given (See Paragraph 3)
99 :	Number of $\tilde{x}$ values
100 :	Number of $\tilde{y}$ values
101 - 700	Deflection data for a maximum of 150 points

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Note: 23, 24, 25 and 30-36 must be entered in consistent units of length and time.

## 13. OUTLINE OF THE PROGRAM

For the purpose of description, the main program has been divided into 20 parts, as indicated in Figure 4, which shows the flow of the program and the subroutines called.

## 14. SIZE LIMITATIONS OF THE PROGRAM

The program's size limitations are as follows:

- a. Box size – the half wing must be enclosed in a rectangle that contains no more than 50 boxes in each direction. (The use of a large number of boxes is not recommended, because the time required is roughly proportional to the cube of the number of boxes along the root chord. The possibility of 50 boxes in each direction is intended to allow a large range in aspect ratio.)
- b. Number of modes – ten at most.
- c. Number of points at which the deflections are given for one mode – 150 at most.
- d. Terms in the deflection polynomial – this is  $\sum_{nm} d_{nm} x^n y^{2m}$ , where  $0 \leq n \leq 4$ ,  $0 \leq m \leq 4$ .

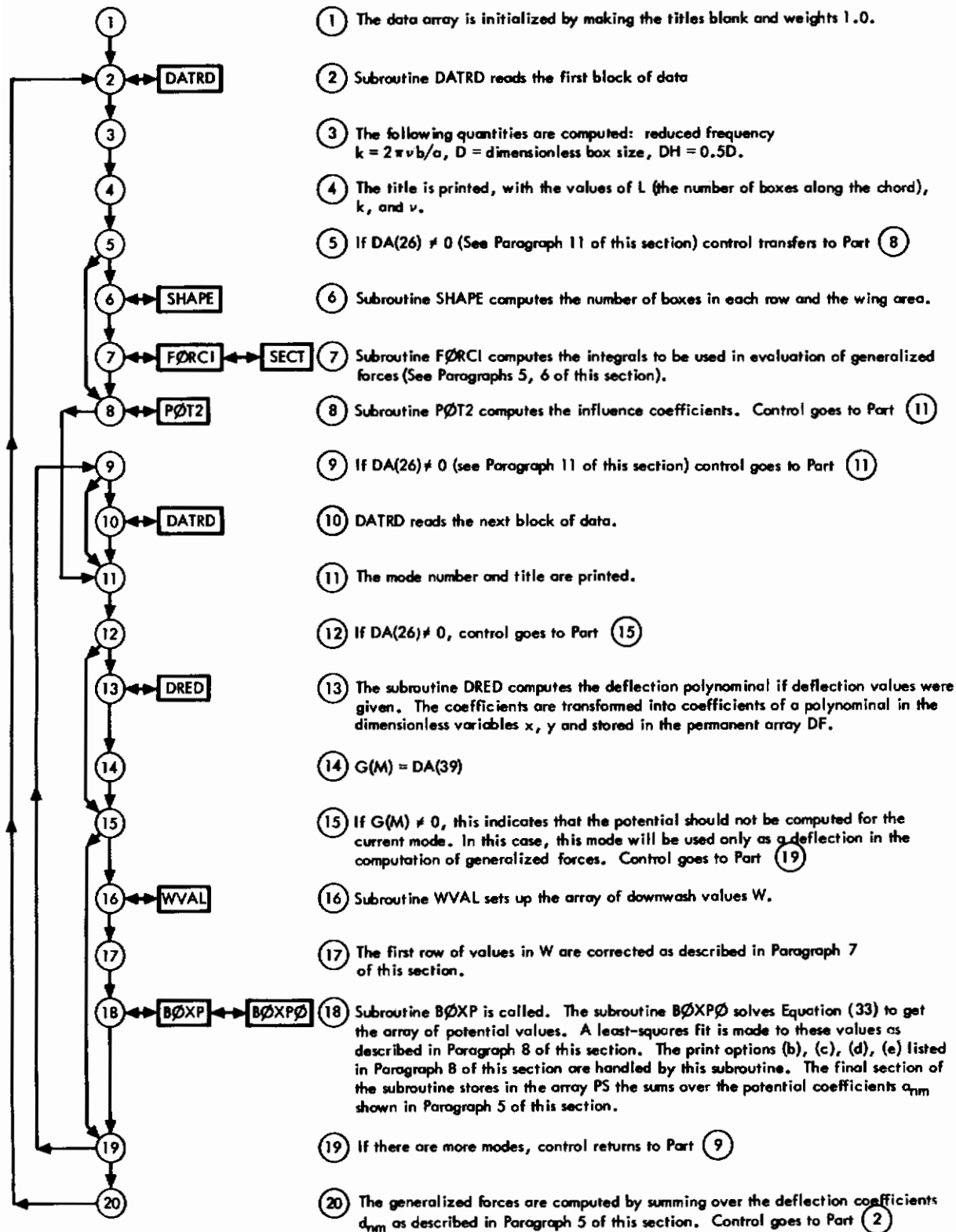


Figure 4. Program Flow Diagram

# *Contrails*

## 4. RESULTS

### 1. THE ASPECT RATIO 1.5 DELTA WING

The computer program was run for the plunging and pitching modes (pitch axis at  $x = 0$ ) at the reduced frequencies  $k = 0.2, 0.5, 0.8, 1.0$ . Forty boxes along the root chord were used, which leads to about 300 boxes on the half-wing.

Theoretical values for comparison were calculated from Davies' formulas (Reference 12). These are analytic expressions of the solution of Equation (5) for the potential and generalized forces for the delta wing in rigid modes of oscillation, expressed as series in  $k$ . Figures 5 through 7 show the values of generalized forces  $L_{11}$  (lift due to plunge),  $L_{21}$  (lift due to pitch), and  $L_{22}$  (moment due to pitch). Note that the vertical scales have been expanded in the portion of interest, especially for  $L_{11}$ . Most of the values agree to within 2 or 3 percent.

The differences indicate the errors introduced by the box method in the solution of Equation (5), as distinguished from the errors inherent in this equation.

Figure 8 gives the chordwise distribution of values of  $\bar{\phi}$  for the plunging mode at  $k = 0.5$ , for  $y = y_{\max}/3 = 0.125$ .

### 2. THE ASPECT RATIO 2.0 RECTANGULAR WING

The plunging and pitching modes were again used at  $k = 0.3, 0.6, 0.9$ . Twenty-five boxes were allowed along the chord, giving 625 boxes on the half-wing. The values of  $L_{21}$  and  $L_{22}$  are shown in Figures 9 and 10, with values from Landahl (Reference 6, page 84) for comparison. Landahl's values were obtained by a method of solution of Equation (5) which applies only to a rectangular wing in modes of oscillation with a deflection independent of  $y$ .

### 3. THE ASPECT RATIO 3.0 RECTANGULAR WING

Finally, for the aspect ratio 3.0 rectangular wing, a comparison is made with experimental pressure values. These values were given in Reference 13 for a 5-percent thickness wing oscillating in an elastic bending mode. At Mach 1, the reduced frequency was 0.24. The chordwise pressure distribution at  $y = y_{\max}/2$  is shown in Figure 11.



# Contrails

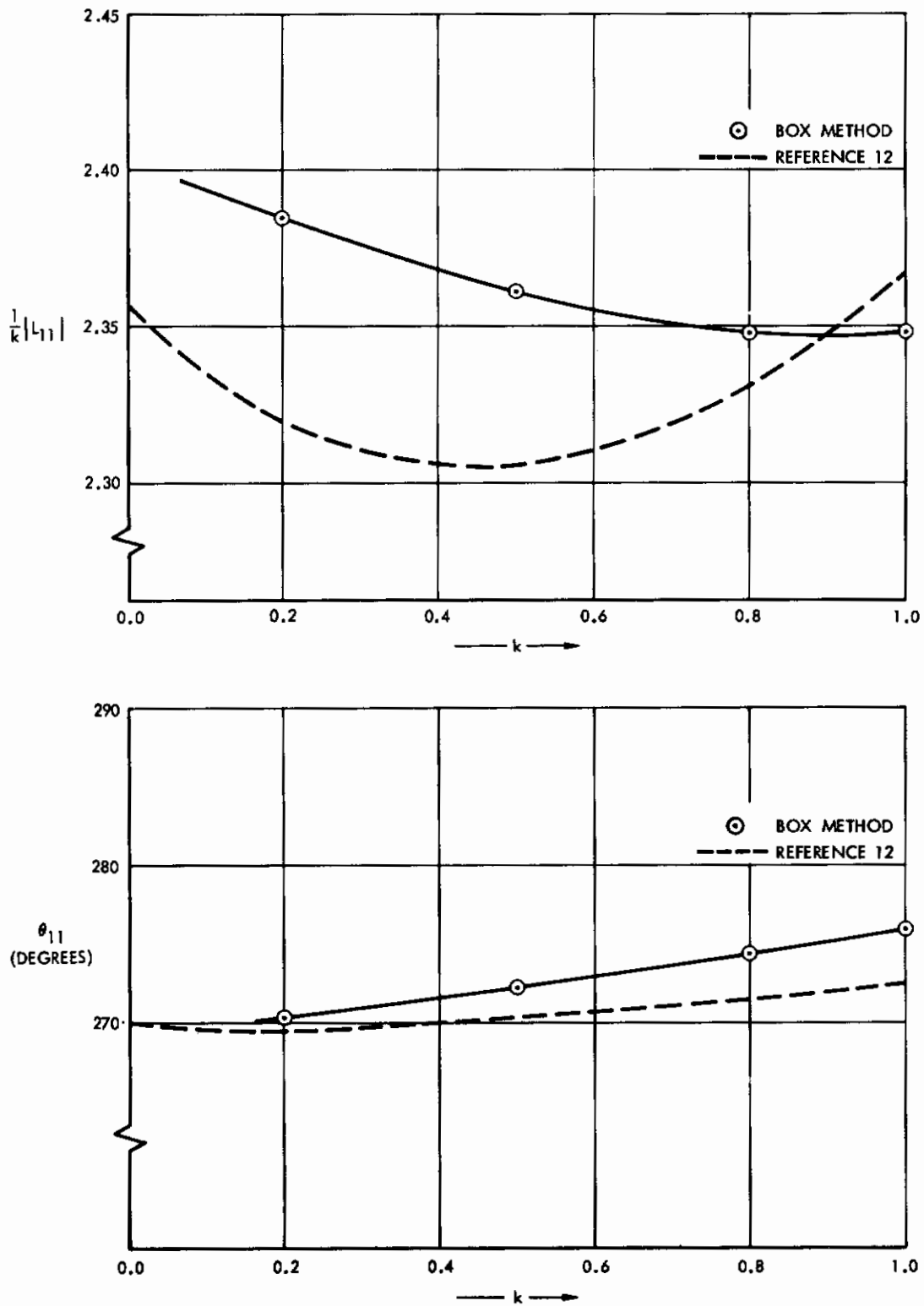


Figure 5. Lift Due to Translation for an Aspect Ratio 1.5 Delta Wing (Compared With Reference 12)



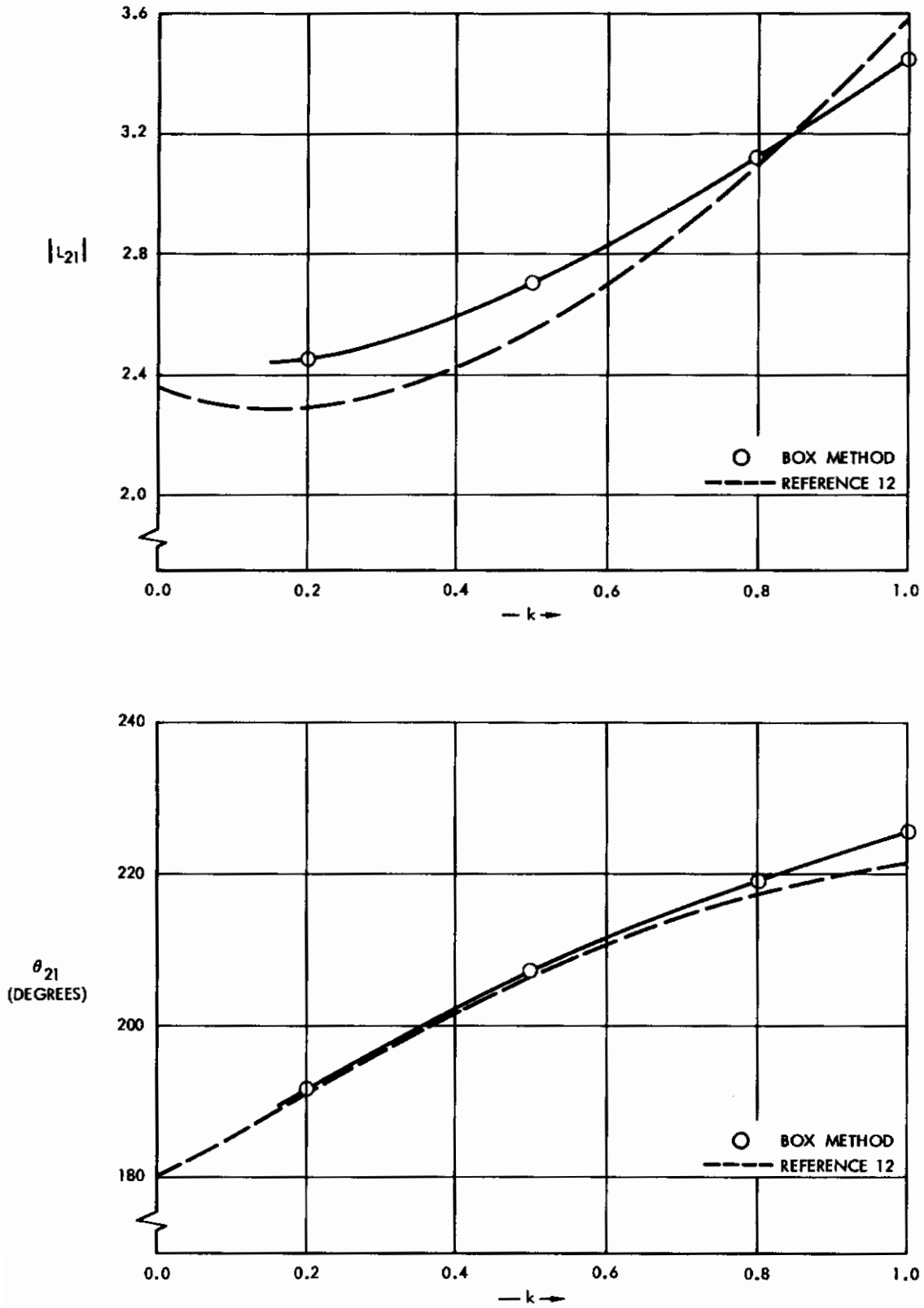


Figure 6. Lift Due to Pitch for an Aspect Ratio 1.5 Delta Wing (Compared With Reference 12)

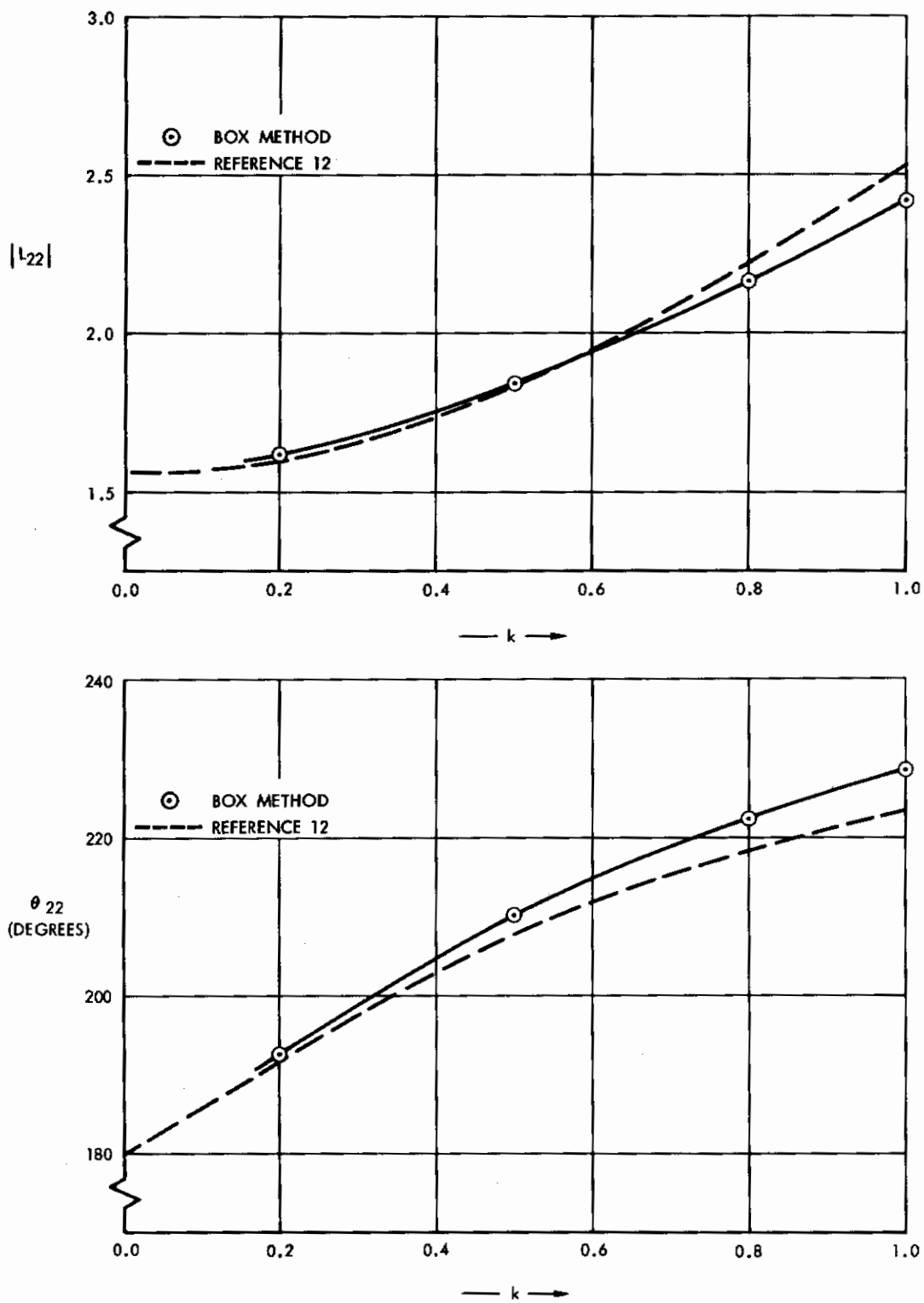


Figure 7. Moment Due to Pitch for an Aspect Ratio 1.5 Delta Wing (Compared With Reference 12)

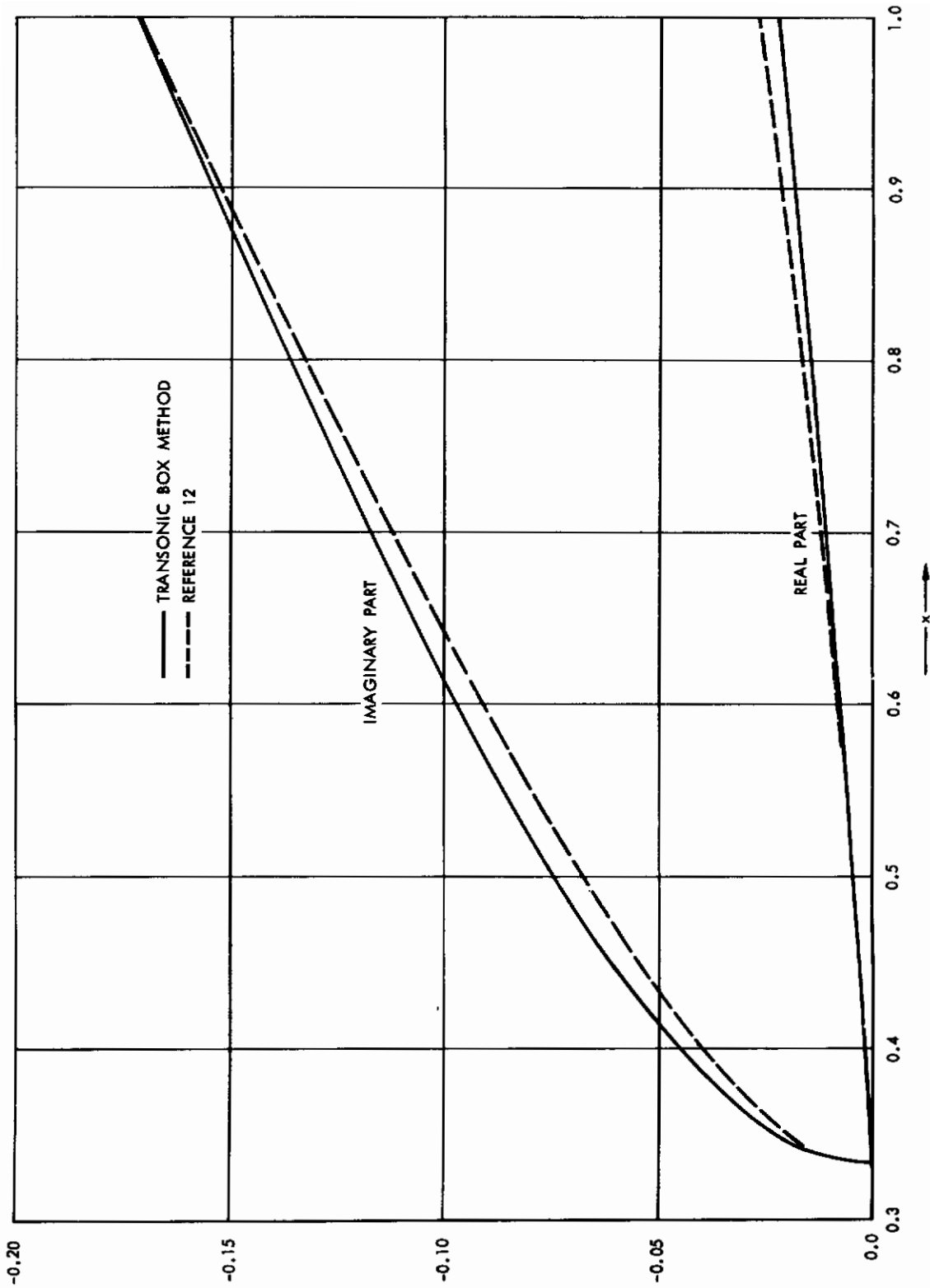


Figure 8. Real and Imaginary Parts of the Unsteady Potential  $\bar{\varphi}$  in the  
 Plunging Mode for an Aspect Ratio 1.5 Delta Wing at  
 $k = 0.5$ ; Chordwise Distribution for  $\gamma = 0.125 = g \text{ Max}/3$   
 (Compared With Reference 12)

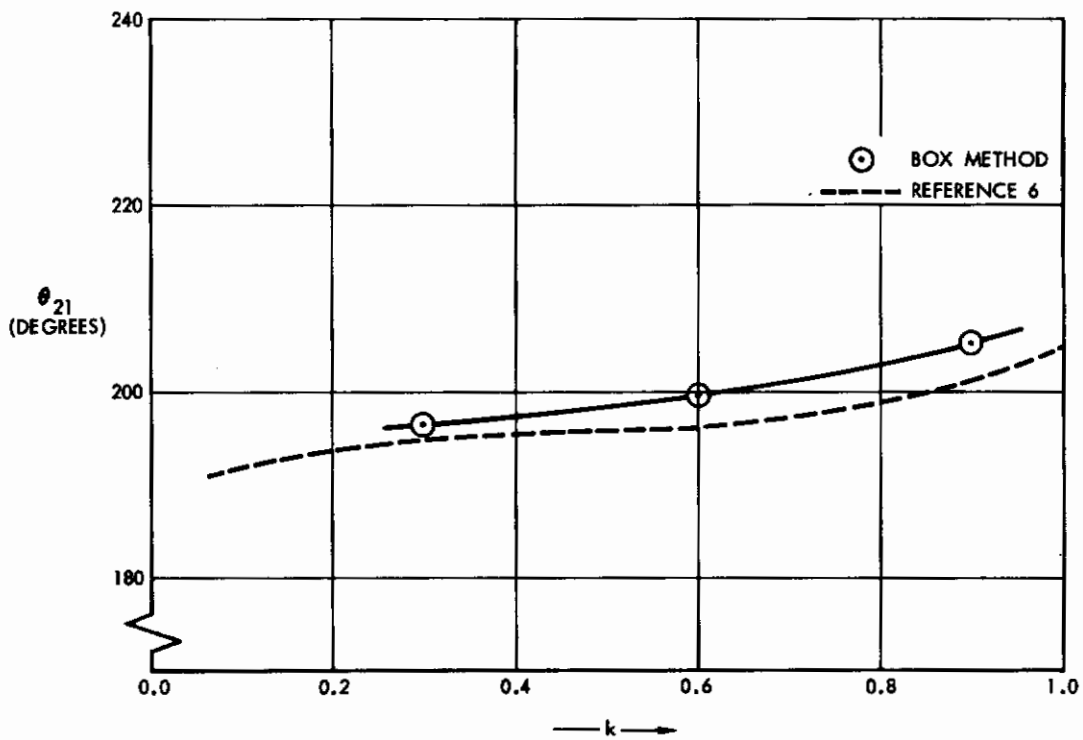
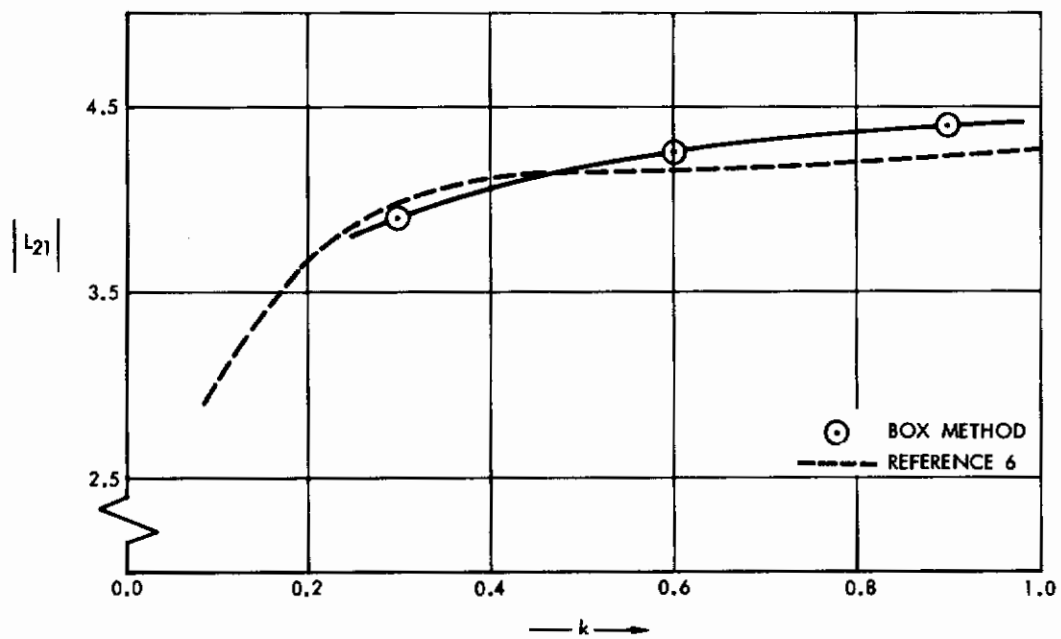


Figure 9. Lift Due to Pitch for an Aspect Ratio 2.0 Rectangular Wing (Compared With Reference 6)

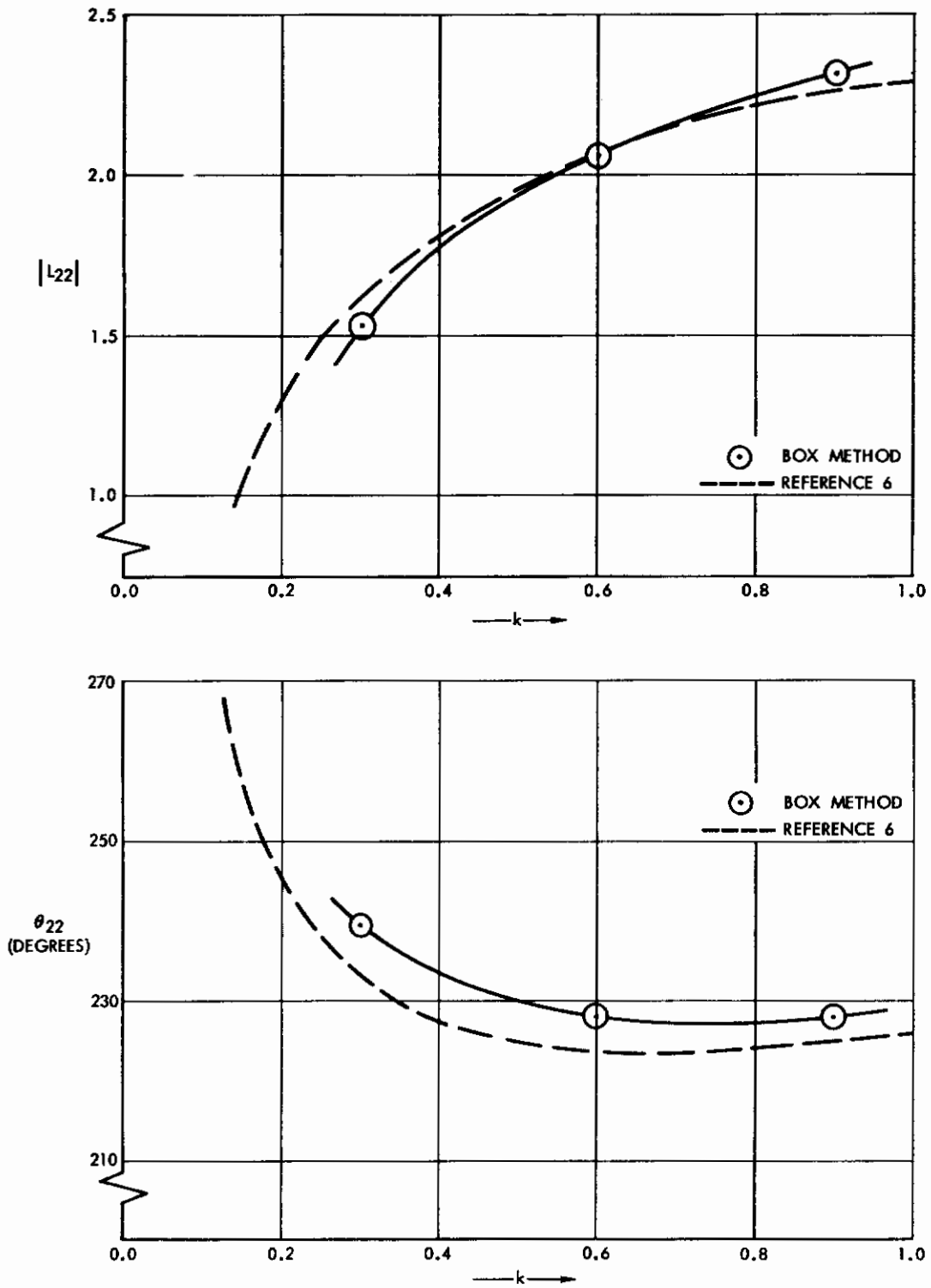


Figure 10. Moment Due to Pitch for an Aspect Ratio 2.0 Rectangular Wing (Compared With Reference 6)

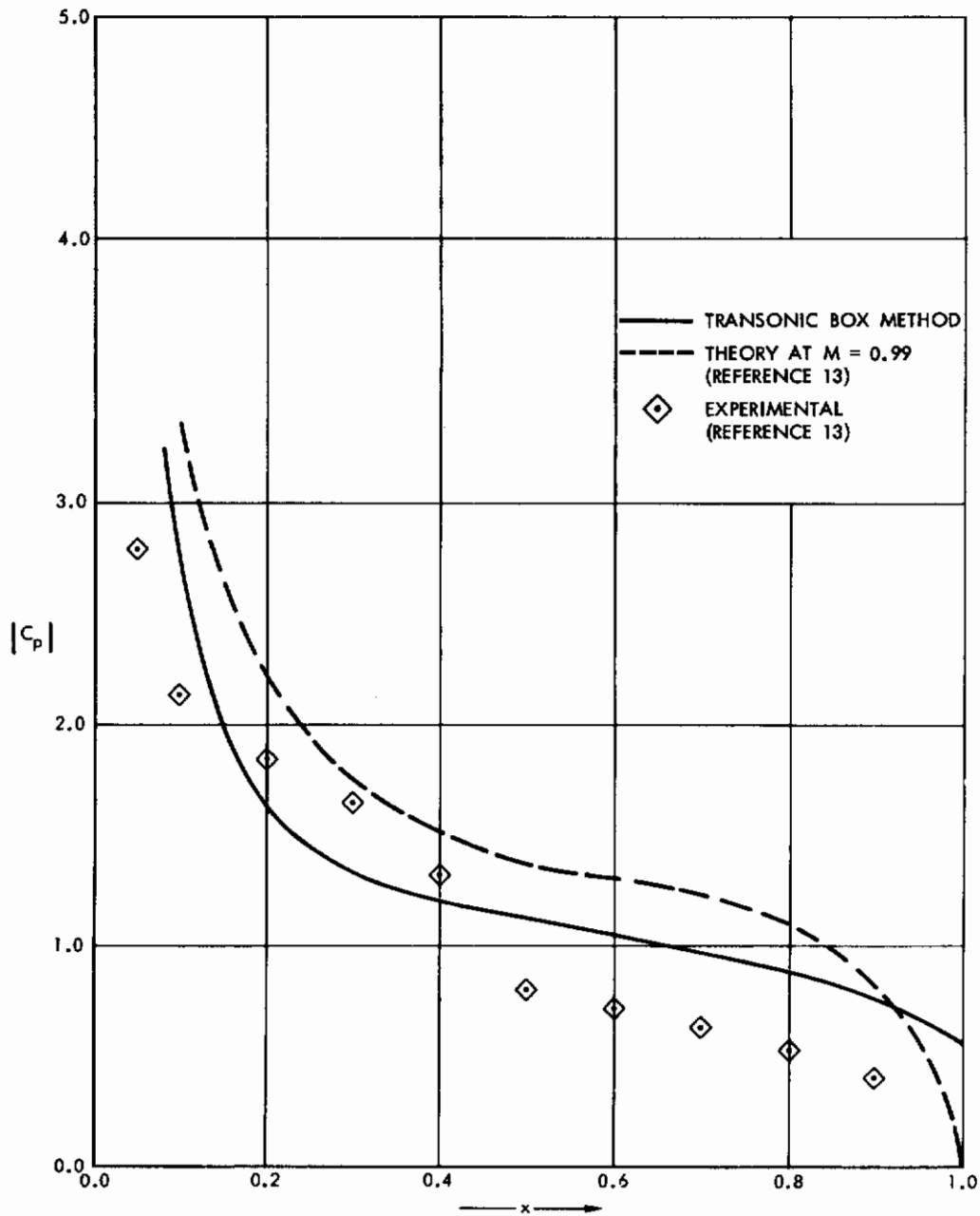


Figure 11. Chordwise Pressure Distribution on an Aspect Ratio 3.0 Rectangular Wing in an Elastic Mode;  $y = y_{\max}/2 = 0.75$  (Compared With Reference 13)

The variation of the experimental values from the computed values is of the type that thickness effects should be expected to cause: the measured pressure is (1) smaller near the leading edge, (2) larger before the point of maximum thickness ( $x = 0.5$ ), and (3) smaller beyond this point. The experimentally determined values of local Mach number along this chord range from 0.84 to 1.35, which indicates that the thickness has a considerable effect on the flow.

The theoretical curve given in Reference 13 was obtained from the subsonic kernel function method, applied at  $M = 0.99$ . This curve is included to show how another theoretical method compares with the experimental values.

#### 4. COMPUTER RUNNING TIME

The results described in this section required about 20 minutes total computer time on the IBM 7094. With nonessential output omitted, this time could have been reduced. All optional output was given, resulting in about 40,000 lines of output.

# *Contrails*



## 5. CONCLUSIONS

A procedure has been developed for predicting unsteady aerodynamic forces and pressures on an oscillating wing by the use of the transonic box method. The results obtained by this method agree quite well with theoretical values from other methods that are applicable only to special planforms. The box method has the advantage of applicability to a general planform. The only other method of this generality at Mach 1 is the sonic limit of the subsonic kernel function method (see Reference 16) that has not been very successful.

The comparison with experimental values in Figure 11 indicates that the most serious limitation of the method is that thickness is neglected. Thickness may be incorporated into a box program by using modified forms for sources and doublets, depending on the local Mach number (see Reference 14). This was not accomplished under the present program. Other possible extensions of the transonic box method, that would not require much change in the existing computer program, are described in Paragraph 10 of Section 2.

# *Contrails*

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## APPENDIX I. PROPERTIES OF SOURCE AND DOUBLET DISTRIBUTIONS

### 1. BOUNDARY BEHAVIOR OF A DOUBLET DISTRIBUTION

We wish to evaluate

$$\bar{\varphi}_d(x, y, 0+) = \lim_{z \rightarrow 0+} \iint_{\xi > 0} d\xi d\eta \rho(\xi, \eta) \bar{\varphi}_1(x-\xi, y-\eta, z).$$

The integrand is zero for  $\xi > x$ . If we define  $\rho(\xi, \eta) = 0$  for  $\xi < 0$ , then

$$\bar{\varphi}_d(x, y, 0+) = \lim_{z \rightarrow 0+} \iint_{\xi < x} d\xi d\eta \rho(\xi, \eta) \bar{\varphi}_1(x-\xi, y-\eta, z).$$

Put

$$\xi = x - z^2 u$$

$$\eta = y - z v$$

Then, using Equation (14), we have

$$\begin{aligned} \bar{\varphi}_d(x, y, 0+) &= \lim_{z \rightarrow 0+} \iint_{u > 0} dudv \rho(x - z^2 u, y - z v) \cdot z^3 \bar{\varphi}_1(z^2 u, z v, z) \\ &= \lim_{z \rightarrow 0+} \iint_{u > 0} dudv \rho(x - z^2 u, y - z v) \cdot \frac{ik}{2\pi} \frac{1}{z} e^{-\frac{1}{2} ik \left( z^2 u + \frac{1+v^2}{u} \right)} \end{aligned}$$

Let  $(x, y)$  be a point at which  $\rho$  is continuous. Then the value of  $\rho$  in the integrand approaches  $\rho(x, y)$  as  $z \rightarrow 0$ . Taking the limit under the integral sign,

$$\bar{\varphi}_d(x, y, 0+) = \rho(x, y) \cdot \frac{ik}{2\pi} \iint_{u > 0} dudv \cdot \frac{1}{z} e^{-\frac{1}{2} ik \frac{1+v^2}{u}}$$

Let  $v = s \sqrt{u}$ . Then

$$\bar{\varphi}_d(x, y, 0+) = \rho(x, y) \cdot \frac{ik}{2\pi} \int_0^{\infty} \frac{du}{u^{3/2}} e^{-\frac{1}{2} ik/u} \int_{-\infty}^{\infty} ds e^{-\frac{1}{2} iks^2}$$

These integrals may be reduced to a standard form by rotating the paths of integration in the complex plane into positions in which the exponents are negative, then making the substitutions

$$u = \frac{ik}{2p}$$

$$s = \sqrt{\frac{2}{ik}} q$$

The result is

$$\begin{aligned} \bar{\varphi}_d(x, y, 0+) &= \rho(x, y) \cdot \frac{1}{\pi} \int_0^{\infty} e^{-p} p^{-\frac{1}{2}} dp \int_{-\infty}^{\infty} e^{-q^2} dq \\ &= \rho(x, y) \end{aligned}$$

since both integrals have the value  $\sqrt{\pi}$  (see Reference 12, formulas 507, 512). Consequently, Equation (18) is valid at any point of continuity of  $\rho(x, y)$ .

## 2. BOUNDARY BEHAVIOR OF A SOURCE DISTRIBUTION

To evaluate  $\bar{\varphi}_{sz}(x, y, 0+)$ , note that

$$\begin{aligned} \frac{\partial \bar{\varphi}_s}{\partial z} &= \iint_{\xi > 0} d\xi d\eta \rho(\xi, \eta) \frac{\partial}{\partial z} \bar{\varphi}_0(x-\xi, y-\eta, z) \\ &= \iint_{\xi > 0} d\xi d\eta \rho(\xi, \eta) \bar{\varphi}_1(x-\xi, y-\eta, z) \\ &= \bar{\varphi}_d \end{aligned}$$

# Contrails

Hence, by the result of the preceding section, if  $(x, y)$  is a point at which  $\rho(x, y)$  is continuous,

$$\bar{\varphi}_{sz}(x, y, 0+) = \bar{\varphi}_d(x, y, 0+) = \rho(x, y)$$

which verifies Equation (17).

# *Contrails*



## APPENDIX II. EXPRESSIONS FOR THE INFLUENCE COEFFICIENTS

Equation (31) may be expressed more conveniently in terms of

$$u = (x_i - \xi)/d$$

$$v = (y_j - \eta)/d$$

$$m = |j - j'|$$

$$n = i - i'$$

$$l = kd$$

(d = box side length). By Equation (14) we have

$$A(n, m) = \frac{ik}{2\pi} \iint \frac{du dv}{u} e^{-\frac{1}{2} i l \left( u + \frac{v^2}{u} \right)} \quad (39)$$

$$\begin{aligned} & |v-m| < \frac{1}{2} \\ & |u-n| < \frac{1}{2} \\ & u \geq 0 \end{aligned}$$

It is assumed that  $l$  is small. If  $l < 0.1$ , the following approximation gives an error of less than 0.1 percent in the value of A:

$$\begin{aligned} e^{-\frac{1}{2} i l u} &= e^{-\frac{1}{2} i l n} e^{-\frac{1}{2} i l (u - n)} \\ &\cong e^{-\frac{1}{2} i l n} \left[ 1 - \frac{1}{2} i l (u - n) - \frac{1}{8} l^2 (u - n)^2 \right] \end{aligned}$$

# Contrails

This reduces Equation (39) to

$$\begin{aligned}
 A(n, m) = \frac{ik}{2\pi} e^{-\frac{1}{2} i\ell n} & \left\{ \left( 1 + \frac{1}{2} i\ell n - \frac{1}{8} \ell^2 n^2 \right) \iint \frac{du dv}{u^2} e^{-\frac{1}{2} i\ell v^2/u} \right. \\
 & + \left( -\frac{1}{2} i\ell + \frac{1}{4} \ell^2 n \right) \iint \frac{du dv}{u} e^{\frac{1}{2} i\ell v^2/u} \\
 & \left. - \frac{1}{8} \ell^2 \iint du dv e^{-\frac{1}{2} i\ell v^2/u} \right\} \quad (40)
 \end{aligned}$$

where the limits of integration are the same as in Equation (39).

The following formula expresses these double integrals in terms of single integrals:

$$\begin{aligned}
 \int_{u_1}^{u_2} du \int_{v_1}^{v_2} dv \frac{1}{u^p} e^{-\frac{1}{2} i\ell v^2/u} &= \frac{1}{3-2p} \int_{u_1}^{u_2} \frac{v du}{u^p} e^{-\frac{1}{2} i\ell v^2/u} \Bigg|_{v=v_1}^{v_2} \\
 &+ \frac{2}{3-2p} \int_{v_1}^{v_2} \frac{dv}{u^{p-1}} e^{-\frac{1}{2} i\ell v^2/u} \Bigg|_{u=u_1}^{u_2}
 \end{aligned}$$

Equation (40) becomes

$$A(n, m) = \frac{ik}{2\pi} e^{-\frac{1}{2} i\ell n} \cdot \begin{cases} A_n(n + \frac{1}{2}, m) - A_n(n - \frac{1}{2}, m), & n > 0 \\ A_0(\frac{1}{2}, m), & n = 0 \end{cases} \quad (41)$$

# *Contrails*

where

$$\begin{aligned}
 A_n(u, m) &= v \int_0^u du e^{-\frac{1}{2} i \ell v^2 / u} \left[ -\frac{1}{u^2} \left( 1 + \frac{1}{2} i \ell n - \frac{1}{8} \ell^2 n^2 \right) \right. \\
 &\quad \left. + \frac{1}{u} \left( -\frac{1}{2} i \ell + \frac{1}{4} \ell^2 n \right) - \frac{1}{24} \ell^2 \right] \Big|_{v = m - \frac{1}{2}}^{m + \frac{1}{2}} \\
 &\quad + \left[ -\frac{2}{u} \left( 1 + \frac{1}{2} i \ell n - \frac{1}{8} \ell^2 n^2 \right) + 2 \left( -\frac{1}{2} i \ell + \frac{1}{4} \ell^2 n \right) \right. \\
 &\quad \left. - \frac{1}{12} \ell^2 u \right] \int_{m - \frac{1}{2}}^{m + \frac{1}{2}} dv e^{-\frac{1}{2} i \ell v^2 / u} \\
 &= B_n(u, m) + C_n(u, m)
 \end{aligned} \tag{42}$$

$B_n$  and  $C_n$  denote the contributions of the terms containing the  $u$  - integrals and  $v$  - integrals, respectively.

$B_n(u, m)$  may be expressed in terms of the sine and cosine integrals

$$S(x) = \int_1^{\infty} \frac{\sin xt}{t} dt$$

$$C(x) = \int_1^{\infty} \frac{\cos xt}{t} dt$$

( $C(x)$  and  $S(x)$  are evaluated by the subroutine CIN.) The resulting formula for  $B_n$  is

# Contrails

$$\begin{aligned}
 B_n(u, m) = & \left\{ \left[ \frac{2i}{\ell v} \left( 1 + \frac{1}{2} i \ell n - \frac{1}{8} \ell^2 n^2 \right) - \frac{1}{24} \ell^2_{uv} \right] e^{-\frac{1}{2} i \ell v^2 / u} \right. \\
 & + \left[ v \left( -\frac{1}{2} i \ell + \frac{1}{4} \ell^2 n \right) + \frac{1}{48} i \ell^3 v^3 \right] \left[ C \left( \frac{\ell v^2}{2u} \right) \right. \\
 & \left. \left. - i S \left( \frac{\ell v^2}{2u} \right) \right] \right\} \Bigg|_{v = m - \frac{1}{2}}^{m + \frac{1}{2}} \quad (43)
 \end{aligned}$$

To evaluate  $C_n(u, m)$ , put  $v = s + m$ . Expanding part of the exponential gives the approximation

$$\begin{aligned}
 \int_{m-\frac{1}{2}}^{m+\frac{1}{2}} dv e^{-\frac{1}{2} i \ell v^2 / u} &= \int_{-\frac{1}{2}}^{\frac{1}{2}} ds e^{-\frac{1}{2} i \ell (m^2 + 2ms + s^2) / u} \\
 &\cong e^{-\frac{1}{2} i \ell m^2 / u} \int_{-\frac{1}{2}}^{\frac{1}{2}} ds e^{-i \ell ms / u} \left( 1 - \frac{1}{2} i \ell s^2 / u \right)
 \end{aligned}$$

and performing the integration gives

$$\begin{aligned}
 \int_{m-\frac{1}{2}}^{m+\frac{1}{2}} dv e^{-\frac{1}{2} i \ell v^2 / u} &= e^{-\frac{1}{2} i \ell m^2 / u} \left\{ \frac{\sin(\ell m / 2u)}{\ell m / 2u} \left( 1 - \frac{1}{8} \frac{i \ell}{u} \right) \right. \\
 &\quad \left. + \frac{i u}{\ell m^2} \left[ \frac{\sin(\ell m / 2u)}{\ell m / 2u} - \cos(\ell m / 2u) \right] \right\}, \quad m \neq 0
 \end{aligned}$$

For small values of  $\ell m / 2u$ , the trigonometric functions of this argument are expanded in power series. To sufficient accuracy,

$$\int_{m-\frac{1}{2}}^{m+\frac{1}{2}} dv e^{-\frac{1}{2} i \ell v^2 / u} = e^{-\frac{1}{2} i \ell m^2 / u} \left[ 1 - \frac{1}{6} \left( \frac{\ell m}{2u} \right)^2 - \frac{1}{24} \frac{i \ell}{u} \right], \quad \ell m / 2u < 0.2$$

# Contrails

It may be verified that this is valid for  $m = 0$ .

Combining these expressions,

$$C_n(u, m) = e^{-\frac{1}{2} i \ell m^2 / u} \left[ -\frac{2}{u} \left( 1 + \frac{1}{2} i \ell n - \frac{1}{8} \ell^2 n^2 \right) + 2 \left( -\frac{1}{2} i \ell + \frac{1}{4} \ell^2 n \right) - \frac{1}{12} \ell^2 u \right]$$
$$\left\{ \begin{array}{l} \left( 1 - \frac{1}{8} \frac{i \ell}{u} \right) \frac{\sin(\ell m / 2u)}{\ell m / 2u} + \frac{i u}{\ell m^2} \left[ \frac{\sin(\ell m / 2u)}{\ell m / 2u} - \cos(\ell m / 2u) \right], \\ \ell m / 2u > 0.2 \\ \\ 1 - \frac{1}{24} \frac{i \ell}{u} - \frac{1}{6} \left( \frac{\ell m}{2u} \right)^2, \quad \ell m / 2u < 0.2 \end{array} \right. \quad (44)$$

The subroutine POT2 evaluates the influence coefficients according to Equations (41), (42), (43), and (44).

# *Contrails*

**APPENDIX III. COMPUTER PROGRAM LISTINGS**





00000770  
00000780  
00000790  
00000800  
00000810  
00000820  
00000830  
00000840  
00000850  
00000860  
00000870  
00000880  
00000890  
00000900  
00000910  
00000920  
00000930  
00000940  
00000950  
00000960  
00000970  
00000980  
00000990  
00001000  
00001010  
00001020  
00001030  
00001040  
00001050  
00001060  
00001070  
00001080  
00001090  
00001100  
00001110  
00001120  
00001130

```
35 CALL SHAPE
   CALL FORCI
27 LIM=ML(L)
   IF (LIM-50) 22,22,101
22 LIM2=2*LIM
   LPOT= MINO(L,15)
   CALL POT2(100,LIM2,LPOT,CK,D)
   M=0
   K=DA(28)
   GO TO 4
```

C  
C PRELIMINARY CALCULATIONS ARE FINISHED.  
C THE NEXT SECTION IS GONE THROUGH FOR EACH MODE.  
C

```
3 IF (DA(26)) 4,28,4
28 CALL DATRD(DA)
4 K=K-1
   M=M+1
   WRITE
0 (6, 48)M
48 FORMAT (1H115X,8HMODE NO.13)
   WRITE
0 (6, 49)(DA(I),I=13,22)
49 FORMAT (1H010X,12A6)
   IF (DA(26)) 29,7,29
7 CALL DRED
9 G(M)=DA(39)
29 IF (G(M)) 74,72,74
72 CALL WVAL
   IF (ML) 24,26,24
24 LIM2=2*ML
00 25 J=1,LIM2
25 W(J,1,1)=W(J,1,1)*2.0/3.14159265
   LEADING EDGE CORRECTION
26 CONTINUE
   CALL BOXP
74 IF (K) 6,6,5
```

C

```

00001140
00001150
00001160
00001170
00001180
00001200
00001220
00001230
00001240
00001250
00001260
00001270
00001280
00001290
00001300
00001310
00001320
00001330
00001340
00001350
00001360
00001370
00001380
00001390
00001400
00001410
00001420
00001430
00001440
00001460
00001470
00001480
00001490
00001500
00001510
00001520

5 IF (M-10) 3,6,6
C
C
C
FINAL SECTION OF PROGRAM - COMPUTATION OF GENERALIZED FORCES
6 WRITE (6,46)
46 FORMAT (1H110X,18HGENERALIZED FORCES/1H05X,5HM0DES/4X,11H0SC, DEF00001200
1L.8X,9HREAL PART10X,9HIMAG PART10X,10HABS. VALUE6X,11HPHASE ANGLE)00001210
AC=8.0/AREA
D0 12 M1=1,M
IF (G(M1)) 12,14,12
14 D0 18 M2=1,M
S1=0.0
S2=0.0
N1=5*(M1-1)
N2=5*(M2-1)
D0 8 J=1,5
J1=J+N1
J2=J+N2
D0 8 I=1,5
S1=S1+PS(1,I,J1)*DF(I,J2)
8 S2=S2+PS(2,I,J1)*DF(I,J2)
S1=AC*S1
S2=AC*S2
S3= SQRT(S1**2+S2**2)
S4= ATAND(S2,S1)
18 WRITE
0 (6,47)M1,M2,S1,S2,S3,S4
47 FORMAT (1H02I6,1P3E19.5,0P1F16.4)
12 CONTINUE
WRITE (6,43)
43 FORMAT(1H1)
GO TO 16
C
C
C
ERROR EXITS
83 IPR=27
84 WRITE

```

00001530  
00001540  
00001550  
00001560  
00001580  
00001600  
00001610

```
0 (6, 45)IPR
45 FORMAT(1H010X, 8H8AD DATAI4)
GO TO 102
101 WRITE ( 6, 56)
56 FORMAT(1H010X, 42HLATERAL LIMIT ON NUMBER OF BOXES EXCEEDED.)
102 STOP
END
```

```

SUBROUTINE SHAPE
SUBROUTINE SHAPE
DIMENSION A(2,100,16),W(2,50,50),DA(700),PS(2,5,50),DF(5,50)
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),CDE(5,5)
COMMON A,W,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,CDE,M,L,NS,D,DI,CK,IEDG
COMMON AREA,DH
IEDG=0
NS=DA(29)
IF (NS) 81,81,1
1 IF (NS-3) 2,2,81
2 NSP=NS+1
IF (DA(24)) 82,82,3
3 DO 4 I=1,NSP
XEDG(I)=DA(2*I+27)/DA(24)
4 YEDG(I)=DA(2*I+28)/DA(24)
XEDG=0.0
XEDG(NS+2)=1.0
YEDG(NS+2)=YEDG(NS+1)
Y1=DA(30)
IF (Y1) 83,20,20
20 K=0
X=DH
AREA=0.0
DO 15 I=1,L
7 IF (X-XEDG(K+1)) 8,8,9
8 ML(I)=0.5+DI*(YEDG(K)+F*(X-XEDG(K))/G)
GO TO 15
9 K=K+1
G=XEDG(K+1)-XEDG(K)
F=YEDG(K+1)-YEDG(K)
IF (G) 84,12,12
12 IF (F) 85,13,13
13 AREA=AREA+G*(YEDG(K+1)+YEDG(K))
GO TO 7
15 X=X+D
IF (XEDG(NS+1)-1.0) 17,16,84
16 IF (YEDG(NS+1)-YEDG(NS)) 17,17,18
17 IEDG=1
00001640
00001650
00001660
00001670
00001680
00001690
00001700
00001710
00001720
00001730
00001740
00001750
00001760
00001770
00001780
00001790
00001800
00001810
00001820
00001830
00001840
00001850
00001860
00001870
00001880
00001890
00001900
00001910
00001920
00001930
00001940
00001950
00001960
00001970
00001980
00001990
00002000

```

SUBROUTINE SHAPE

```
18 RETURN  
81 IPR=29  
   GO TO 86  
82 IPR=24  
   GO TO 86  
83 IPR=30  
   GO TO 86  
84 K= MING(K,NS)  
   IPR=2*K+29  
   GO TO 86  
85 IPR=2*K+30  
86 WRITE  
   G (6, 4) IPR  
41 FORMAT(1H010X,8HBAD DATA13)  
STOP  
END
```

00002010  
00002020  
00002030  
00002040  
00002050  
00002060  
00002070  
00002080  
00002090  
00002100  
00002110  
00002120  
00002130  
00002140  
00002160  
00002170

```

SUBROUTINE FORCI
CONTROL SUBROUTINE FOR CALCULATION OF INTEGRALS USED
TO FIND GENERALIZED FORCES
SUBROUTINE FORCI
DIMENSION A(2,100,16),W(2,50,50),DA(700),PS(2,5,50),DF(5,50)
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)
COMMON A,W,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,D,DI,CK,IEDG
COMMON AREA,DH
COMMON YMAX2,N
YMAX2=YEDG(NS+1)**2
DO 3 J=1,9
DO 2 I=1,9
2 AXY(I,J)=0.0
3 AY(J)=0.0
4 IF (DA(30)) 4,5,4
N=0
CALL SECT
SECT DOES THE CALCULATIONS FOR EACH SECTION OF THE PLANFORM
DO 7 I=1,NS
IF (YEDG(I)-YEDG(I+1)) 6,7,7
6 N=I
CALL SECT
7 CONTINUE
RETURN
END

```

```

00002200
00002210
00002220
00002230
00002240
00002250
00002260
00002270
00002280
00002290
00002300
00002310
00002320
00002330
00002340
00002350
00002360
00002370
00002380
00002390
00002400
00002410
00002420
00002430
00002440
00002450
00002460
00002470

```



SUBROUTINE SECT

```

H(4)=H(3)
H(5)=H(2)
H(6)=H(1)
C
C GAUSSIAN POINTS AND WEIGHTS FOR THE INTERVAL (0,1)
C
X2=XEDG(N+1)
X1=XEDG(N)
Y2=YEDG(N+1)
Y1=YEDG(N)
IF (N) 4,2,4
2 X1=0.0
Y1=0.0
4 DY=Y2-Y1
DG 19 J=1,6
V=U(J)
G=H(J)*DY
IF (Y2**2-YMAX2) 7,6,7
6 G=2.0*V*G
V=V*V
7 Y=Y2-V*DY
X0=X2+V*(X1-X2)
XOP=1.0-X0
G=G* SQRT(XOP)
YQ=Y*Y
IF (IEDG) 8,9,8
8 G=G* SQRT(1.0-YQ/YMAX2)
9 E=2.0*XOP*G
IF (DA(30)) 10,10,11
10 G=G* SQRT(1.0+X0)
11 DG 17 I=1,6
U2=U(I)**2
X=X0+XOP*U2
F=E*H(I)*U2
IF (DA(30)) 12,12,13
12 F=F* SQRT(X+X0)
13 YP=1.0
00002870
00002880
00002890
00002900
00002910
00002920
00002930
00002940
00002950
00002960
00002970
00002980
00002990
00003000
00003010
00003020
00003030
00003040
00003050
00003060
00003070
00003080
00003090
00003100
00003110
00003120
00003130
00003140
00003150
00003160
00003170
00003180
00003190
00003200
00003210
00003220
00003230

```



SUBROUTINE SECT

00003240  
00003250  
00003260  
00003270  
00003280  
00003290  
00003300  
00003310  
00003320  
00003330  
00003340  
00003350  
00003360  
00003370

```
DO 16 M=1,9  
XP=YP  
DO 15 L=1,9  
  15 XY(L,M)=AXY(L,M)+XP*F  
  16 XP=X*XP  
  17 CONTINUE  
  18 YP=YQ*YP  
  19 CONTINUE  
  20 YP=1.0  
  21 DO 18 M=1,9  
  22 AY(M)=AY(M)+YP*G  
  23 YP=YQ*YP  
  24 CONTINUE  
  25 RETURN  
  26 END
```

```

SUBROUTINE DRED
SUBROUTINE DRED
DIMENSION A(2,100,16),W(2,50,50),DA(700),PS(2,5,50),DF(5,50)
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)
COMMON A,W,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,O,DI,CK,IEDG
COMMON AREA,DH
DIMENSION IXB(5),B(25),C(25,25),G(25)
COMMON C,B,G,IXB
NP=DA(98)
IF (NP) 83,24,30
C
C A POLYNOMIAL FOR THE DEFLECTION IS FITTED TO VALUES
C OF DEFLECTION AT GIVEN POINTS.
C
30 NX=DA(99)
NY=DA(100)
IF (NX) 81,81,31
31 IF (NY) 82,82,32
32 IF (150-NP) 83,38,38
38 IF (NP-NX) 81,34,34
34 IF (NP-NY) 82,35,35
35 MX= MINO(NX,5)
MY= MINO(NY,5)
IY=MY
IXY=MX+MY
D0 2 J=1,5
D0 1 I=1,5
1 COE(I,J)=0.0
2 IXB(J)=MX
NC=MX*MY
3 D0 5 I=1,NC
D0 4 J=1,NC
4 C(I,J)=C.0
5 B(I)=0.0
KP=100
D0 11 IP=1,NP
X=DA(KP+1)/DA(24)
Y2=(DA(KP+2)/DA(24))**2
00003400
00003410
00003420
00003430
00003440
00003450
00003460
00003470
00003480
00003490
00003500
00003510
00003520
00003530
00003540
00003550
00003560
00003570
00003580
00003590
00003600
00003610
00003620
00003630
00003660
00003662
00003664
00003670
00003680
00003690
00003700
00003710
00003720
00003730
00003740
00003750
00003760

```

SUBROUTINE DRED

```
DEF=DA(KP+3)
WT=DA(KP+4)
IF (WT) 84,84,6
6 YP=1.0
  K=1
  DG 8 J=1,IY
  XYP=YP
  JX=IXB(J)
  DG 7 I=1,JX
  G(K)=XYP
  XYP=X*XYP
7 K=K+1
8 YP=Y2*YP
  DG 10 I=1,NC
  DG 9 J=1,NC
9 C(I,J)=C(I,J)+G(I)*G(J)*WT
10 B(I)=B(I)+G(I)*DEF*WT
11 KP=KP+4
  K=MSIMER(25,NC,1,C,B)
  IF (K-1) 22,22,15
15 DG 16 I=1,IY
  IP=IY+1-I
  IF (IXB(IP)+IP-IXY) 16,17,17
16 CONTINUE
  IXY=IXY-1
  GO TO 15
17 IXB(IP)=IXB(IP)-1
  IF (IXB(IP)) 18,18,19
18 IY=IP-1
19 NC=0
  DG 20 I=1,IY
  NC=NC+IXB(I)
  GO TO 3
22 K=1
  DG 23 J=1,IY
  JX=IXB(J)
  DG 23 I=1,JX
```

00003770  
00003780  
00003790  
00003800  
00003810  
00003820  
00003830  
00003840  
00003850  
00003860  
00003870  
00003880  
00003890  
00003900  
00003910  
00003920  
00003930  
00003940  
00004030  
00004040  
00004050  
00004060  
00004070  
00004080  
00004090  
00004100  
00004110  
00004120  
00004130  
00004140  
00004150  
00004160  
00004170  
00004180  
00004190  
00004200  
00004210

SUBROUTINE DRED

```

C0E(I,J)=8(K)
23 K=K+1
61 IF (DA(87)) 61,66,61
61 WRITE ( 6,41)
D0 64 J=1,IY
JX=IXB(J)
D0 63 I=1,JX
WRITE
0 (6, 42)I,J,C0E(I,J)
63 CONTINUE
64 CONTINUE
66 CONTINUE
GO TO 28
24 YP=1.0
K=1
D0 27 J=1,5
XYP=YP
D0 26 I=1,5
C0E(I,J)=XYP*DA(K+45)
K=K+1
26 XYP=XYP*DA(24)
27 YP=YP*DA(24)**2
28 K=25*(M-1)
D0 29 I=1,25
K=K+1
29 DF(K,1)=C0E(I,1)
RETURN
81 IPR=99
GO TO 85
82 IPR=100
GO TO 85
83 IPR=98
GO TO 85
84 IPR=K+4
85 WRITE
0 (6, 45)IPR
STOP
00004220
00004230
00004240
00004250
00004270
00004280
00004290
00004300
00004310
00004320
00004330
00004340
00004350
00004360
00004370
00004380
00004390
00004400
00004410
00004420
00004430
00004440
00004450
00004460
00004470
00004480
00004490
00004500
00004510
00004520
00004530
00004540
00004550
00004560
00004570
00004580
00004600
```

```
SUBROUTINE DRED
  41 FORMAT(1H010X,56HCOMPUTED DEFLECTION = SUM OF DEF(N,M)*X**(N-1)*Y*00004610
    1*(2M-2)/1H010X,54H(IN DIMENSIONLESS COORDINATES - DISTANCE/CHORD L000004620
    2ENGTH)/1H09X,1HN7X,1HM16X,8HDEF(N,M))
  42 FORMAT(3X,2I8,1PE25.5)
  45 FORMAT(1H010X,8HBAD DATAI4)
  END
```

```

SUBROUTINE MVAL
SUBROUTINE MVAL
EVALUATION OF THE UPWASH ARRAY
DIMENSION A(2,100,16),W(2,50,50),DA(700),PS(2,5,50),DF(5,50)
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)
COMMON A,W,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,D,DI,CK,IEDG
COMMON AREA,DH
DIMENSION G(5,5),H(5,5)
COMMON G,H
J1=5*(M-1)
DO 3 J=1,5
  J1=J1+1
  CI=1.0
  DO 2 I=1,5
    G(I,J)=CI*DF(I+1,J1)
    H(I,J)=CK*DF(I,J1)
  2 CI=CI+1.0
  3 G(5,J)=0.0

G AND H ARE THE COEFFICIENTS OF THE REAL AND IMAGINARY PARTS
OF THE UPWASH
X=DH
DO 10 I=1,L
  JL=ML(I)
  IF (JL) 5,10,5
  5 Y=DH
  DO 9 J=1,JL
    Y2=Y*Y
    W(1,J,I)=0.0
    W(2,J,I)=0.0
    YP=1.0
    DO 8 J1=1,5
      XYP=Y*Y
      DO 7 I1=1,5
        W(1,J,I)=W(1,J,I)+G(I1,J1)*XYP

```

```

00004690
00004700
00004710
00004720
00004730
00004740
00004750
00004760
00004770
00004780
00004790
00004800
00004810
00004820
00004830
00004840
00004850
00004860
00004870
00004880
00004890
00004900
00004910
00004920
00004930
00004940
00004950
00004960
00004970
00004980
00004990
00005000
00005010
00005020
00005030
00005040
00005050

```

SUBROUTINE WVAL

```
W(2,J,I)=W(2,J,I)+H(I1,J1)*XYP
7 XYP=X*XYP
8 YP=Y2*YP
9 Y=Y+D
10 X=X+D
RETURN
END
```

00005060  
00005070  
00005080  
00005090  
00005100  
00005110  
00005120

```

00005150
00005160
00005170
00005180
00005190
00005200
00005210
00005220
00005230
00005240
00005250
00005270
00005280
00005290
00005300
00005310
00005320
00005330
00005340
00005350
00005360
00005370
00005380
00005390
00005400
00005410
00005420
00005430
00005440
00005450
00005460
00005470
00005480
00005490
00005500
00005510
00005520

SUBROUTINE BOXP
SUBROUTINE BOXP
DIMENSION A(2,100,16),S(2,50,50),DA(700),PS(2,5,50),DF(5,50)
DIMENSION ML(50),AXY(9,9),AY(9),XEDG(5),YEDG(5),COE(5,5)
DIMENSION EDG(50),PR(2,50),PSI(2,50),G(20),XO(50),IXB(4)
DIMENSION C(20,20),B(20,2)
COMMON A,S,DA,PS,DF,ML,AXY,AY,XEDG,YEDG,COE,M,L,NS,D,DI,CK,IEDG
COMMON AREA,DH
COMMON B,C,EDG,PR,PSI,G,XO,IXB
C
IF (DA(88)) 71,75,71
71 WRITE ( 6,47)
47 FORMAT(1H11OX,43HTHE UPWASH ARRAY (REAL AND IMAGINARY PARTS))
D0 74 I=1,L
JL=ML(I)
IF (JL) 73,74,73
73 WRITE
0 (6, 42)I
IF (I-1) 77,77,79
77 D0 78 J=1,JL
S1=S(1,J,I)*1.57079633
S2=S(2,J,I)*1.57079633
78 WRITE
0 (6, 142)S1,S2
G0 T0 74
79 WRITE
0 (6, 142)(S(1,J,I),S(2,J,I),J=1,JL)
142 FORMAT(1H 1P2E24.5)
74 CONTINUE
C THESE ARE THE UPWASHES
C
75 CONTINUE
CALL BOXPO
C BOXPO COMPUTES THE POTENTIAL VALUES IN EACH BOX.
C THEY ARE STORED IN THE ARRAY S.
C
IF (DA(89)) 91,95,91
91 WRITE ( 6,143)

```



```

SUBROUTINE BQXP
143 FORMAT(1H11OX,46HTHE POTENTIAL ARRAY (REAL AND IMAGINARY PARTS))
DO 94 I=1,L
  JL=ML(I)
  IF (JL) 93,94,93
93 WRITE
   G (6, 42)I
  WRITE
   G (6, 142)(S(1,J,I),S(2,J,I),J=1,JL)
94 CONTINUE
  THESE ARE THE POTENTIALS
95 CONTINUE
  FIT OF A SERIES TO THE POTENTIAL VALUES
  JL=ML(L)
  DY=D/YEDG(NS+1)
  Y=0.5*DY
  DO 201 J=1,JL
  IF (IEDG) 202,203,202
202 EDG(J)= SQRT(1.0-Y*Y)
  Y=Y+DY
  GO TO 201
203 EDG(J)=1.0
201 CONTINUE
  N=0.5+DI*YEDG
  IF (N) 4,4,2
  2 DO 3 I=1,N
  3 X0(I)=0.0
  4 X1=0.0
  N1=N
  Y1=YEDG
  DO 8 K=1,NS
  X2=XEDG(K+1)
  Y2=YEDG(K+1)
  N=DI*Y2+0.5
  IF (N1-N) 5,7,7
  5 N1=N1+1
00005540
00005550
00005560
00005570
00005580
00005590
00005600
00005610
00005620
00005630
00005640
00005650
00005660
00005670
00005680
00005690
00005700
00005710
00005720
00005730
00005740
00005750
00005760
00005770
00005780
00005790
00005800
00005810
00005820
00005830
00005840
00005850
00005860
00005870
00005880
00005890
00005900

```

00005910  
00005920  
00005930  
00005940  
00005950  
00005960  
00005970  
00005980  
00005990  
00006000  
00006010  
00006020  
00006030  
00006040  
00006050  
00006060  
00006070  
00006080  
00006090  
00006100  
00006110  
00006120  
00006130  
00006140  
00006150  
00006160  
00006170  
00006180  
00006190  
00006200  
00006210  
00006220  
00006230  
00006240  
00006250  
00006260  
00006270

SUBROUTINE BOXP

```

00 6 I=N1,N
    Y=D*( FLOAT(I)-0.5)
    6 X0(I)=X1+(X2-X1)*(Y-Y1)/(Y2-Y1)
    7 X1=X2
      Y1=Y2
    8 N1=N
      AS=0.0
    00 119 I=1,L
      JL=ML(I)
    217 IF (JL) 119,119,217
    118 D0 118 J=1,JL
      AS=AMAX1(AS, ABS(S(1,J,I)), ABS(S(2,J,I)))
    119 CONTINUE
      IX=5
      IY=4
      IXY=9
    13 D0 14 I=1,IY
    14 IXB(I)=IX
    16 NC=0
    17 D0 17 I=1,IY
      NC=NC+IXB(I)
    18 D0 19 I=1,NC
      D0 18 J=1,NC
    19 C(I,J)=0.0
      B(I,1)=0.0
    20 D0 25 I=1,L
      X1=0.5*D
      JL=ML(I)
      Y=0.5*D
      IF (JL) 25,25,20
    601 D0 24 J=1,JL
      XR=X1-X0(J)
      IF (YEDG) 601,601,602
    602 XR=XR*(X1+X0(J))
      XR= SQRT(XR)

```

SUBROUTINE BQXP

```

YP=1.0
K=1
D0 23 N1=1,IY
XP=XR*YP
JX=IXB(N1)
D0 22 N=1,JX
G(K)=XP*EDG(J)
XP=X1*XP
22 K=K+1
23 YP=Y*Y*YP
Y=Y+D
D0 24 N1=1,NC
D0 124 N=1,NC
124 C(N1,N)=C(N1,N)+G(N1)*G(N)
D0 24 N=1,2
24 B(N1,N)=B(N1,N)+G(N1)*S(N,J,I)
25 X1=X1+D
K=MSIMER(20,NC,2,C,B)
IF (K-1) 30,30,29
C
C
C
C
IF XSIMEQ FAILS, THE FOLLOWING SECTION REDUCES THE NUMBER OF TERMS
IN THE SERIES.
29 D0 61 I=1,IY
IP=IY+1-I
IF (IXB(IP)+IP-IXY) 61,62,62
61 CONTINUE
IXY=IXY-1
G0 T0 29
62 IXB(IP)=IXB(IP)-1
IF (IXB(IP)) 63,63,16
63 IY=IP-1
G0 T0 16
C
30 K=1
AC=0.0
D0 133 I=1,IY
00006280
00006290
00006300
00006310
00006320
00006330
00006340
00006350
00006360
00006370
00006380
00006390
00006400
00006410
00006420
00006430
00006440
00006540
00006550
00006560
00006570
00006580
00006590
00006600
00006610
00006620
00006630
00006640
00006650
00006660
00006670
00006680
00006690
00006700
00006710
00006720
00006730

```



```

SUBROUTINE B0XP
WRITE ( 6,149)
113 IF (DA(91)) 114,116,114
C
C
C
PRINTOUT OF VALUES OF POTENTIAL AND PRESSURE
114 WRITE ( 6,41)
41 FORMAT(1H010X,9HPOTENTIAL45X,8HPRESSURE)
X1=0.5*D
D0 39 I=1,L
JL=ML(I)
Y=0.5*D
IF (JL) 39,39,34
34 D0 38 J=1,JL
XR=X1-X0(J)
XQ=0.5
IF (YEDG) 608,608,609
608 XQ=X1
XR=XR*(X1+X0(J))
609 XQ=XQ/XR
XR= SQRT(XR)
D0 37 N=1,2
PSI(N,J)=0.0
PR(N,J)=0.0
K=1
YP=EDG(J)
D0 37 N1=1,IY
XPI=XR*YP
JX=IXB(N1)
D0 36 M1=1,JX
PSI(N,J)=PSI(N,J)+C(K,N)*XPI
PR(N,J)=PR(N,J)+B(K,N)*XPI
XPI=X1*XPI
36 K=K+1
37 YP=Y*Y*YP
PR(1,J)=2.0*(PR(1,J)+PSI(1,J)*XQ-PSI(2,J)*CK)
PR(2,J)=2.0*(PR(2,J)+PSI(2,J)*XQ+PSI(1,J)*CK)
38 Y=Y+D
00007150
00007170
00007180
00007190
00007200
00007210
00007230
00007240
00007250
00007260
00007270
00007280
00007290
00007300
00007310
00007320
00007330
00007340
00007350
00007360
00007370
00007380
00007390
00007400
00007410
00007420
00007430
00007440
00007450
00007460
00007470
00007480
00007490
00007500
00007510
00007520
00007530

```



SUBROUTINE BCXP

216X,10HIMAG. PART)  
49 FFORMAT(1H 218,1P2E25.5)  
149 FFORMAT(1H1)  
END

00007910  
00007920  
00007930  
00007940

```

SUBROUTINE BOXPO
C SOLUTION OF SIMULTANEOUS EQUATIONS FOR THE POTENTIAL
COMMON A(2,100,16),S(2,50,50),DA(700),PS(2,5,50),DF(5,50),ML(50)
COMMON AX(9,9),AY(9),XEDG(5),YEDG(5),M,L,NS,D,OI,CK,IEDG
COMMON AREA,DH,E(2,50,50)
I1=0
DO 9 I=1,L
KO=MAXO(1,I-14)
JL=ML(I)
IF (JL.EQ.O) GO TO 9
IF (I1.EQ.O) GO TO 6
SUBTRACTION OF CONTRIBUTIONS OF PRECEDING ROWS TO UPWASH
DO 5 J=1,JL
DO 5 K=KO,I1
KL=ML(K)
K1=I+1-K
IF (KL.EQ.O) GO TO 5
DO 4 N=1,KL
N1=N+J
N2=IABS(N-J)+1
A1=A(1,N1,K1)+A(1,N2,K1)
A2=A(2,N1,K1)+A(2,N2,K1)
S(1,J,I)=S(1,J,I)-A1*S(1,N,K)+A2*S(2,N,K)
4 S(2,J,I)=S(2,J,I)-A2*S(1,N,K)-A1*S(2,N,K)
5 CONTINUE
C SETTING UP MATRIX FOR SIMULTANEOUS EQUATIONS
DO 8 J=1,JL
DO 8 K=1,J
N1=J+K
N2=IABS(J-K)+1
E(1,J,K)=A(1,N1,1)+A(1,N2,1)
E(2,J,K)=A(2,N1,1)+A(2,N2,1)
E(1,K,J)=E(1,J,K)
8 E(2,K,J)=E(2,J,K)
C SOLUTION OF EQUATIONS
K=MSIMEC(50,JL,1,E,S(1,1,I))
IF (K.NE.1) GO TO 12
9 I1=I1+1
10020
10030
10040
10050
10060
10070
10080
10090
10100
10110
10120
10130
10140
10150
10160
10170
10180
10190
10200
10210
10220
10230
10240
10250
10260
10270
10280
10290
10300
10310
10320
10330
10340
10350
10360
10370
10380

```



SUBROUTINE BOXPO

10390  
10400  
10410  
10420  
10430  
10440

RETURN  
12 WRITE ( 6,41)  
41 FORMAT(1H010X,59HSOLUTION OF SIMULTANEOUS EQUATIONS FOR THE POTENT  
IAL FAILED)  
STOP  
END



SUBROUTINE POT2

```

CN=1.0
K=I
C3=0.0
C4=0.0
C7=0.0
C8=0.0
DO 2 J=1,N
A1=DM/CN
C1=CM* COS(A1)
C2=-CM* SIN(A1)
C5=CM*CIN(A1,C6)
C6=-CM*C6
C9=C1-C3
C10=C2-C4
C11=C5-C7
C12=C6-C8
A(1,K)=B3*C9-B4*C10-B5*C3-B1*C11-B2*C12
A(2,K)=B4*C9+B3*C10-B5*C4+B2*C11-B1*C12
23 C3=C1
C4=C2
C7=C5
C8=C6
B1=B1-D1
B3=B3-D3
B4=B4-D4
D4=D4+DD4
CN=CN+2.0
2 K=K+M2
CM=CM+1.0
DM=DM+DDM
3 DDM=DDM+DD
DO 5 L=1,2
KI=1
DO 5 J=1,N
DO 4 I=1,M1
K=KI+M-I
4 A(L,K)=A(L,K)-A(L,K-1)

```

SUBROUTINE PGT2

```

A(L,K1)=2.0*A(L,K1)
5 K1=K1+M2
CM=0.0
DM=0.0
DDM=DK
D0 12 I=1,M
C7=0.0
C8=0.0
C9=0.0
C10=0.0
P1=0.0
P2=0.0
CN=1.0
B6=0.5*DK12
K=I
D0 10 J=1,N
A1=CM/CN
A2=DM/CN
IF (A1-0.2) 7,7,8
7 B1=2.0-A1**2/3.0
B2=-DK/(6.0*CN)
G0 T0 9
8 B3= SIN(A1)/A1
B1=2.0*B3
B2=(B3- COS(A1))/A2-DH/CN*B3
9 B3= COS(A2)/CN
B4= SIN(A2)/CN
C3=B1*B3+B2*B4
C4=B2*B3-B1*B4
B5=DH*CN
C1=B5*C4-2.0*C3
C2=-2.0*C4-B5*C3
C5=C1-C7
C6=C2-C8
P3=P2-B6*CN
P4=P3+2.0*DK12*(CN-1.0)
A(1,K)=A(1,K)+C5-P1*C6+P3*C3-P4*C9
00008710
00008720
00008730
00008740
00008750
00008760
00008770
00008780
00008790
00008800
00008810
00008820
00008830
00008840
00008850
00008860
00008870
00008880
00008890
00008900
00008910
00008920
00008930
00008940
00008950
00008960
00008970
00008980
00008990
00009000
00009010
00009020
00009030
00009040
00009050
00009060
00009070

```

```
00009080
00009090
00009100
00009110
00009120
00009130
00009140
00009150
00009160
00009170
00009180
00009190
00009200
00009210
00009220
00009230
00009240
00009250
00009260
00009270
00009280
00009290
00009300
00009310
00009320
00009330
00009340
00009350
00009360

SUBROUTINE PGT2
A(2,K)=A(2,K)+C6+P1+C5+P3*C4-P4*C10
P1=P1+DH
P2=P2+CN*DK4
CN=CN+2*0
C7=C1
C8=C2
C9=C3
C10=C4
86=B6+DK12
10 K=K+M2
CM=CM+DK
DM=DM+DDM
12 DDM=DDM+DD
D3=CK/(2*0*3.14159265)
M1=M2-M
K=1
A1=0.0
DG 14 J=1,N
C1=D3* SIN(A1)
C2=-D3* COS(A1)
DG 13 I=1,M
DF =A(1,K)*C1+A(2,K)*C2
A(2,K)=A(2,K)+C1-A(1,K)*C2
A(1,K)=DF
13 K=K+1
K=K+M1
14 A1=A1+DH
RETURN
END
```

```

CIN(X,S)
SINE AND COSINE INTEGRAL SUBROUTINE
C
C IF CALLED BY THE STATEMENT C=CIN(X,S)
C AND S ARE THE INTEGRALS OVER T FROM 1 TO INFINITY OF
C COS(XT)/T AND SIN(XT)/T
C
FUNCTION CIN(XI,S)
SG=1.0
X=XI
IF (X) 1,2,2
1 SG=-SG
X=-X
2 X2=X*X
IF (X-1.0) 3,3,4
C
FOR ABS(X) LESS THAN 1 A SERIES EXPANSION IS USED
C
3 V=((X2/98.0-0.6)*.05*X2+1.0)*X2/18.0-1.0)*X+1.57079633
U=((X2/45.0-1.0)*X2/24.0+1.0)*X2/4.0-.577215665-ALOG(X)
GO TO 5
C
FOR ABS(X) GREATER THAN 1 APPROXIMATIONS OF HASTINGS ARE USED
C
4 P=((X2+19.394119)*X2+47.411538)*X2+8.493336)/(((X2+21.361055)
1 *X2+70.376496)*X2+30.038227)*X)
Q=((X2+21.383724)*X2+49.719775)*X2+5.089504)/(((X2+27.177958)
1 *X2+119.918932)*X2+76.707876)*X2)
CO=COS (X)
SI=SIN (X)
U=Q*CO-P*SI
V=P*CO+Q*SI
5 S=V*SG
CIN=U
RETURN
END
00019050
00019060
00019070
00019080
00019090
00019100
00019110
00019120
00019130
00019140
00019150
00019160
00019170
00019180
00019190
00019200
00019210
00019220
00019230
00019240
00019250
00019260
00019270
00019280
00019290
00019300
00019310
00019320
00019330
00019340
00019350
00019360
00019370
00019380
00019390

```

```

      DATRD
      CARD-READ SUBROUTINE 'DATRD(DATA(I))'
      SUBROUTINE DATRD(DATA)
      DIMENSION DRBU(14),DATA(1)
      DATA ATEST/5HALPHA/,DTEST/1H /,ETEST/1H-/
      1 READ ( 5,2) EMIN,ALP,IND,(DRBU(I),I=1,12)
      2 FORMAT(A1,A5,I6,I2A6)
      IF (ALP.EQ.ATEST) GO TO 9
C
      WRITE (99,2) EMIN,ALP,IND,(DRBU(I),I=1,12)
      CARD IS WRITTEN IN INTERNAL BUFFER
      REWIND 99
      IF (ALP.NE.DTEST) GO TO 8
C
      NUMERIC CARD
C
      READ (99,990) (DRBU(I),I=1,5)
      REWIND 99
      DO 5 I=1,5
      IF(DRBU(I))4,6,4
      DATA(IND)=DRBU(I)
      IND=IND+1
      4   GOTO 11
      5
C
      TEST FOR BLANK FIELD
      IF(SIGN (1.0,DRBU(I)))5,5,4
C
      ALPHA CARD
C
      DO 10 I=1,10
      DATA(IND)=DRBU(I)
      IND=IND + 1
      10  IF (EMIN.NE.ETEST) GO TO 1
      11
C
      RETURN IF COLUMN 1 CONTAINS A MINUS SIGN
C
      13  RETURN
      C

```

```

00021410
00021420
00021430
00021435
00021450
00021455
00021460
00021465
00021470
00021475
00021480
00021490
00021500
00021510
00021520
00021530
00021540
00021550
00021560
00021570
00021580
00021590
00021600
00021610
00021620
00021630
00021640
00021650
00021770
00021780
00021790
00021800
00021805
00021810
00021815
00021820
00021825

```

00021830  
00021835  
00021840  
00021845  
00021850  
00021855  
00021860  
00021870  
00021880  
00021890  
00021900  
00021910

```

          DATRD
C      BAD CARD
C      8 READ (99,992) DRBU
        WRITE ( 6,993) DRBU
        WRITE ( 6,991)
        REWIND 99
        STOP
        590 FORMAT(12X,5E12.0)
        991 FORMAT(38H BAD DATA ON THIS CARD. JOB TERMINATED )
        992 FORMAT(14A6)
        993 FORMAT(12HOCARD IMAGE=14A6)
        END

```



SIMULTANEOUS EQUATION SUBROUTINE

```

$BMAP SIMR          K=MSIMER(N,L,LB,A,B)
* SOLVES THE SYSTEM OF EQUATIONS A*X=B.
* TO USE, SET K=MSIMER(N,L,LR,A,B)
* WHERE N IS THE NUMBER OF ROWS FOR WHICH A IS
* DIMENSIONED, AND L IS THE NUMBER OF EQUATIONS.
* LB IS THE NUMBER OF COLUMNS IN B.
* K=1 DENOTES SUCCESSFUL SOLUTION
* K=2 FOR A SINGULAR OR ILL-CONDITIONED MATRIX
* K=3 IF IMPROPER DATA IS GIVEN.
* TO AVOID THIS SIGNAL, L MUST BE POSITIVE AND AT MOST 100,
* N MUST NOT BE LESS THAN L, A MUST NOT INCLUDE A ROW
* OF ZEROS.
* A IS DESTROYED.
* IF K=1, THE SOLUTION IS RETURNED IN B
ENTRY
*
* MSIMER SAVE 1,2,3,4,5,6,7
*
* PROLOGUE
CLA* 3,4
PAX 0,1
TXL E1,1,0
TXH E1,1,100
PCD 0,1
STD A4-1
STD A6-1
STD 2A6-1
STD A16
STD A16+1
STD 2A21
STD A12+1
STD 2A26
STD 2A26+1
STD 2A37
STD 2A37+1

```

```

00050000
00050020
00050030
00050040
00050050
00050060
00050070
00050080
00050090
00050100
00050110
00050120
00050130
00050140
00050150
00050160
00050170
00050180
00050190
00050200
00050210
00050220
00050230
00050240
00050250
00050260
00050270
00050280
00050290
00050300
00050310
00050320
00050330
00050340
00050350
00050360
00050370

```

SIMULTANEOUS EQUATION SUBROUTINE

STD	A12+2	00050380
STD	A20	00050390
STD	A22	00050400
STD	A25+1	00050410
STD	A25+2	00050420
STD	A33	00050430
STD	A36+1	00050440
STD	A36+2	00050450
SXD	A52+1	00050460
TXI	*+1,1,1	00050470
SCD	A32+1	00050480
TXI	*+1,1,-2	00050490
SCD	A22+1,1	00050500
CLA*	5,4	00050510
PAX	0,7	00050520
SXA	1A6-1,7	00050530
SXA	A14,7	00050540
SXA	1A21-1,7	00050550
SXA	A26,7	00050560
SXA	A37,7	00050570
CLA*	4,4	00050580
PAX	0,1	00050590
TXL	E1,1,0	00050600
TXH	E1,1,**	00050610
SXA	A2,1	00050620
SXA	A5-1,1	00050630
SXD	A9,1	00050640
SXD	A12,1	00050650
TXI	*+1,1,-1	00050660
SXD	A7,1	00050670
SXD	A36,1	00050680
SXD	A18,1	00050690
SXD	A38-1,1	00050700
SXD	A21-1,1	00050710
SXD	A23,1	00050720
SXD	A25,1	00050730
SXD	A31,1	00050740

A52



SIMULTANEOUS EQUATION SUBROUTINE

1A6	FMP	B,6	00051120
	STG	B,6	00051130
	LDQ	T	00051140
	TNX	2A6,4,1	00051150
2A6	TXI	1A6,6,-N	00051160
	TNX	A7,1,1	00051170
	TXI	**1,3,1	00051180
	TXI	A2,5,1	00051190
A7	TXH	B7,2,L -1	00051200
*			00051210
*			00051220
*			00051230
			00051240
			00051250
			00051260
			00051270
			00051280
A8			00051290
			00051300
			00051310
			00051320
			00051330
			00051340
A9			00051350
			00051360
			00051370
			00051380
			00051390
A10			00051400
			00051410
			00051420
*			00051430
*			00051440
*			00051450
			00051460
			00051470
			00051480

		SEARCH FOR MAXIMUM PIVOT IN COLUMN
	PXA	0,2
	PAX	0,1
	PXA	0,3
	PAX	0,6
	PXA	0,0
	LDQ	A,6
	LRS	0
	TLQ	**3
	XCA	
	SXA	A10,1
	TXI	**1,1,1
	TXH	**2,1,L
	TXI	A8,6,-1
	LDQ	TGL
	TLQ	**2
	TRA	E3
	AXT	**1
	SXD	**1,2
	TNX	A17,1,**
		ROW INTERCHANGE
	PXA	0,3
	PAX	0,6
	PAX	0,7

SIMULTANEOUS EQUATION SUBROUTINE

	SCD	**1,1	00051490
	TXI	**1,7,**	00051500
	PXA	0,2	00051510
	PAX	0,4	00051520
A11	CLA	A,6	00051530
	LDQ	A,7	00051540
	STQ	A,7	00051550
	STQ	A,6	00051560
	TXI	**1,4,1	00051570
A12	TXH	A13,4,L	00051580
	TXI	**1,6,-N	00051590
	TXI	A11,7,-N	00051600
A13	PXA	0,5	00051610
	PAX	0,6	00051620
	PAX	0,7	00051630
A14	AXT	LB,4	00051640
	SCD	**1,1	00051650
	TXI	**1,6,**	00051660
A15	CLA	B,7	00051670
	LDQ	B,6	00051680
	STQ	B,6	00051690
	STQ	B,7	00051700
	TNX	A17,4,1	00051710
A16	TXI	**1,6,-N	00051720
	TXI	A15,7,-N	00051730
*			00051740
*			00051750
*			00051760
A17	CLA	=1,0	00051770
	FDP	A,3	00051780
	STQ	AM	00051790
A18	TXH	A21,2,L -1	00051800
	PXA	0,2	00051810
	PAX	0,4	00051820
	PXA	0,3	00051830
	PAX	0,6	00051840
A20	TXI	**1,6,-N	00051850

DIVISION OF ROW BY PIVOT

SIMULTANEOUS EQUATION SUBROUTINE

	LDQ	A,6	00051860
	FMP	AM	00051870
	STG	A,6	00051880
	TXI	**+1,4,1	00051890
	TXL	A20,4,L -1	00051900
A21	PXA	0,5	00051910
	PAX	0,6	00051920
	AXT	LB,4	00051930
1A21	LDQ	B,6	00051940
	FMP	AM	00051950
	STG	B,6	00051960
	TNX	**+2,4,1	00051970
2A21	TXI	1A21,6, -N	00051980
*			00051990
*	ROW REDUCTION		00052000
*			00052010
	PXA	0,2	00052020
	PAX	0,1	00052030
	PAX	0,4	00052040
	TNX	A31,1,1	00052050
	PXA	0,3	00052060
	PAX	0,6	00052070
	PAX	0,7	00052080
	STA	A29	00052090
	SXA	A26+1,5	00052100
	SXA	3A26,5	00052110
	TXI	**+1,6,1	00052120
A22	TXI	**+1,7, -N	00052130
	TXI	**+1,3, -N+1	00052140
	SXA	A28,7	00052150
A23	TXH	A26,2,L -1	00052160
	SXA	A27,3	00052170
A24	LDQ	A,6	00052180
	FMP	A,7	00052190
	CHS		00052200
	FAD	A,3	00052210
	STG	A,3	00052220

SIMULTANEOUS EQUATION SUBROUTINE

A25	TXI	*+1,4,1	00052230
	TXH	A27,4,L -1	00052240
	TXI	*+1,3, -N	00052250
	TXI	A24,7, -N	00052260
A27	AXT	**3	00052270
A26	AXT	LB,4	00052280
	AXT	**7	00052290
	TXI	*+1,7,1	00052300
	SXA	*-2,7	00052310
1A26	LDQ	A,6	00052320
	FMP	B,5	00052330
	CHS	B,7	00052340
	FAD	B,7	00052350
	STG	B,7	00052360
	TNX	3A26,4,1	00052370
2A26	TXI	*+1,5, -N	00052380
3A26	TXI	1A26,7, -N	00052390
	AXT	**5	00052400
A28	TNX	A29,1,1	00052410
	AXT	**7	00052420
	PXA	0,2	00052430
	PAX	0,4	00052440
	TXI	*+1,3,1	00052450
	TXI	A23,6,1	00052460
A29	AXT	**3	00052470
A31	TXH	A43,2,L -1	00052480
	PXA	0,2	00052490
	PAX	0,1	00052500
	PAX	0,4	00052510
	PXA	0,3	00052520
	PAX	0,6	00052530
	PAX	0,7	00052540
A32	TXI	*+1,3, -N-1	00052550
	TXI	*+1,6,-1	00052560
A33	TXI	*+1,7, -N	00052570
	SXA	A37+1,5	00052580
	SXA	A41,5	00052590

SIMULTANEOUS EQUATION SUBROUTINE

A34	SXA	A40,3	00052600
A35	SXA	A39,7	00052610
	SXA	A38,3	00052620
	LDQ	A,6	00052630
	FMP	A,7	00052640
	CHS		00052650
	FAD	A,3	00052660
	STO	A,3	00052670
A36	TXI	**1,4,1	00052680
	TXH	A37,4,L -1	00052690
	TXI	**1,3, -N	00052700
A37	TXI	A35,7, -N	00052710
	AXT	LB,4	00052720
	AXT	**7	00052730
	TXI	**1,7,-1	00052740
	SXA	**2,7	00052750
1A37	LDQ	A,6	00052760
	FMP	B,5	00052770
	CHS		00052780
	FAD	B,7	00052790
	STG	B,7	00052800
	TNX	A41,4,1	00052810
2A37	TXI	**1,5, -N	00052820
	TXI	1A37,7, -N	00052830
A41	AXT	**5	00052840
	TXI	**1,1,1	00052850
A38	TXH	A40,1,L -1	00052860
A39	AXT	**3	00052870
	AXT	**7	00052880
	PXA	O,2	00052890
	PAX	O,4	00052900
	TXI	**1,3,-1	00052910
	TXI	A34,6,-1	00052920
A40	AXT	**3	00052930
	TXI	**1,5,-1	00052940
	TXI	A7,2,1	00052950
A43	CLA	=1	00052960





COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

```

$IBMAP SIMC
* SOLVES THE SYSTEM OF COMPLEX EQUATIONS A*X=B.
* TO USE, SET K=MSIMEC(N,L,LB,A,B)
* WHERE N IS THE NUMBER OF ROWS FOR WHICH A IS
* DIMENSIONED, AND L IS THE NUMBER OF EQUATIONS.
* LB IS THE NUMBER OF COLUMNS IN B.
* K=1 DENOTES SUCCESSFUL SOLUTION
* K=2 FOR A SINGULAR OR ILL-CONDITIONED MATRIX
* K=3 IF IMPROPER DATA IS GIVEN.
* TO AVOID THIS SIGNAL, L MUST BE POSITIVE AND AT MOST 100,
* N MUST NOT BE LESS THAN L, A MUST NOT INCLUDE A ROW
* OF ZEROS.
* A IS DESTROYED.
* IF K=1, THE SOLUTION IS RETURNED IN B
*
* ENTRY
*
* MSIMEC SAVE 1,2,3,4,5,6,7
*
* PROLOGUE
*
* CLA* 3,4
* ALS 1
* PAX 0,1
* TXL E1,1,1
* TXH E1,1,200
* PCD 0,1
* STD A4-1
* STD A6-1
* STD 2A6-1
* STD A16
* STD A16+1
* STD 2A21
* STD A12+1
* STD 2A26
* STD 2A26+1
* STD 2A37

```

```

00053240
00053260
00053270
00053280
00053290
00053300
00053310
00053320
00053330
00053340
00053350
00053360
00053370
00053380
00053390
00053400
00053410
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00053550
00053560
00053570
00053580
00053590
00053600
00053610

```

COMPLEX SIMULTANECUS EQUATION SUBROUTINE

STD	2A37+1	00053620
STD	A12+2	00053630
STD	A20	00053640
STD	A22	00053650
STD	A25+1	00053660
STD	A25+2	00053670
STD	A33	00053680
STD	A36+1	00053690
STD	A36+2	00053700
SXD	A52,1	00053710
TXI	*+1,1,2	00053720
SCD	A32,1	00053730
TXI	*+1,1,-4	00053740
SCD	A22+1,1	00053750
CLA*	5,4	00053760
PAX	0,7	00053770
SXA	1A6-1,7	00053780
SXA	A14,7	00053790
SXA	1A21-1,7	00053800
SXA	A26,7	00053810
SXA	A37,7	00053820
CLA*	4,4	00053830
ALS	1	00053840
PAX	0,1	00053850
TXL	E1,1,1	00053860
TXH	E1,1,**	00053870
SXA	A2,1	00053880
SXA	A5-1,1	00053890
SXD	A9,1	00053900
SXD	A12,1	00053910
TXI	*+1,1,-2	00053920
SXD	A7,1	00053930
SXD	A36,1	00053940
SXD	A18,1	00053950
SXD	A38-1,1	00053960
SXD	A21-1,1	00053970
SXD	A23,1	00053980

A52

COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

	SXD	A25,1	00053990
	SXD	A31,1	00054000
	SCD	A51,1	00054010
	SCD	A51+1,1	00054020
	TXI	*+1,1,2	00054030
	CLA	6,4	00054040
	PAC	0,3	00054050
	CLA	7,4	00054060
	PAC	0,5	00054070
A51	TXI	*+1,3,2L -2	00054080
*	FXI	*+1,5,2L -2	00054090
*			00054100
*			00054110
A2			00054120
	AXT	2L,2	00054130
	PXA	0,3	00054140
	PAX	0,6	00054150
	PAX	0,7	00054160
	PXA	0,0	00054170
A3	LDQ	A,6	00054180
	LRS	0	00054190
	TLQ	*+2	00054200
	XCA		00054210
	LDQ	A+1,6	00054220
	LRS	0	00054230
	TLQ	*+2	00054240
	XCA		00054250
A4	TNX	A4,2,2	00054260
	TXI	A3,6,2N	00054270
	TZE	E1	00054280
	ST0	T	00054290
	CLA	=1,0	00054300
	FDP	T	00054310
	STQ	T	00054320
A5	AXT	2L,2	00054330
	FMP	A,7	00054340
	ST0	A,7	00054350

NORMALIZATION OF ROWS

COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

LDQ	T		00054360
FMP	A+1,7		00054370
STG	A+1,7		00054380
LDQ	T		00054390
TNX	A6,2,2		00054400
TXI	A5,7,2N		00054410
PXA	0,5	A6	00054420
PAX	0,6		00054430
AXT	LB,4		00054440
FMP	B,6	1A6	00054450
STG	B,6		00054460
LDQ	T		00054470
FMP	B+1,6		00054480
STG	B+1,6		00054490
LDQ	T		00054500
TNX	2A6,4,1		00054510
TXI	1A6,6,2N		00054520
TNX	A7,1,2	2A6	00054530
TXI	*+1,3,2		00054540
TXI	A2,5,2	A7	00054550
TXH	B7,2,2L -2	*	00054560
		*	00054570
		*	00054580
		*	00054590
			00054600
			00054610
			00054620
			00054630
			00054640
		A8	00054650
			00054660
			00054670
			00054680
			00054690
			00054700
			00054710
			00054720

PXA	0,2	
PAX	0,1	
PXA	0,3	
PAX	0,6	
PXA	0,0	
LDQ	A,6	
LRS	0	
TLC	*+3	
XCA		
SXA	A10,1	
LDQ	A+1,6	
LRS	0	
TLC	*+3	

SEARCH FOR MAXIMUM PIVOT IN COLUMN

COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

	XCA			00054730
	SXA	A10,1		00054740
	TXI	**+1,1,2		00054750
	TXH	**+2,1,2L		00054760
A9	TXI	A8,6,-2		00054770
	LDQ	TOL		00054780
	TLQ	**+2		00054790
	TRA	E3		00054800
A10	AXT	**+1		00054810
	SXD	**+1,2		00054820
	TNX	A17,1,**		00054830
*				00054840
*				00054850
*				00054860
		ROW INTERCHANGE		00054870
	PXA	0,3		00054880
	PAX	0,6		00054890
	PAX	0,7		00054900
	SCD	**+1,1		00054910
	TXI	**+1,7,**		00054920
	PXA	0,2		00054930
A11	PAX	0,4		00054940
	CLA	A,6		00054950
	LDQ	A,7		00054960
	STG	A,7		00054970
	STQ	A,6		00054980
	CLA	A+1,6		00054990
	LDQ	A+1,7		00055000
	STG	A+1,7		00055010
	STQ	A+1,6		00055020
	TXI	**+1,4,2		00055030
A12	TXH	A13,4,2L		00055040
	TXI	**+1,6,2N		00055050
	TXI	A11,7,2N		00055060
A13	PXA	0,5		00055070
	PAX	0,6		00055080
	PAX	0,7		00055090
A14	AXT	LB,4		00055090

COMPLEX SIMULTANEGUS EQUATION SUBROUTINE

	SCD	**1,1	00055100
	TXI	**1,6,**	00055110
A15	CLA	B,7	00055120
	LDQ	B,6	00055130
	STG	R,6	00055140
	STQ	B,7	00055150
	CLA	B+1,7	00055160
	LDQ	B+1,6	00055170
	STG	B+1,6	00055180
	STQ	B+1,7	00055190
A16	TNX	A17,4,1	00055200
	TXI	**1,6,2N	00055210
	TXI	A15,7,2N	00055220
*			00055230
*			00055240
*			00055250
A17	NZT	A,3	00055260
	TRA	B1	00055270
	LDQ	A,3	00055280
	FMP	A,3	00055290
	STG	T	00055300
	LDQ	A+1,3	00055310
	FMP	A+1,3	00055320
	FAD	T	00055330
	STG	T	00055340
	CLA	A,3	00055350
	FDP	T	00055360
	STQ	AM	00055370
	CLA	A+1,3	00055380
	FDP	T	00055390
	STQ	AN	00055400
A18	TXH	A21,2,2L -2	00055410
	PXA	O,2	00055420
	PAX	O,4	00055430
	PXA	O,3	00055440
	PAX	O,6	00055450
A20	TXI	**1,6,2N	00055460

DIVISION OF ROW BY PIVOT

COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

LDQ	A,6	00055470
FMP	AN	00055480
STG	T	00055490
LDQ	A+1,6	00055500
FMP	AM	00055510
FSB	T	00055520
LDQ	A+1,6	00055530
STG	A+1,6	00055540
FMP	AN	00055550
STG	T	00055560
LDQ	A,6	00055570
FMP	AM	00055580
FAD	T	00055590
STG	A,6	00055600
TXI	**1,4,2	00055610
TXL	A20,4,2L -2	00055620
PXA	0,5	00055630
PAX	0,6	00055640
AXT	LB,4	00055650
LDQ	B,6	00055660
FMP	AN	00055670
STG	T	00055680
LDQ	B+1,6	00055690
FMP	AM	00055700
FSB	T	00055710
LDQ	B+1,6	00055720
STG	B+1,6	00055730
FMP	AN	00055740
STG	T	00055750
LDQ	B,6	00055760
FMP	AM	00055770
FAD	T	00055780
STG	B,6	00055790
TXN	**2,4,1	00055800
TXI	1A21,6,2N	00055810
		00055820
		00055830

\* \* \* ROW REDUCTION



COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

	PXA	O,2	00055840
	PAX	O,1	00055850
	PAX	O,4	00055860
	TNX	A31,1,2	00055870
	PXA	O,3	00055880
	PAX	O,6	00055890
	PAX	O,7	00055900
	STA	A29	00055910
	SXA	A26+1,5	00055920
	SXA	3A26,5	00055930
	TXI	*+1,6,2	00055940
A22	TXI	*+1,7,2N	00055950
	TXI	*+1,3,2N -2	00055960
A23	SXA	A28,7	00055970
	TXH	A26,2,2L -2	00055980
A24	SXA	A27,3	00055990
	LDQ	A,6	00056000
	FMP	A,7	00056010
	STG	T	00056020
	LDQ	A+1,6	00056030
	FMP	A+1,7	00056040
	FSB	T	00056050
	FAD	A,3	00056060
	STG	A,3	00056070
	LDQ	A,6	00056080
	FMP	A+1,7	00056090
	STG	T	00056100
	LDQ	A+1,6	00056110
	FMP	A,7	00056120
	FAD	T	00056130
	CHS	A+1,3	00056140
	FAD	A+1,3	00056150
	STG	A+1,3	00056160
	TXI	*+1,4,2	00056170
	TXH	A27,4,2L -2	00056180
A25	TXI	*+1,3,2N	00056190
			00056200

COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

A27	TXI	A24,7,2N	00056210
A26	AXT	**3	00056220
	AXT	LB,4	00056230
	AXT	**7	00056240
	TXI	*+1,7,2	00056250
	SXA	*-2,7	00056260
1A26	LDQ	A,6	00056270
	FMP	B,5	00056280
	STG	T	00056290
	LDQ	A+1,6	00056300
	FMP	B+1,5	00056310
	FSB	T	00056320
	FAD	B,7	00056330
	STG	B,7	00056340
	LDQ	A,6	00056350
	FMP	B+1,5	00056360
	STG	T	00056370
	LDQ	A+1,6	00056380
	FMP	B,5	00056390
	FAD	T	00056400
	CHS		00056410
	FAD	B+1,7	00056420
	STG	B+1,7	00056430
	TNX	3A26,4,1	00056440
2A26	TXI	*+1,5,2N	00056450
	TXI	1A26,7,2N	00056460
3A26	AXT	**5	00056470
	TNX	A29,1,2	00056480
A28	AXT	**7	00056490
	PXA	0,2	00056500
	PAX	0,4	00056510
	TXI	*+1,3,2	00056520
	TXI	A23,6,2	00056530
A29	AXT	**3	00056540
A31	TXH	A43,2,2L -2	00056550
	PXA	0,2	00056560
	PAX	0,1	00056570

COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

	PAX	O,4	00056580
	PXA	O,3	00056590
	PAX	O,6	00056600
	PAX	O,7	00056610
A32	TXI	*+1,3,2N +2	00056620
	TXI	*+1,6,-2	00056630
A33	TXI	*+1,7,2N	00056640
	SXA	A37+1,5	00056650
	SXA	A41,5	00056660
	SXA	A40,3	00056670
	SXA	A39,7	00056680
A34	SXA	A38,3	00056690
A35	LDQ	A,6	00056700
	FMP	A,7	00056710
	STG	T	00056720
	LDQ	A+1,6	00056730
	FMP	A+1,7	00056740
	FSB	T	00056750
	FAD	A,3	00056760
	STG	A,3	00056770
	LDQ	A,6	00056780
	FMP	A+1,7	00056790
	STG	T	00056800
	LDQ	A+1,6	00056810
	FMP	A,7	00056820
	FAD	T	00056830
	CHS		00056840
	FAD	A+1,3	00056850
	STG	A+1,3	00056860
	TXI	*+1,4,2	00056870
A36	TXH	A37,4,2L -2	00056880
	TXI	*+1,3,2N	00056890
	TXI	A35,7,2N	00056900
A37	AXT	LB,4	00056910
	AXT	** ,7	00056920
	TXI	*+1,7,-2	00056930
	SXA	*-2,7	00056940

COMPLEX SIMULTANEOUS EQUATION SUBROUTINE

1A37	LDQ	A,6	00056950
	FMP	B,5	00056960
	STG	T	00056970
	LDQ	A+1,6	00056980
	FMP	B+1,5	00056990
	FSB	T	00057000
	FAD	B,7	00057010
	STG	B,7	00057020
	LDQ	A,6	00057030
	FMP	B+1,5	00057040
	STG	T	00057050
	LDQ	A+1,6	00057060
	FMP	B,5	00057070
	FAD	T	00057080
	CHS		00057090
	FAD	B+1,7	00057100
	STG	B+1,7	00057110
	TNX	A41,4,1	00057120
2A37	TXI	*+1,5,2N	00057130
	TXI	1A37,7,2N	00057140
A41	AXT	**5	00057150
	TXI	*+1,1,2	00057160
A38	TXH	A40,1,2L -2	00057170
A39	AXT	**3	00057180
	AXT	**7	00057190
	PXA	0,2	00057200
	PAX	0,4	00057210
	TXI	*+1,3,-2	00057220
	TXI	A34,6,-2	00057230
A40	AXT	**3	00057240
	TXI	*+1,5,-2	00057250
	TXI	A7,2,2	00057260
A43	CLA	=1	00057270
	TRA	MSIMEC+1	00057280
B1	STZ	AM	00057290
	CLA	=1.0	00057300
	FDP	A+1,3	00057310



# *Contrails*

APPENDIX IV. SAMPLE DATA SHEETS

The following pages are sample data sheets for a computer run on three modes at three frequencies. The potential will not be computed for the first mode. The generalized forces found will be  $L_{21}$ ,  $L_{22}$ ,  $L_{23}$ ,  $L_{31}$ ,  $L_{32}$ ,  $L_{33}$ .

Of the first fourteen cards, the cards numbered 6, 9, 10, 11, 12, 13, 14 do nothing and are included only to indicate how all data is entered. Cards 1 through 14 are complete in this respect, and all later cards are of the same type as one of the first fourteen. The data used in the least squares surface fit for the deflection is entered in locations 101 through 700.

Card number 22 represents 56 cards for the intermediate data points which would have to be included in an actual run.

FORTRAN FIXED IO DIGIT DECIMAL DATA

DECK NO. \_\_\_\_\_ PROGRAMMER \_\_\_\_\_ DATE \_\_\_\_\_ PAGE \_\_\_\_\_ of \_\_\_\_\_ JOB NO. \_\_\_\_\_

NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1		
13		title - columns 13 - 72 (optional)
25		
37		
49		
61		
1		
13		mode title - columns 13 - 72 (optional)
25		
37		
49		
61		
1		
13		frequency, $\nu$ use constant
25		root chord length, b unite here and
37		speed of sound, a in items 30 - 36
49		Indicator for Subsequent frequencies
61		
1		
13		no of boxes along root chord, L
25		no. of modes to be given
37		no. of sections of leading edge given, NS
49		
61		



FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO.	PROGRAMMER	DATE	PAGE	of	JOB NO.
NUMBER	IDENTIFICATION	DESCRIPTION	DO NOT KEY PUNCH		
1					
13					
25					
37					
49					
61					
1					
13					
25					
37					
49					
61					
1					
13					
25					
37					
49					
61					
1					
13					
25					
37					
49					
61					

coordinates of points on leading edge

$\bar{y}_0$

$\bar{x}_1$

$\bar{y}_1$

$\bar{x}_2$

$\bar{y}_2$

$\bar{x}_3$

$\bar{y}_3$

$\bar{x}_4$

$\bar{y}_4$

$\bar{x}_5$

$\bar{y}_5$

$\bar{x}_6$

$\bar{y}_6$

$\bar{x}_7$

$\bar{y}_7$

$\bar{x}_8$

$\bar{y}_8$

$\bar{x}_9$

$\bar{y}_9$

used if  $NS \geq 2$

used if  $NS = 3$

non zero value suppresses calculation of potential for this mode

beginning of list of coefficients of deflection series

$d_{00}$

$d_{10}$

$d_{20}$

$d_{30}$

$d_{40}$

$\sum d_{ms}^2 y_{2m}^2$

## FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO.	PROGRAMMER	DATE	PAGE	of	JOB NO.
NUMBER		IDENTIFICATION		DESCRIPTION DO NOT KEY PUNCH	
1	5 1				
13					d01
25					d11
37					d12
49					etc
61			0 9		
1	8 7				
13					indicator for deflection coefficient printout
25					indicator for upwash array printout
37					indicator for potential array printout
49					indicator for potential coefficient printout
61			1 0		indicator for potential and pressure printout
1	9 8				
13					NP number of points at which deflection is given
25					NX number of x values
37					NY number of y values
49					
61			1 1		
1	1 0 1				
13					$\hat{x}_1$ data for first point
25					$\hat{y}_1$
37					$f(\hat{x}_1, \hat{y}_1)$
49					weight (optional - weight 1.0 used if no entry is made)
61			1 2		

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FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. \_\_\_\_\_ PROGRAMMER \_\_\_\_\_ DATE \_\_\_\_\_ PAGE \_\_\_\_\_ of \_\_\_\_\_ JOB NO. \_\_\_\_\_

NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1		
13		$\bar{x}_2$ data for second point
25		$\bar{y}_2$
37		$r(\bar{x}_2, \bar{y}_2)$
49		weight
61		
1		
13		$\bar{x}_3$ data for third point
25		$\bar{y}_3$
37		etc
49		
61		
1		
13		beginning of data for second mode
25		
37		
49		
61		
1		
13		zero cancels previous value, causing potential to be computed
25		
37		
49		
61		

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## FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO.	PROGRAMMER	DATE	PAGE	of	JOB NO.
NUMBER		IDENTIFICATION		DESCRIPTION DO NOT KEY PUNCH	
1					
13	8 7				
25	1				
37	1				
49	1				
61	1				
1	4 6				
13	0 . 0				
25	0 . 0 . 4				
37					
49					
61					
1					
13					
25					
37					
49					
61					
1					
13					
25					
37					
49					
61					
1					
13					
25					
37					
49					
61					
1					
13					
25					
37					
49					
61					

Last card for second mode  
this entry cancels previous value of 400

the polynomial is  $0.04\tilde{x}$ , normalised for maximum value 1.0

beginning of data for third mode



FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. \_\_\_\_\_ PROGRAMMER \_\_\_\_\_ DATE \_\_\_\_\_ PAGE \_\_\_\_\_ of \_\_\_\_\_ JOB NO. \_\_\_\_\_

NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1		
13		
25		
37		
49		
61		
1		
13		
25		
37		
49		
61		
1		
13		
25		
37		
49		
61		
1		
13		
25		
37		
49		
61		
1		
13		
25		
37		
49		
61		

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FORTRAN FIXED IO DIGIT DECIMAL DATA

DECK NO.	PROGRAMMER	DATE	PAGE	of	JOB NO.	DESCRIPTION DO NOT KEY PUNCH
1						
13	2.3					
25	4 0 . 0					
37						
49						
61						
1						
13						
25						
37						
49						
61						
1						
13						
25						
37						
49						
61						
1						
13						
25						
37						
49						
61						

data for a third frequency

73 . 80

2 5

73 . 80

73 . 80

73 . 80

UNCLASSIFIED

Security Classification

DOCUMENT CONTROL DATA - R&D

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13. ABSTRACT The fundamental equations of the transonic box method were derived, based on the representation of the velocity potential by a doublet distribution. They form the basis of a systematic method of treating an oscillating wing at M=1, analogous to the supersonic Mach box method. A digital computer program, written in Fortran IV, is presented. The program applies to a planar wing of polygonal planform, with a straight trailing edge, and as many as three sweep angles along the leading edge. For a maximum of ten modes of oscillation, the program computes the oscillatory potentials and pressures and a generalized force matrix. Results obtained from the program are compared with existing theoretical and experimental values. Several possible extensions of the method are described.			

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